



UNIVERSITÀ  
DEGLI STUDI DI BARI  
ALDO MORO



## 107° CONGRESSO NAZIONALE della SOCIETÀ ITALIANA DI FISICA

# Quantum decay at short, intermediate and long times: Observation in integrated photonics

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# Decay and quantum evolution

**Exponential decay:** classical statistical law on the behavior of the survival probability  $p(t)$  of an unstable state/system

$$\frac{p(t + \Delta t) - p(t)}{p(t)} = -\gamma \Delta t \quad (\Delta t \rightarrow 0) \xrightarrow{\text{red arrow}} \frac{p(t)}{p(t_0)} = e^{-\gamma(t-t_0)}$$

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**Quantum decay:** determined by Hamiltonian evolution

$$|\psi_0\rangle \xrightarrow{\text{blue arrow}} |\psi(t)\rangle = \exp(-iHt) |\psi_0\rangle$$

## SURVIVAL PROBABILITY

$$p(t) = |\langle\psi_0|\psi(t)\rangle|^2 = |\langle\psi_0|e^{-iHt}|\psi_0\rangle|^2 \xrightarrow{t \rightarrow \infty} 0$$

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**Exponential decay:** classical statistical law on the behavior of the survival probability  $p(t)$  of an unstable state/system

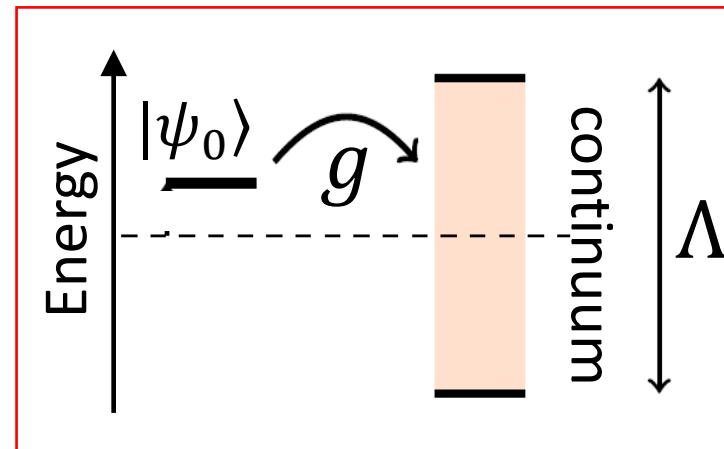
$$\frac{p(t + \Delta t) - p(t)}{p(t)} = -\gamma \Delta t \quad (\Delta t \rightarrow 0) \longrightarrow \frac{p(t)}{p(t_0)} = e^{-\gamma(t-t_0)}$$

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**SURVIVAL PROBABILITY**

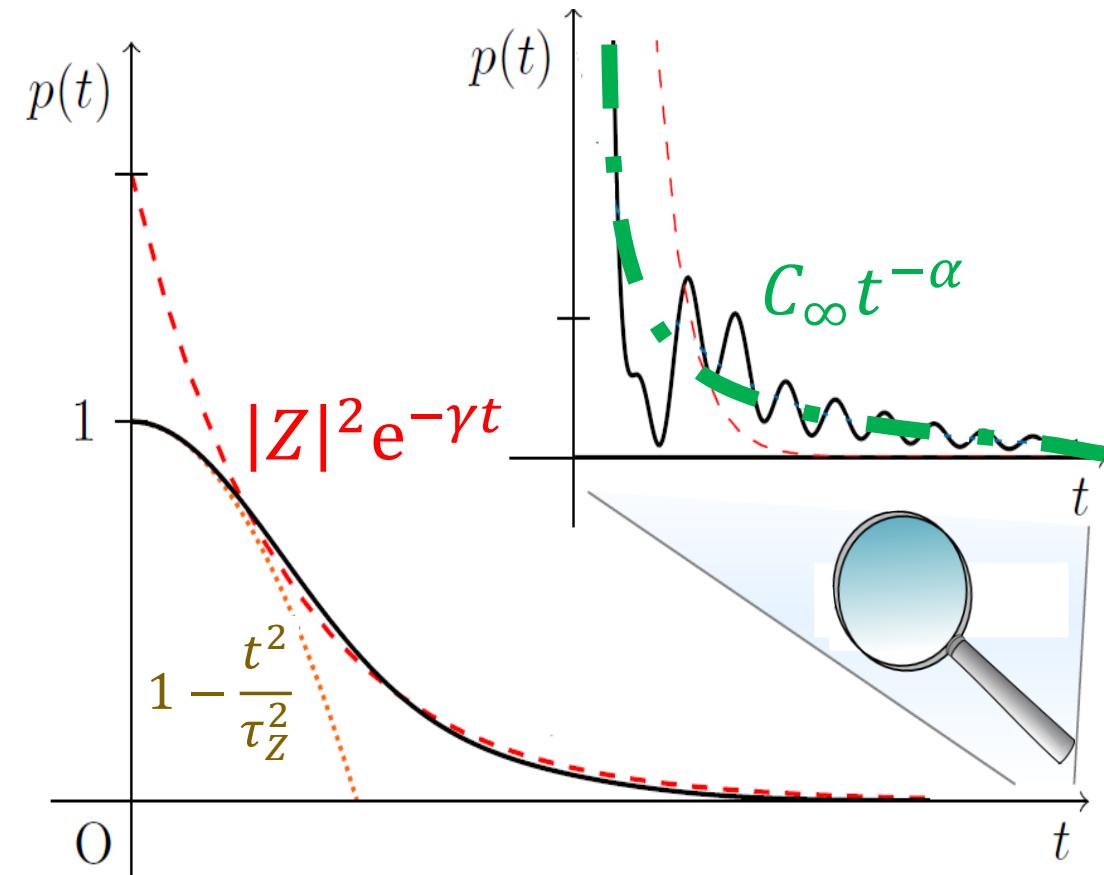
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*is quantum decay always well described by an exponential law?*

# Quantum decay: general results

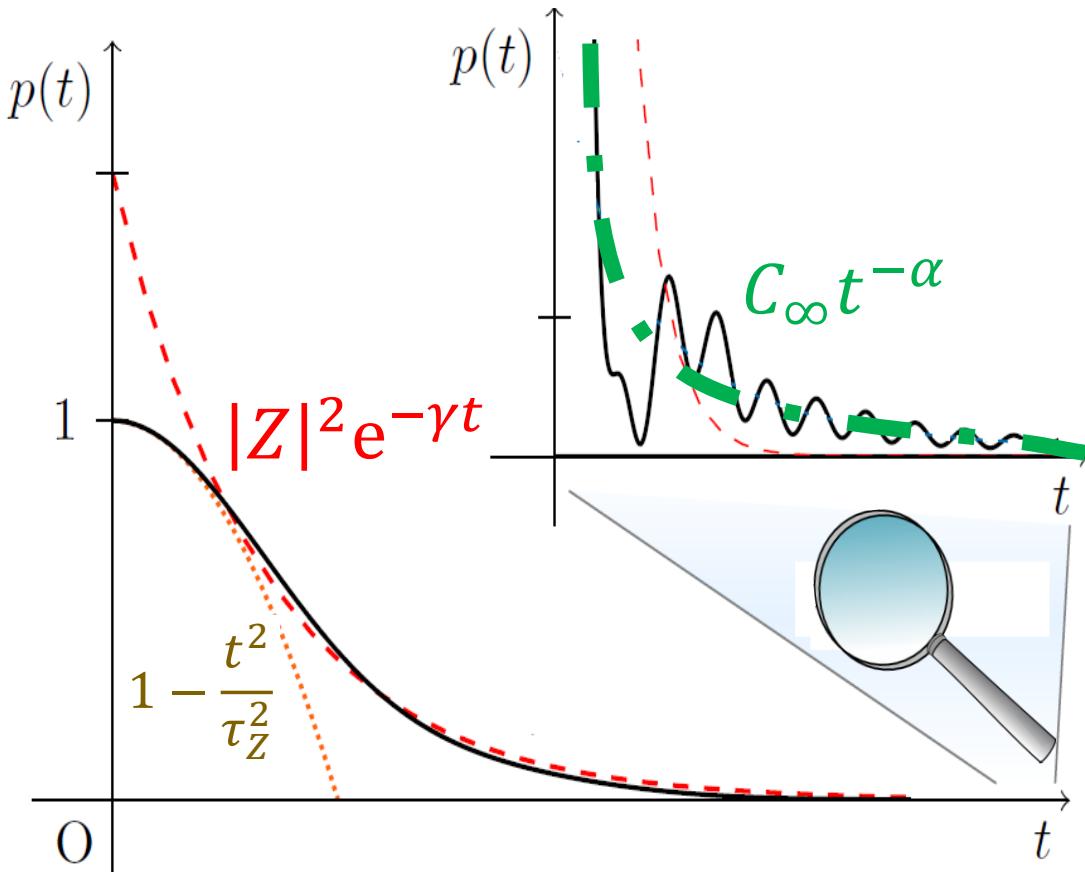
$g$  coupling constant,  $\Lambda$  continuum bandwidth



- **SHORT TIMES ( $t \lesssim \Lambda^{-1}$ ):** quadratic survival probability (Zeno region), observed in different experiments
- **LONG TIMES ( $t \sim g^{-2} \log g$ ):** power-law tail, very elusive in isolated quantum systems (but see Rothe *et al.* PRL 2006)
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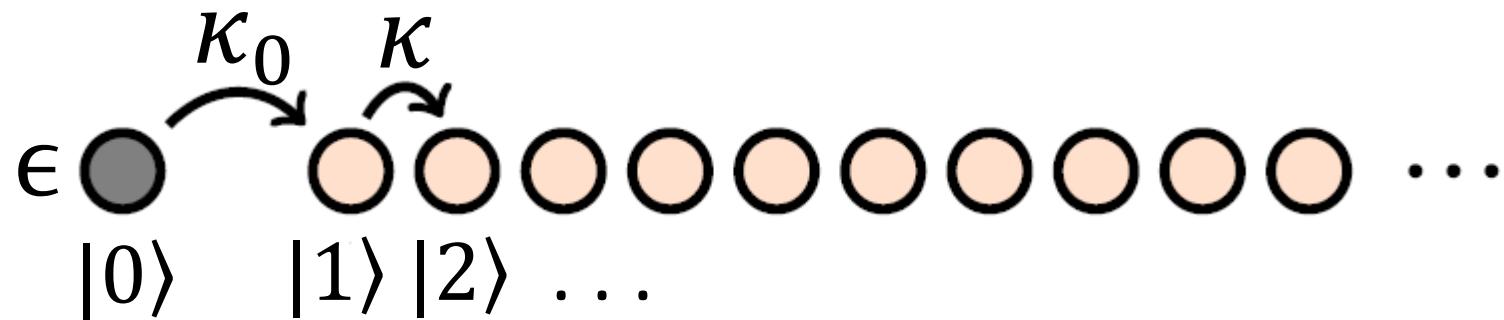
2P  $\rightarrow$  1S electric dipole transition in Hydrogen (natural system, weak coupling)

$$\tau = \gamma^{-1} = 1.6 \times 10^{-9} \text{ s}$$

$$t_{\text{quadratic}} \sim \Lambda^{-1} \sim 10^{-19} \text{ s}$$

$$t_{\text{power}} \sim 100 \tau$$

# One-dimensional hopping model (nearest-neighbor)

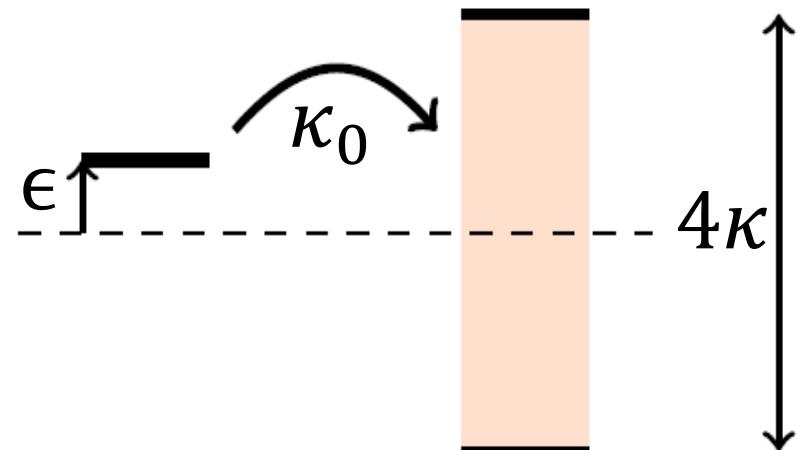


$$H = H_0 + H_1 + H_{\text{int}}$$

$$H_0 = \epsilon |0\rangle\langle 0|$$

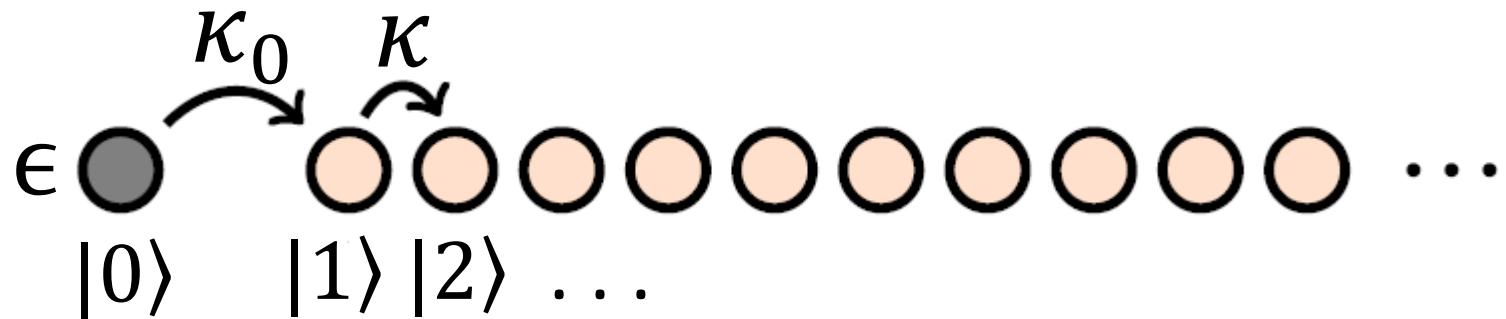
$$H_1 = \kappa \sum_{n=1}^{\infty} |n\rangle\langle n+1| + \text{H. c.}$$

$$H_{\text{int}} = \kappa_0 (|0\rangle\langle 1| + |1\rangle\langle 0|)$$



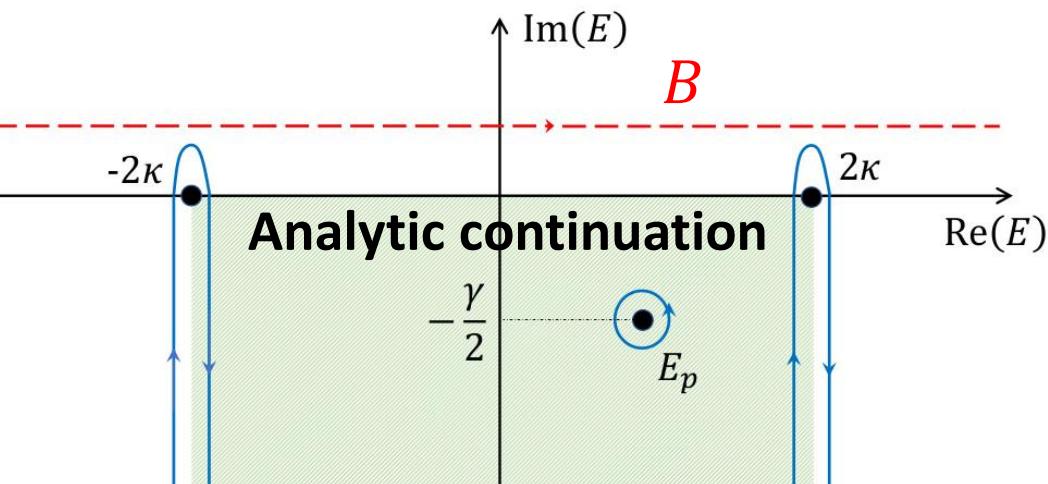
$$\boxed{\chi = \frac{\kappa_0^2}{2\kappa^2}}$$

# One-dimensional hopping model (nearest-neighbor)

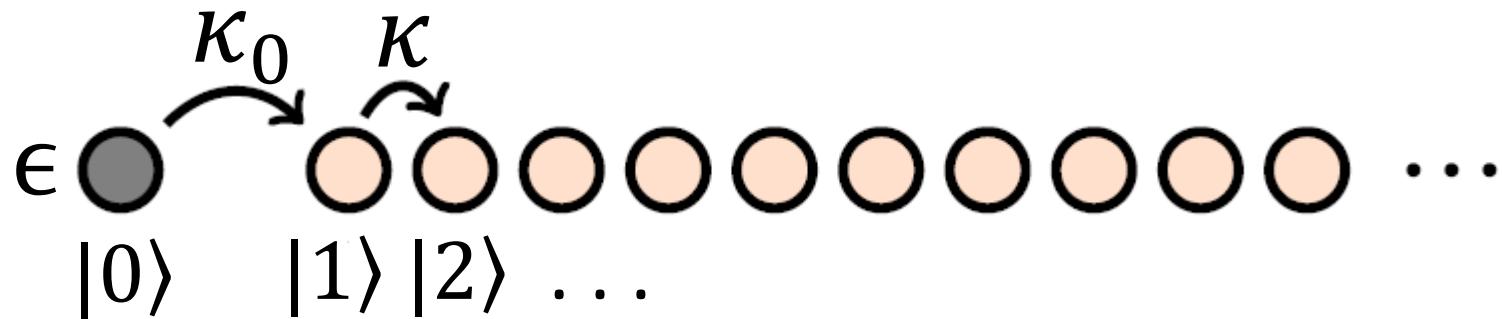


Initial state:  $|\psi_0\rangle = |0\rangle$  **SURVIVAL PROBABILITY**

$$p(t) = \left| \frac{1}{2\pi} \int_B dE e^{-iEt} \langle 0 | (E - H)^{-1} | 0 \rangle \right|^2 = |Z e^{-i\epsilon_p t - \gamma t/2} + \tilde{A}(t)|^2$$

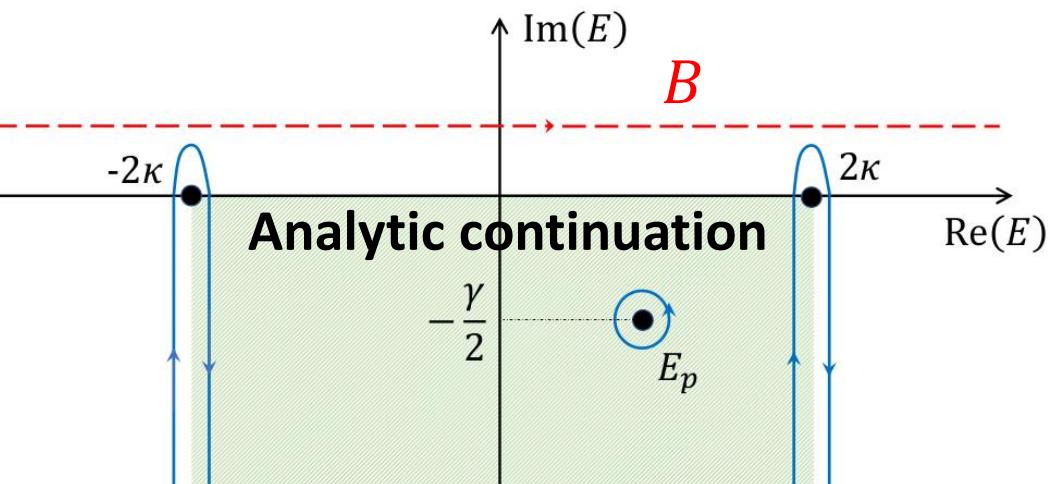


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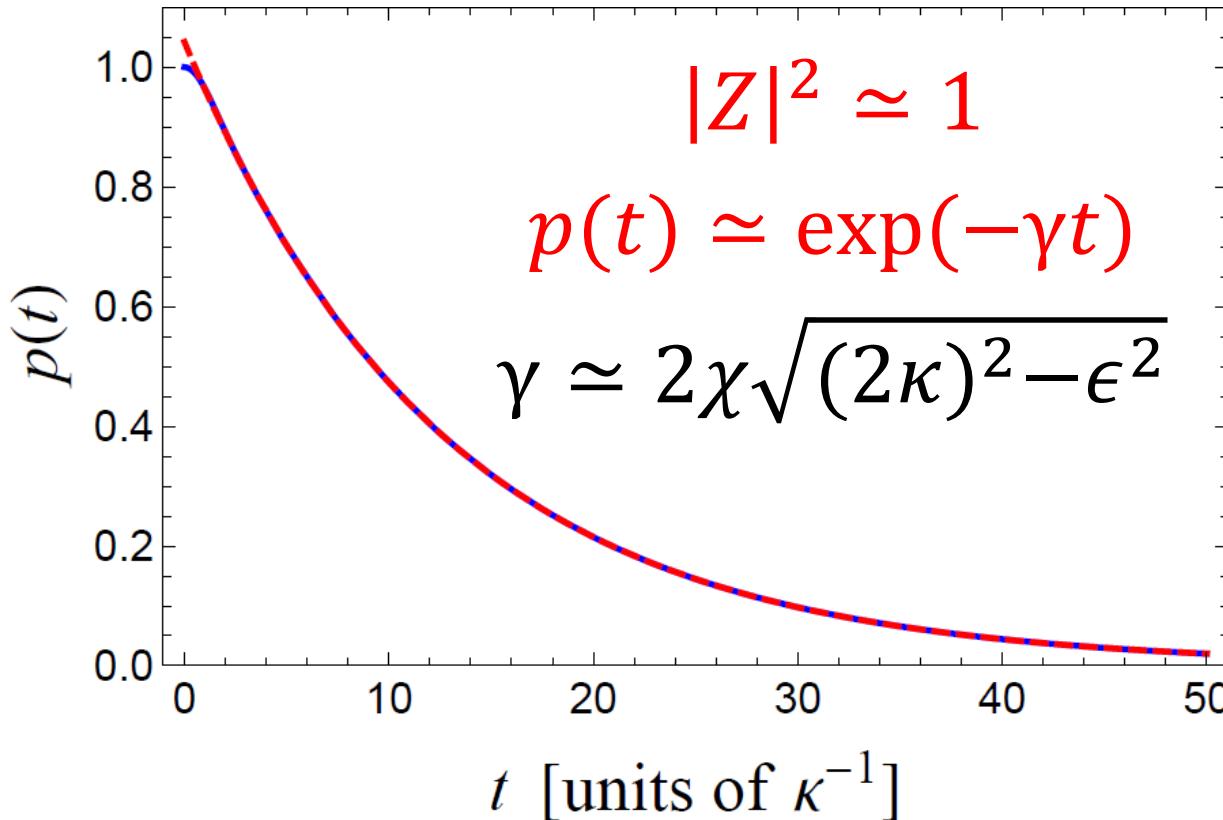


Exponential  
(pole contribution)

Deviations  
(cut contribution)

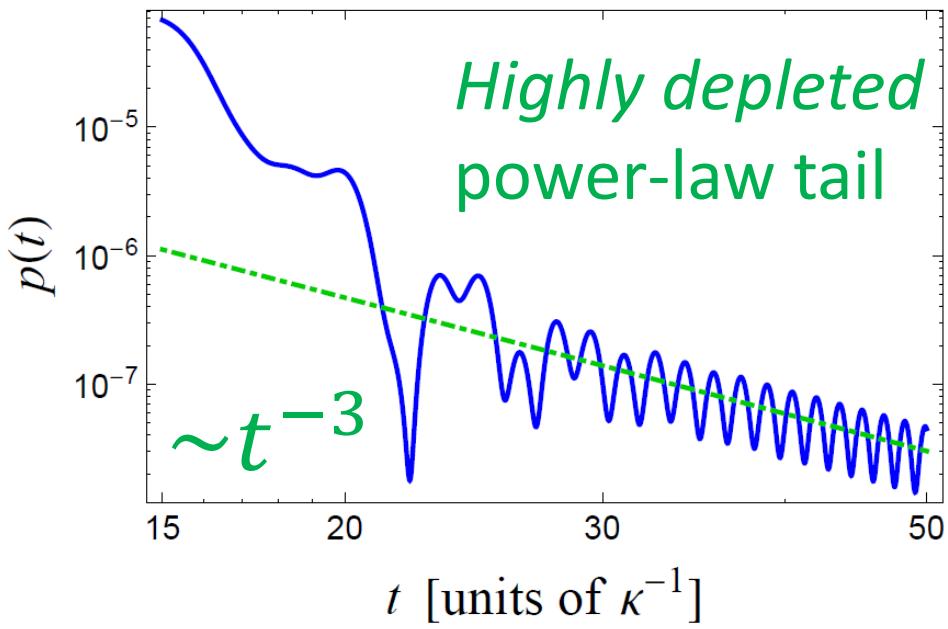
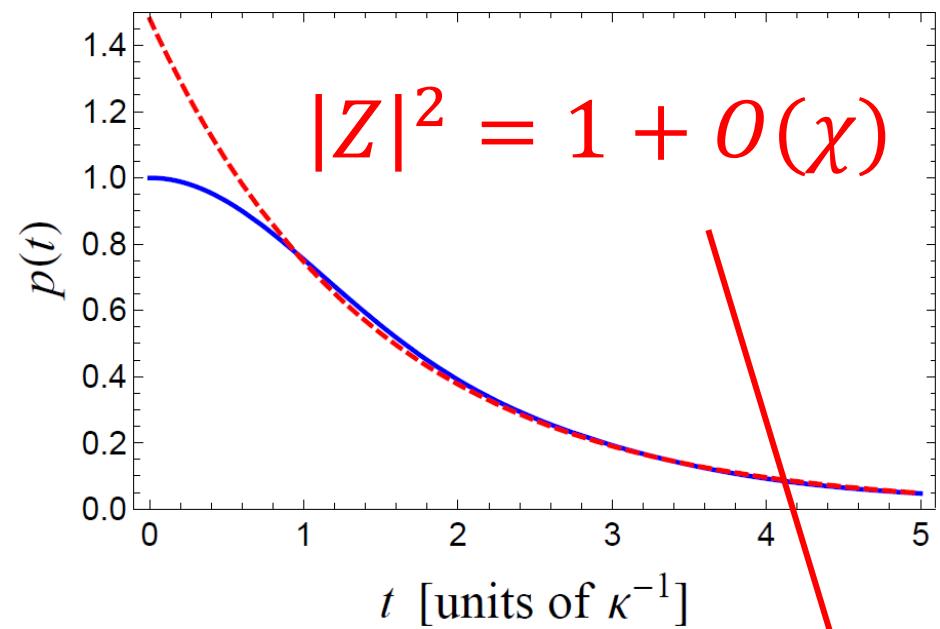
# Decay regimes

$\chi = \frac{\kappa_0^2}{2\kappa^2} \ll 1$  : **exponential decay** with small corrections



# Decay regimes

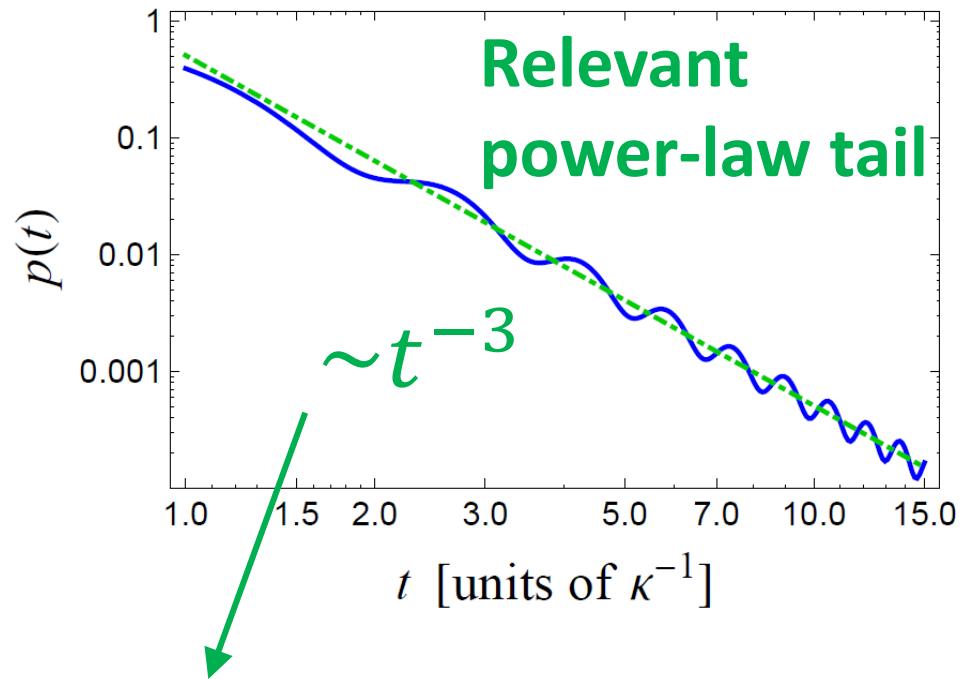
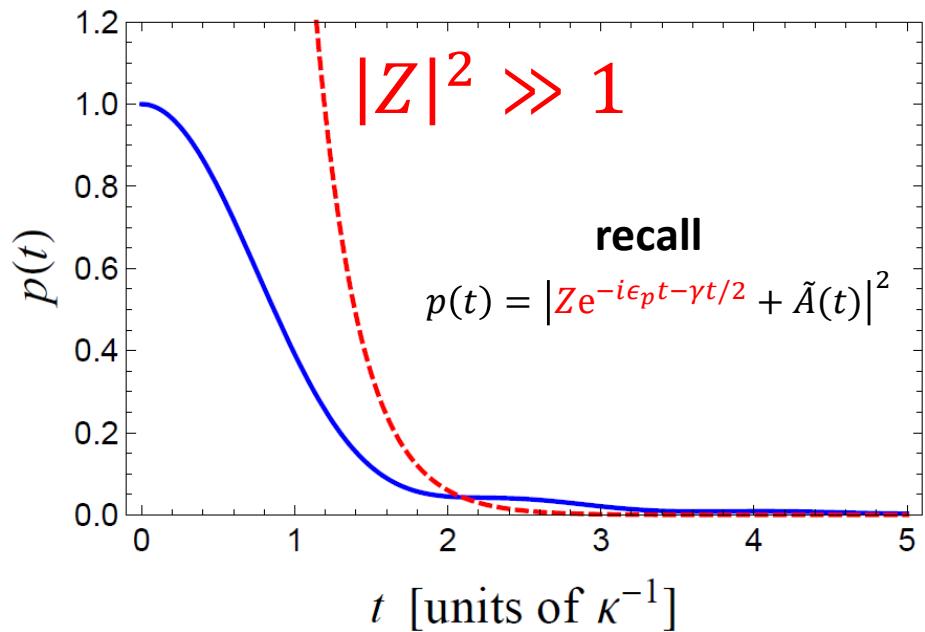
$\chi \lesssim \frac{1 - |\epsilon|/2\kappa}{2 - |\epsilon|/2\kappa}$  : exponential decay with **relevant deviations**



$$|Z|^2 = 1 + \frac{2\chi}{1 - 2\chi} \frac{1 - 3\chi/2 - (\epsilon/2\kappa)^2}{1 - 2\chi - (\epsilon/2\kappa)^2}$$

# Decay regimes

$\frac{1-|\epsilon|/2\kappa}{2-|\epsilon|/2\kappa} \lesssim \chi \leq 1 - |\epsilon|/2\kappa$ : non-exponential decay

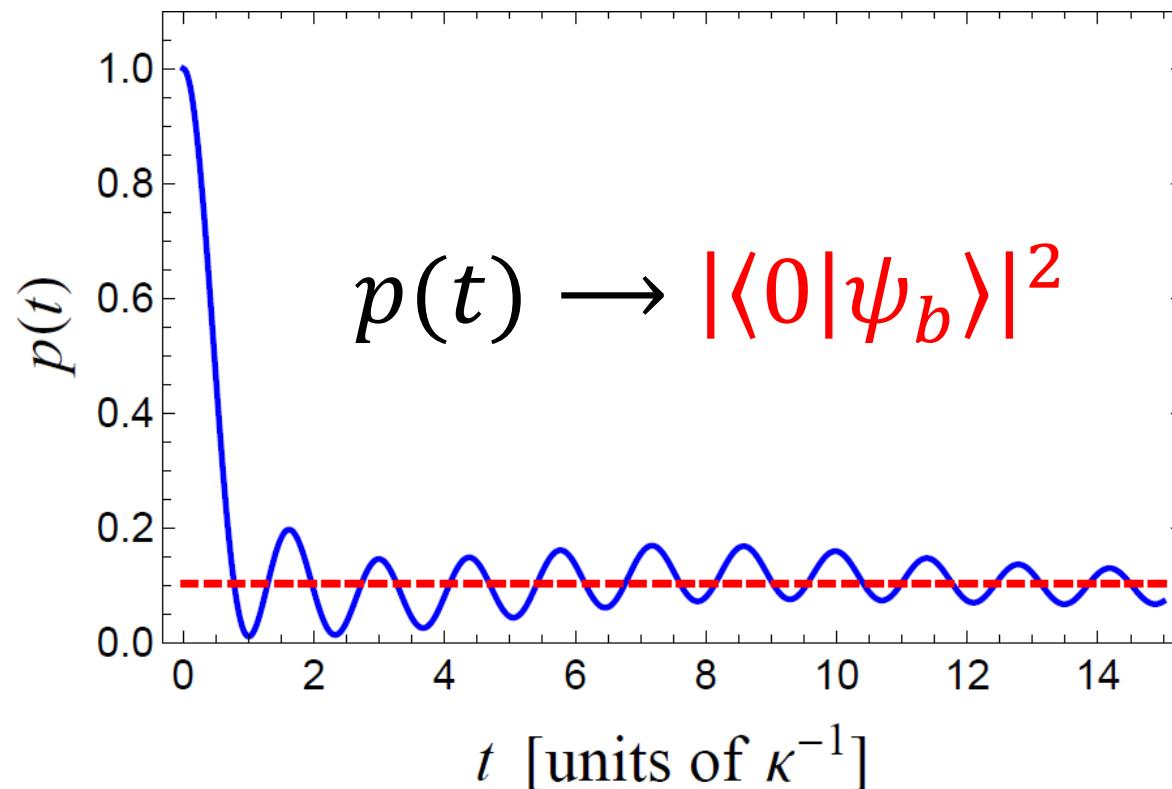


$$p(t) \simeq \frac{\chi^2 \kappa}{\pi t^3} \left| \frac{e^{i(2\kappa t + \pi/4)}}{(2\kappa(1-\chi) + \epsilon)^2} + \frac{e^{-i(2\kappa t + \pi/4)}}{(2\kappa(1-\chi) - \epsilon)^2} \right|^2$$

# Decay regimes

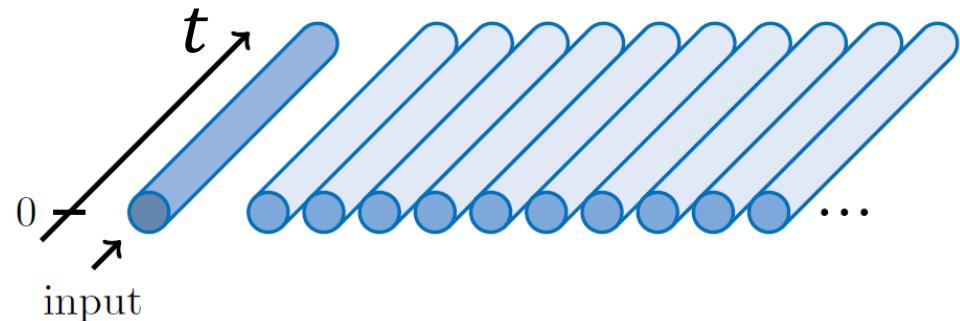
$\chi > 1 - |\epsilon|/2\kappa$ : no instability

Emergence of a **bound state**  $|\psi_b\rangle$  of the total Hamiltonian



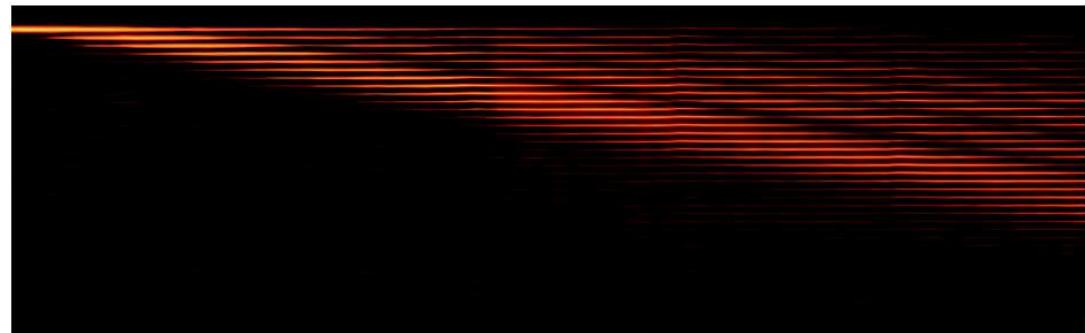
# Experimental implementation of the hopping model

Array of **femtosecond laser-written integrated waveguides**  
(IFN-CNR Milano)



**Quantum-optical analogy** (S. Longhi, 2009): stationary paraxial optics in photonic structures can simulate Schrödinger dynamics.  
**Longitudinal coordinate  $\equiv$  Time**

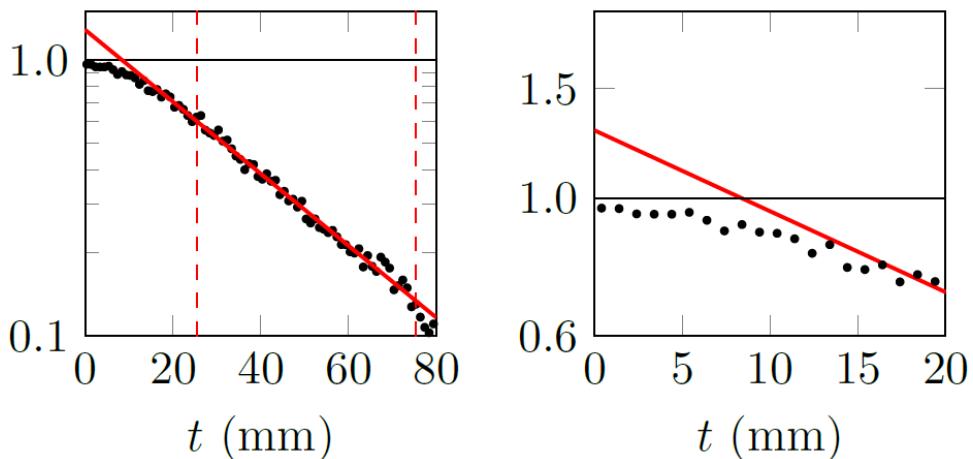
*Scattered laser light  
High dynamic range*



**Next-to-nearest-neighbor** hopping improves prediction accuracy

$$H_2 = q_0(|0\rangle\langle 2| + |2\rangle\langle 0|) + q \sum_{n=1}^{\infty} (|n\rangle\langle n+2| + |n+2\rangle\langle n|)$$

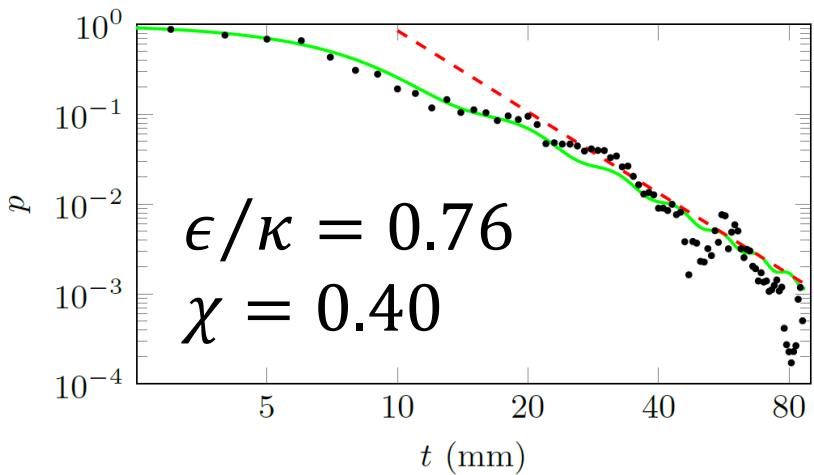
# Observation of different decay regimes



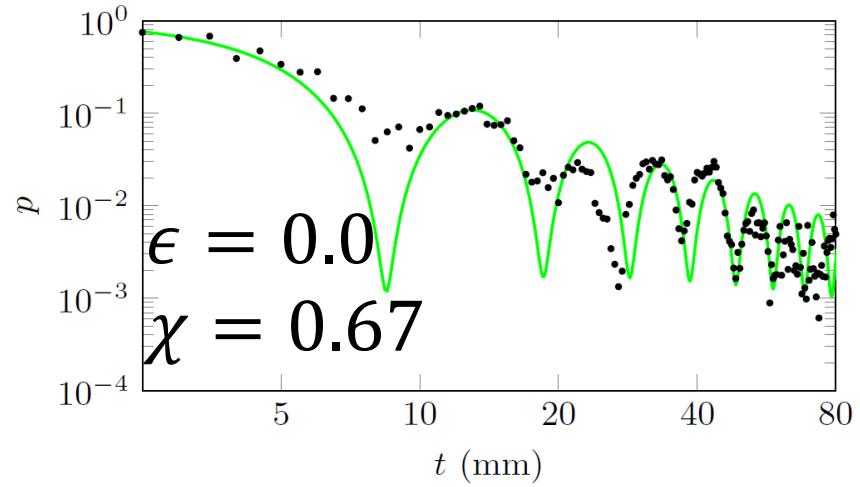
$$\epsilon/\kappa = -0.67$$

$$\chi = 0.071$$

Exponential decay  
with relevant *short-time*  
deviations



Non-exponential decay with  
**detectable power-law** ( $t^{-3}$ ) tail



Non-exponential decay with  
no simple power-law tail

## Conclusions and outlook

- ✓ Models defined on **one-dimensional arrays** represent a controllable testbed to study non-exponential decays and strong-coupling effects in general
- ✓ **Laser-written waveguides** provided a feasible, effective and *fully controllable* experimental platform to obtain evidence of power-law decay

A. Crespi *et al.*, Phys. Rev. Lett. **122**, 130401 (2019)

- Detailed characterization of next-to-nearest neighbor model (article in preparation)
- Effects of disorder on the semi-infinite array
- Observation of the transition *in time* from exponential to power-law decay

**Thank you  
for your attention**