

Quantum computing with semiconductor spins: the flip-flop qubit universal quantum gates set

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SEMICONDUCTOR QUBIT: AN INTRODUCTION

From the original proposals...

PHYSICAL REVIEW A

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JANUARY 1998

Quantum computation with quantum dots

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(Received 9 January 1997; revised manuscript received 22 July 1997)

A silicon-based nuclear spin quantum computer

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Quantum computers promise to exceed the computational efficiency of ordinary classical machines because quantum algorithms allow the execution of certain tasks in fewer steps. But practical implementation of these machines poses a formidable challenge. Here I present a scheme for implementing a quantum-mechanical computer. Information is encoded onto the nuclear spins of donor atoms in doped silicon electronic devices. Logical operations on individual spins are performed using externally applied electric fields, and spin measurements are made using currents of spin-polarized electrons. The realization of such a computer is dependent on future refinements of conventional silicon electronics.

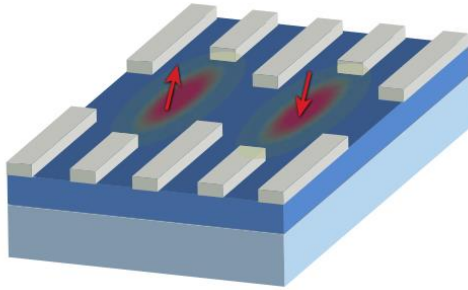
Confinement of spins in semiconductor represents a fruitful platform for UNIVERSAL QUANTUM COMPUTATION

Group III-V
(GaAs)

N-V center
(point defect in
diamond)

Group IV (Si)

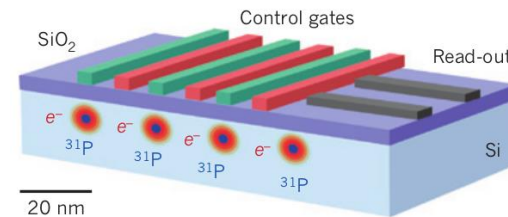
QUANTUM DOT-based qubit



Long coherence time
Fast gate operations
Easy manipulation
Potential for scaling

Group IV semiconductor
Silicon

Nuclear free isotopes
Integrability with CMOS infrastructure



DONOR-based qubit

ARTICLE OPEN

High-fidelity resonant gating of a silicon-based quantum dot hybrid qubit

Dohun Kim^{1,2}, Daniel R Ward¹, Christie B Simmons¹, Don E Savage³, Max G Lagally³, Mark Friesen¹, Susan N Coppersmith¹ and Mark A Eriksson¹

npj | Quantum Information

Review Article | OPEN | Published: 06 September 2017

Interfacing spin qubits in quantum dots and donors—hot, dense, and coherent

L. M. K. Vandersypen[✉], H. Bluhm, J. S. Clarke, A. S. Dzurak, R. Ishihara, A. Morello, D. J. Reilly, L. R. Schreiber & M. Veldhorstnpj Quantum Information 3, Article number: 34 (2017) | Download Citation [↓](#)nature
International journal of science

Letter | Published: 14 February 2018

A programmable two-qubit quantum processor in silicon

T. F. Watson, S. G. J. Philips, E. Kawakami, D. R. Ward, P. Scarlino, M. Veldhorst, D. E. Savage, M. G. Lagally, Mark Friesen, S. N. Coppersmith, M. A. Eriksson & L. M. K. Vandersypen[✉]Nature 555, 633–637 (29 March 2018) | Download Citation [↓](#)

High-fidelity readout and control of a nuclear spin qubit in silicon

Jarryd J. Pla¹, Kuan Y. Tan^{1,†}, Juan P. Dehollain¹, Wee H. Lim^{1,†}, John J. L. Morton², Floris A. Zwanenburg^{1,†}, David N. Jamieson³, Andrew S. Dzurak¹ & Andrea Morello¹

npj | Quantum Information

Review Article | OPEN | Published: 26 June 2017

Quantum information density scaling and qubit operation time constraints of CMOS silicon-based quantum computer architectures

Davide Rotta, Fabio Sebastiano, Edoardo Charbon & Enrico Prati[✉]npj Quantum Information 3, Article number: 26 (2017) | Download Citation [↓](#)nature
COMMUNICATIONS

ARTICLE

DOI: 10.1038/s41467-017-01113-2 OPEN

Coherent coupling between a quantum dot and a donor in silicon

Patrick Harvey-Collard^{1,2}, N. Tobias Jacobson³, Martin Rudolph², Jason Dominguez², Gregory A. Ten Eyck², Joel R. Wendt², Tammy Pluyms², John King Gamble³, Michael P. Lilly⁴, Michel Pioro-Ladrière^{1,5} & Malcolm S. Carroll²nature
COMMUNICATIONS

ARTICLE

DOI: 10.1038/s41467-017-01905-6 OPEN

Silicon CMOS architecture for a spin-based quantum computer

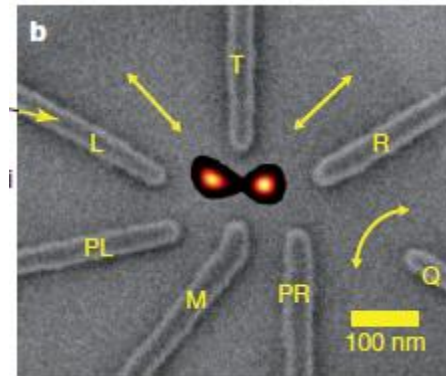
M. Veldhorst^{1,2}, H.G.J. Eenink^{1,2}, C.H. Yang² & A.S. Dzurak²

SOME RECENT RESULTS...

Intel reveals silicon-based quantum processor prototype

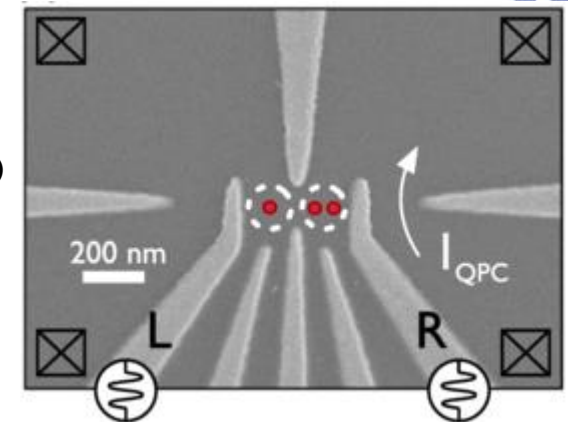
The 'spin qubit processor' is made with isotopically pure silicon wafers, and uses magnetic resonance to manipulate individual electrons.

DOUBLE QD SINGLET-TRIPLET QUBIT



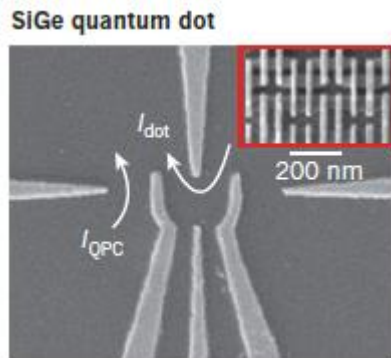
B. M. Maune et al., Nature 481, 344 (2012)

DOUBLE QD HYBRID QUBIT



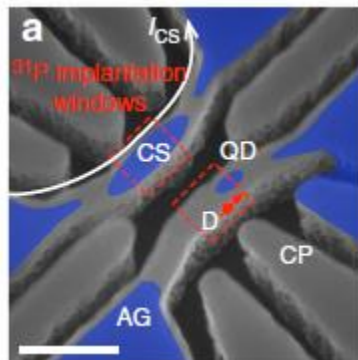
J. C. Abadillo-Uriel et al., Phys. Rev. B 98, 165438 (2018)

QD SINGLE SPIN QUBIT



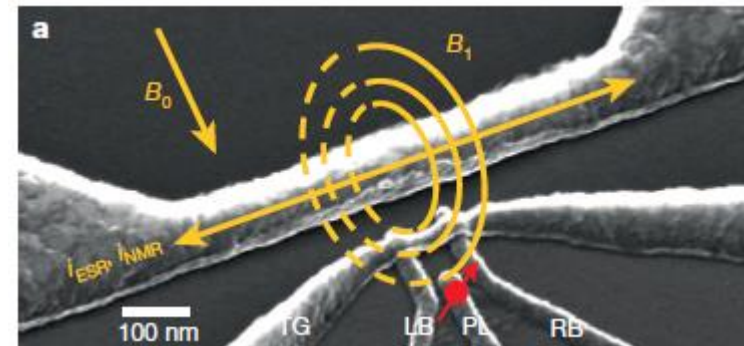
J. J. L. Morton et al., Nature 479, 345–353 (2011)

QD SPIN-DONOR QUBIT



P. Harvey-Collard, Nature Communications 8, 1029 (2017)

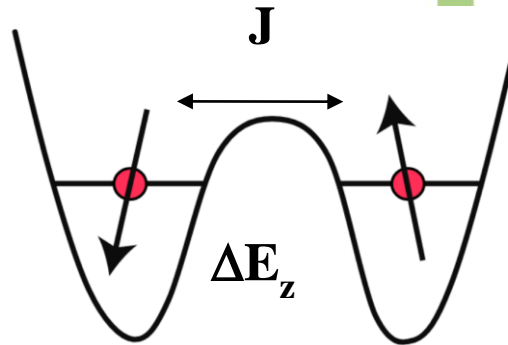
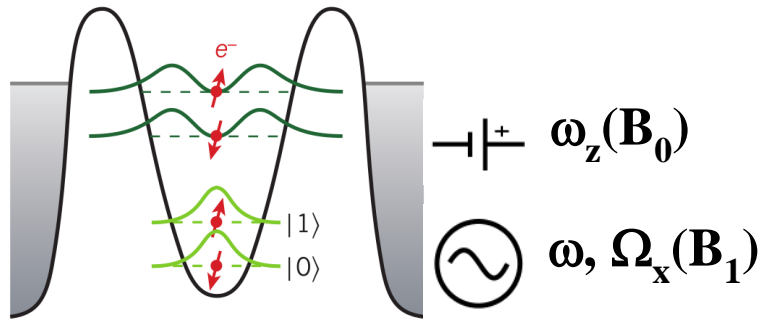
DONOR QUBIT



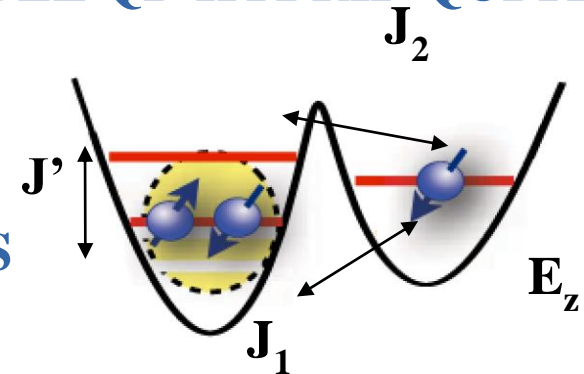
J. J. Pla et al., Nature 496, 334 (2013)

DOUBLE QD SINGLET-TRIPLET QUBIT

QD SINGLE SPIN QUBIT

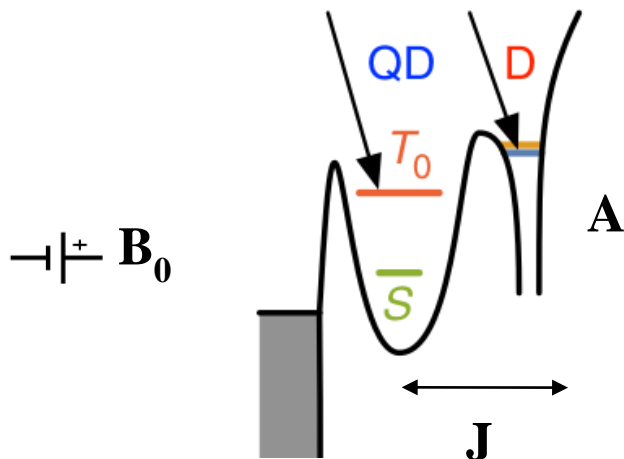


DOUBLE QD HYBRID QUBIT

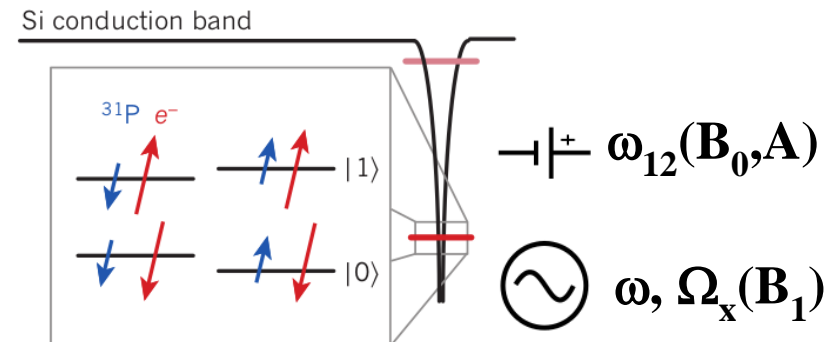


SPIN QUBITS: ENERGY PROFILES & PHYSICAL CONTROLS

QD SPIN-DONOR QUBIT



DONOR QUBIT



Qubit	Material	f(MHz)	$T_2^*(ns)$	$T_2(ns)$	$Q \equiv T_2^*/T_\pi$	Ref.
Single spin	Si/SiGe	~ 5	$\sim 9 \times 10^2$	3.7×10^4	~ 9	E. Kawakami et al., Nat. Nanotech. 9, 666–670 (2014)
Single spin	^{28}Si	~ 0.3	$\leq 1.2 \times 10^5$	1.2×10^6	≤ 80	M. Veldhorst et al., Nature 526, 410–414 (2015)
Donor spin (e^-)	P in ^{nat}Si	~ 3	55	2×10^5	≤ 1	J. J. Pla et al., Nature 489, 541 (2012)
Donor spin (e^-)	P in ^{28}Si	~ 0.2	$\sim 3 \times 10^5$	1×10^6	~ 108	J. T. Muhonen et al., Nat. Nanotech. 9, 986 (2014)
Singlet-Triplet	Si/SiGe	~ 351	$> 1 \times 10^3$	n.a.	n.a.	K. Takeda et al., arXiv:1910.00771 (2019)
Hybrid	Si/SiGe	$\sim 1 \times 10^4$	~ 11	~ 40	~ 250	D. Kim et al., Nature 511, 70–74 (2014)

Table I. Comparison of different spin-based qubits.

	Semiconductor Single-Spin qubit	Semiconductor Hybrid qubit (Steane code)	Semiconductor Hybrid qubit (Surface code)	Superconductor Flux qubit (DWave like)	Superconductor Transmon qubit (IBM like)	Trapped Ion qubit
Mqb_{ph}/cm^2	8000	830	100×10^2	8×10^{-4}	10^{-5}	2×10^{-5}
$A_{chip}(mm^2)$	25	240	20	25×10^7	2×10^{10}	10^{10}
Ref.	R. Li et al., Science Adv. 4, eaar3960 (2018)	D. Rotta et al., Npj Quant. Inf. 3, 26 (2017)	D. Rotta et al., Npj Quantum Inf. 3, 26 (2017)	R. Harris and et al., Phys. Rev. B 82, 024511 (2010)	J. M. Gambetta et al., Npj Quant. Inf. 655, 3:2 (2017).	B. Lekitsch et al., Science Adv. 3, e1601540 (2017)

Table II. Number of physical qubits per unit surface and area covered by 2 billions of physical qubits. The silicon hybrid qubit footprint refers to the 7 nm technology node.

E. Ferraro and E. Prati, *Is all-electrical silicon quantum computing feasible in the long term?*, Physics Letters A 384 (17), 126352 (2020)



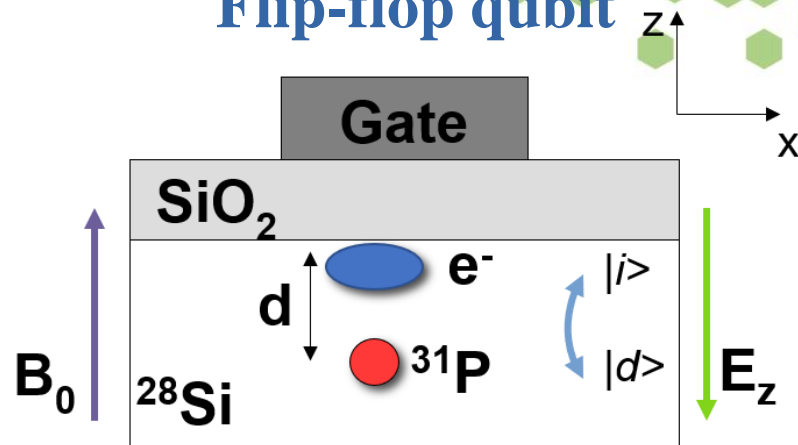
FLIP-FLOP QUBIT

Flip-flop qubit

Encoded states:

$$|0\rangle = |g \downarrow \uparrow\rangle$$

$$|1\rangle = |g \uparrow \downarrow\rangle$$



$$\hat{H}_{B_0} = \gamma_e B_0 \left[\hat{\mathbb{1}} + \left(\frac{\hat{\mathbb{1}} + \hat{\sigma}_z^{id}}{2} \right) \Delta\gamma \right] \hat{S}_z - \gamma_n B_0 \hat{I}_z$$

$$\hat{H}_A = A \left(\frac{\hat{\mathbb{1}} - \hat{\sigma}_z^{id}}{2} \right) \mathbf{S} \cdot \mathbf{I}$$

$$\hat{H}_{Orb} = \frac{V_t}{2} \hat{\sigma}_x^{id} - \frac{d e (\Delta E_z + E_a \cos(\omega_E t))}{2 h} \hat{\sigma}_z^{id}$$

Two eigenstates:

$|g\rangle$: e^- ground state

$|e\rangle$: e^- excited state

Applying a large static magnetic field B_0 , (i.e. $(\gamma_e + \gamma_n)B_0 \gg A$), the eigenstates of the system are the four qubit states $\{|\downarrow\uparrow\rangle, |\downarrow\downarrow\rangle, |\uparrow\downarrow\rangle, |\uparrow\uparrow\rangle\}$. Electrically modulating the hyperfine interaction A by E_z causes an Electron Dipole Spin Resonance (EDSR) transition between the states with antiparallel spins $\{|\downarrow\uparrow\rangle, |\uparrow\downarrow\rangle\}$, that are in turn chosen to encode the qubit.

Universal quantum gate set

$G = \{H, \Lambda(S)\}$, where $\Lambda(S)$ is a two-qubit gate in which the operation S is applied to the target qubit if and only if the control qubit is in the logical state $|1\rangle$, for example the CNOT gate.

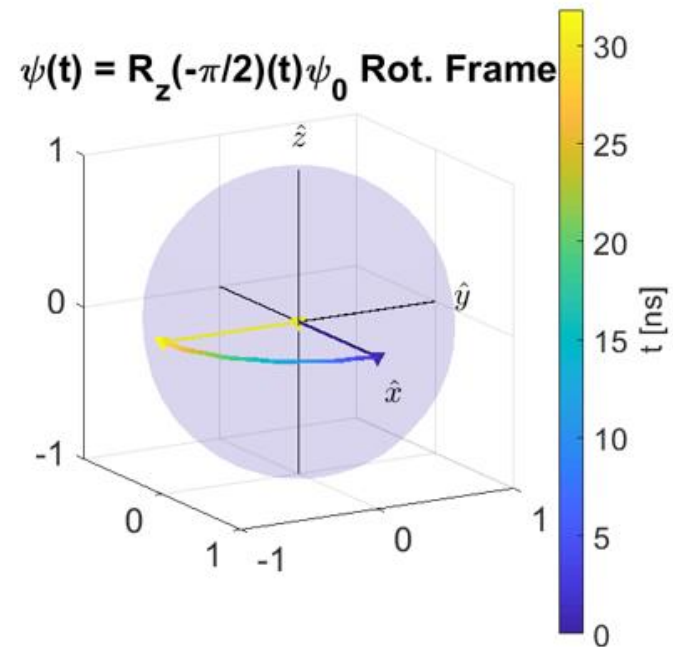
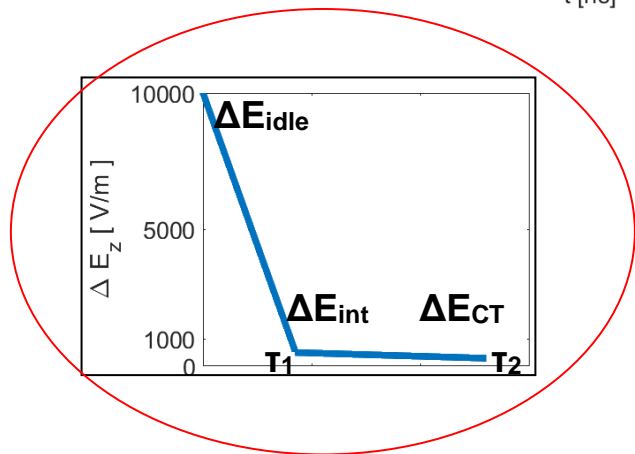
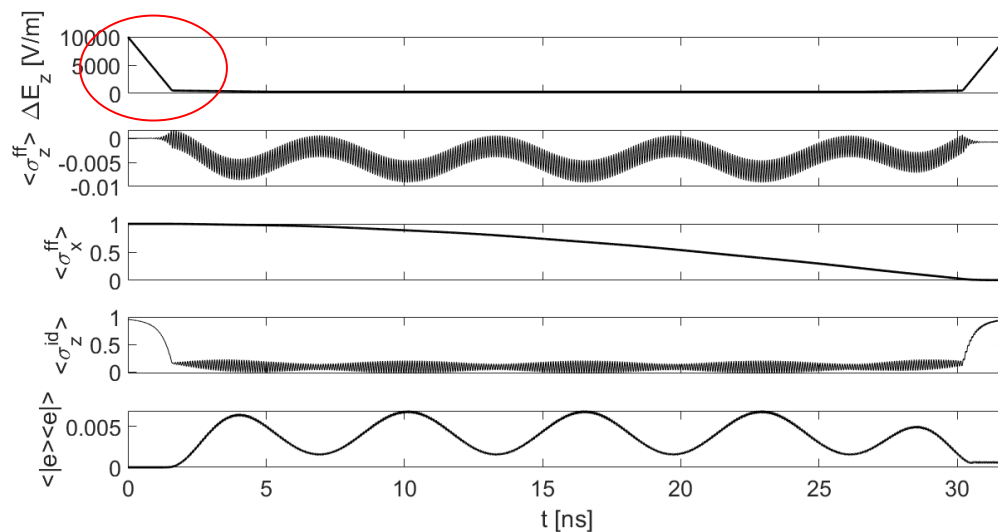
Moreover, a construction of the CNOT gate using only $R_z(-\pi/2)$, H and $\text{sqrt}(i\text{SWAP})$ gates is feasible.

$$G = \{R_z(-\pi/2), H, \text{sqrt}(i\text{SWAP})\}$$

For each gate operation, we consider the effect of the charge noise using the $1/f$ model for the power spectral density, dominant in the low-frequency regime.

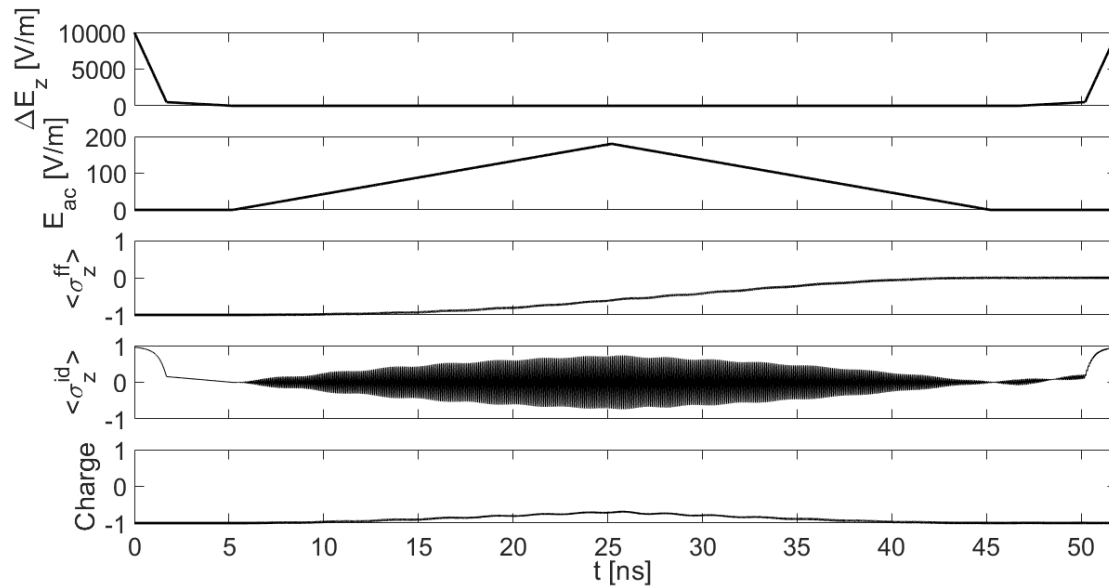
$$R_z(-\pi/2)$$

ΔE_{idle}	ΔE_{int}	ΔE_{ct}	τ_1	τ_2	T	V_t	K	T_{gate}
[V/m]	[V/m]	[V/m]	[ns]	[ns]	[ns]	[GHz]		[ns]
10000	500	290	1.7	3.5	21.6	11.29	$\simeq 20$	31.9

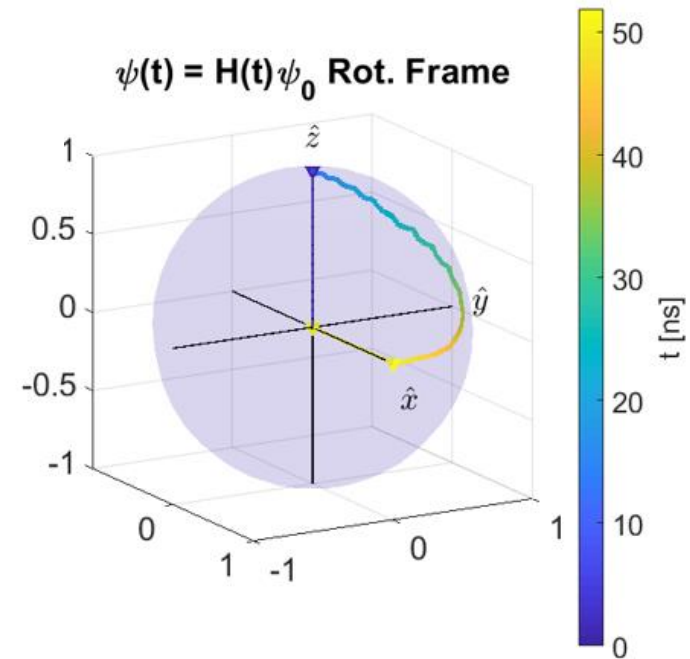


Hadamard

ΔE_{idle}	ΔE_{int}	ΔE_{ct}	E_{ac}	τ_1	τ_2	T	$T_{E_{ac}}^{ON}$	V_t	K	K_E	T_{gate}
[V/m]	[V/m]	[V/m]	[V/m]	[ns]	[ns]	[ns]	[ns]	[GHz]			[ns]
10000	500	0	180	1.7	3.5	41.5	40	11.5	$\simeq 21$	$\simeq 57$	51.9



Acts on the qubit as a rotation of an angle π around the $(\hat{x} + \hat{z})/\sqrt{2}$ axis



1/f noise model

- The 1/f noise model is based on the definition of the Power Spectral Density that is inversely proportional to the frequency and is given by $S(\omega) = \alpha / (\omega t_0)$, where α is the noise amplitude, that does not depend on ω and t_0 is the time unit.
- We generated the 1/f noise in the frequency domain as

$$n(\omega) = m(\omega)^{-1/2} e^{i\varphi(\omega)},$$

where $m(\omega)$ is generated from a standard Gaussian white process and the phase factor $\varphi(\omega) = [0, 2\pi]$ is chosen uniformly.

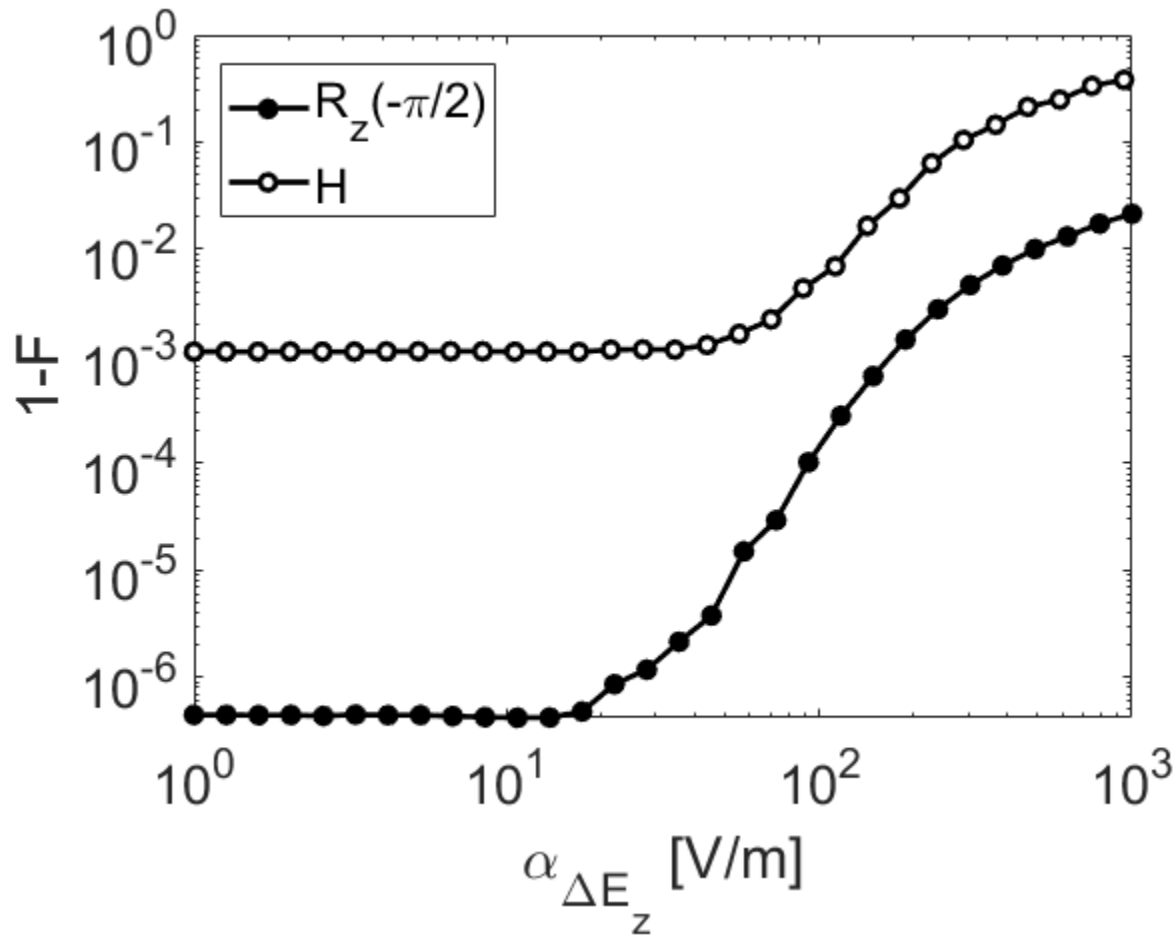
- To obtain the noise in the time domain, we calculate the inverse Fourier transform and then multiply the result by the noise amplitude α

ENTANGLEMENT FIDELITY

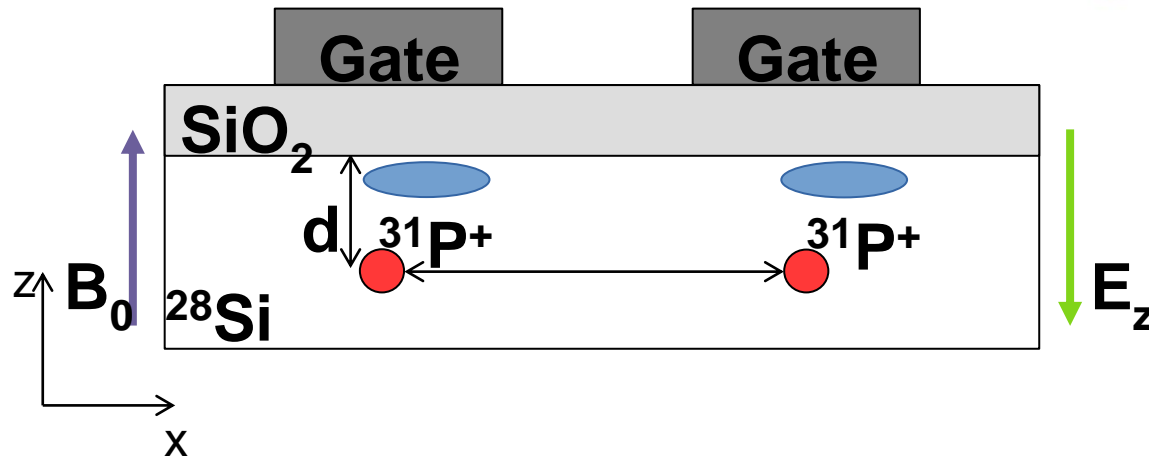
$$F = \text{tr}[\rho^{RS} \mathbf{1}_R \otimes (U_i^{-1} U_d)_S \rho^{RS} \mathbf{1}_R \otimes (U_d^{-1} U_i)_S]$$

$$\rho^{RS} = |\psi\rangle\langle\psi| \text{ with } |\psi\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$

Entanglement Fidelity in presence of 1/f noise model



Two flip-flop qubits: the dipole-dipole interaction



The two-qubit Hamiltonian

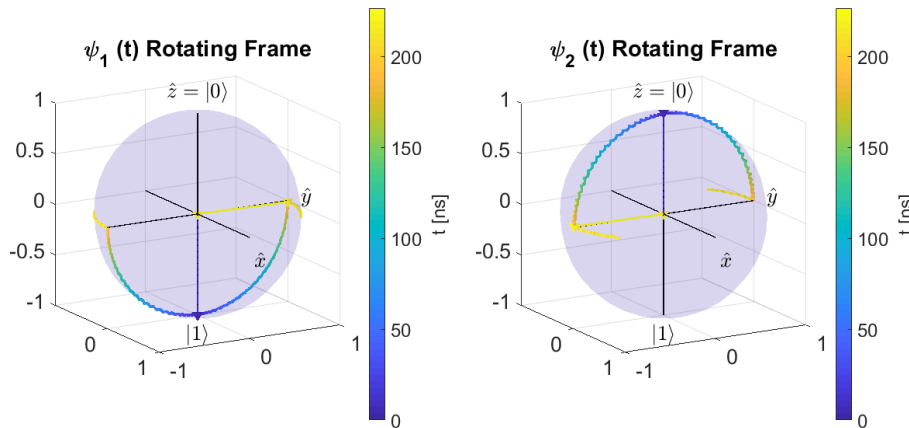
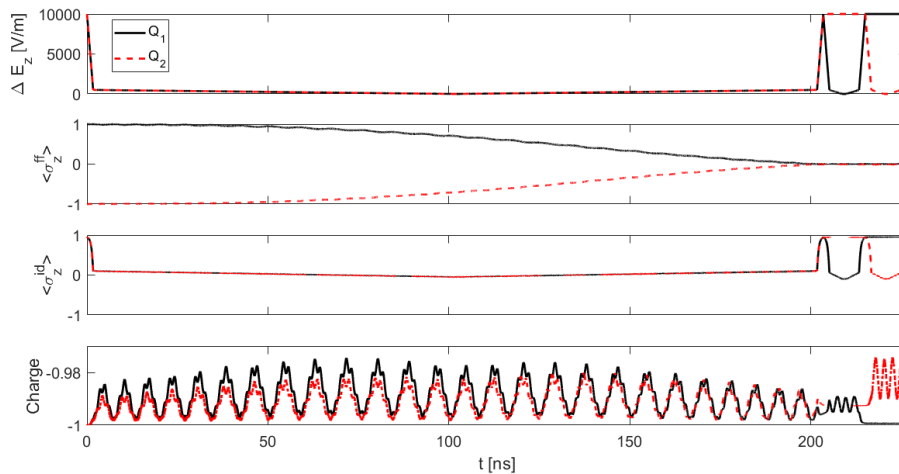
$$\hat{H}_{2,ff} = \hat{H}_{ff,1} \otimes \hat{1} + \hat{1} \otimes \hat{H}_{ff,2} + \hat{V}_{dip}$$

$$\hat{V}_{dip} = \frac{e^2 d^2}{4\pi\epsilon_0\epsilon_r r^3} |i_1 i_2\rangle \langle i_1 i_2|$$

$|i_1\rangle, |i_2\rangle$: e^- state at the interface

$\sqrt{i\text{SWAP}}$ @ $r = 180 \text{ nm}$

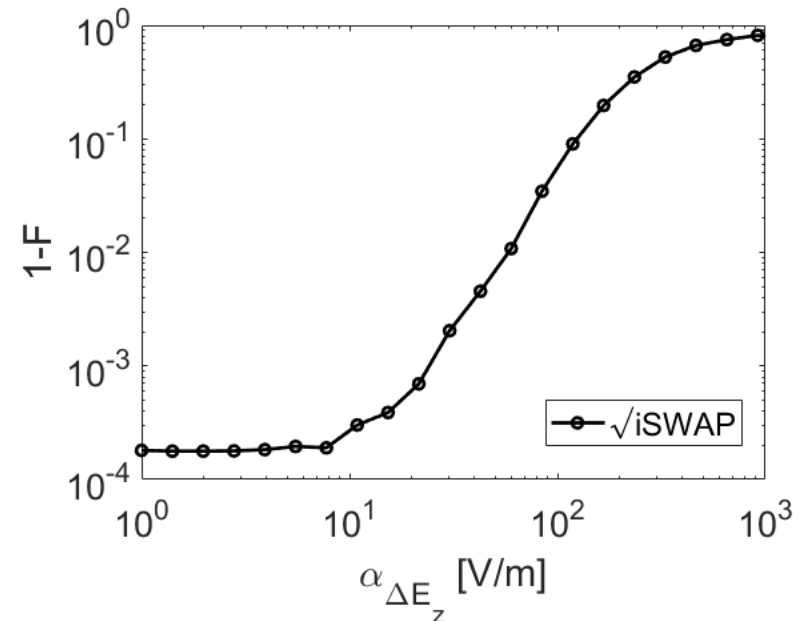
	ΔE_{idle}	ΔE_{int}	ΔE_{ct}	τ_1	τ_2	T	V_t	K	T_{gate}
	[V/m]	[V/m]	[V/m]	[ns]	[ns]	[ns]	[GHz]		[ns]
$Q_1 Q_2$	10000	500	0	1.7	99	2	11.58	$\simeq 21$	203.4
$Q_1(Q_2)$	10000	500	0	1.7	3.5	1.2	11.58	$\simeq 21$	11.6



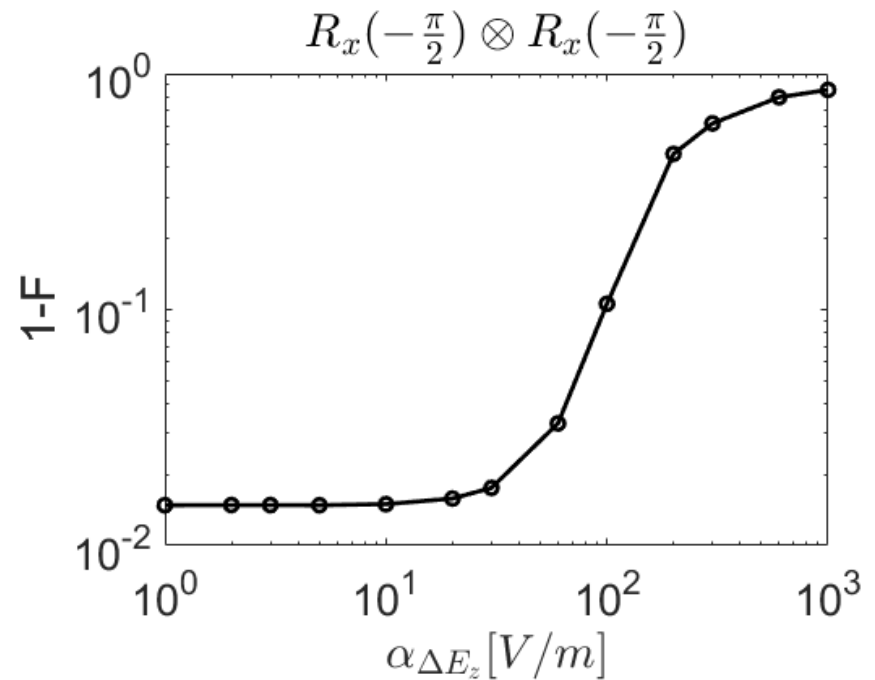
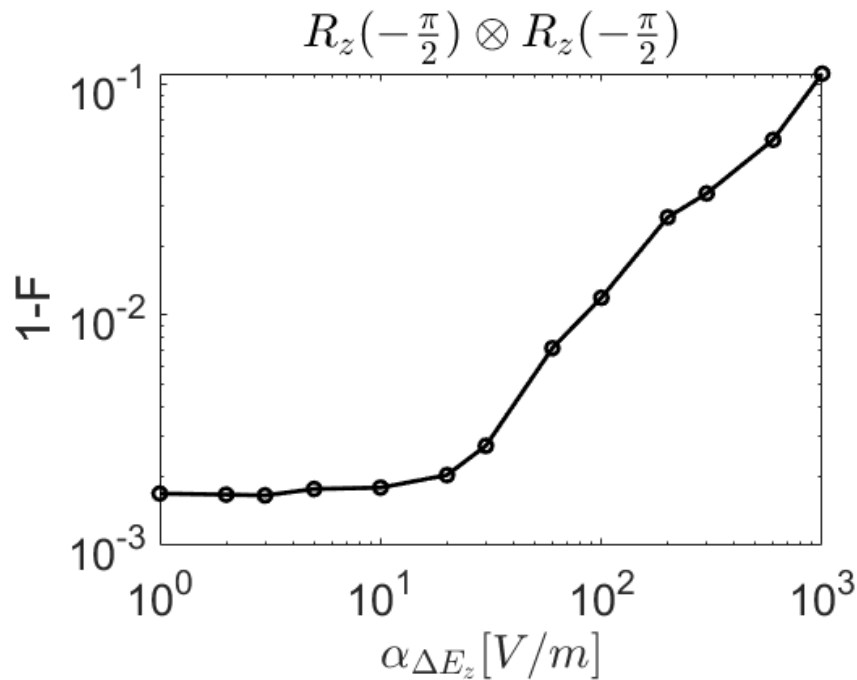
$$\sqrt{i\text{SWAP}} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \frac{1}{\sqrt{2}} & \frac{i}{\sqrt{2}} & 0 \\ 0 & \frac{i}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$|\Psi_1(0)\rangle = |1\rangle$$

$$|\Psi_2(0)\rangle = |0\rangle$$



Transversal gates for Quantum Error Correction: some preliminary results @ r=360 nm



The infidelity is compared with its ideal value where the qubits does not interact

Conclusions

Quantum computation by **flip-flop qubits**, donor based qubits in which the logical states are encoded in the donor nuclear and its bound electron.

A **universal set of quantum gates** composed by $\{R_z(-\pi/2), H, \text{sqrt}(i\text{SWAP})\}$ has been presented and the noise effect on the entanglement fidelity has been studied.

The noise model adopted shows a **1/f spectrum**, typical of qubit sensitive to charge noise.

Results are very promising in correspondance to a realistic noise level around 50 V/m, we obtain $F \geq 99.999\%$ for the $R_z(-\pi/2)$ and 99.8% for the H gate. The two-qubit gate $\text{sqrt}(i\text{SWAP})$ may be realized with a fidelity above 99.5%.

Study on parallel one and two qubit operations for **QEC** code.