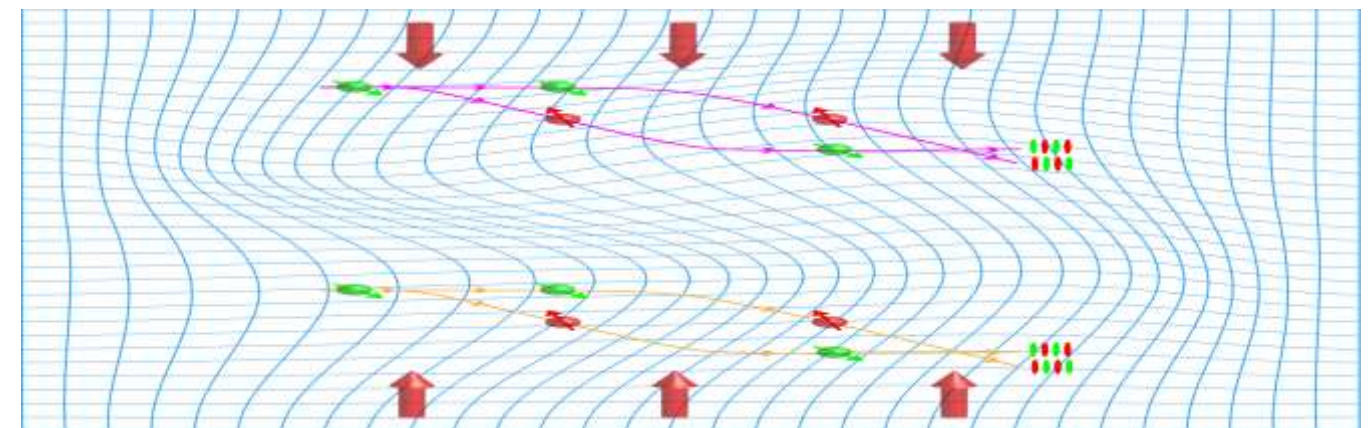
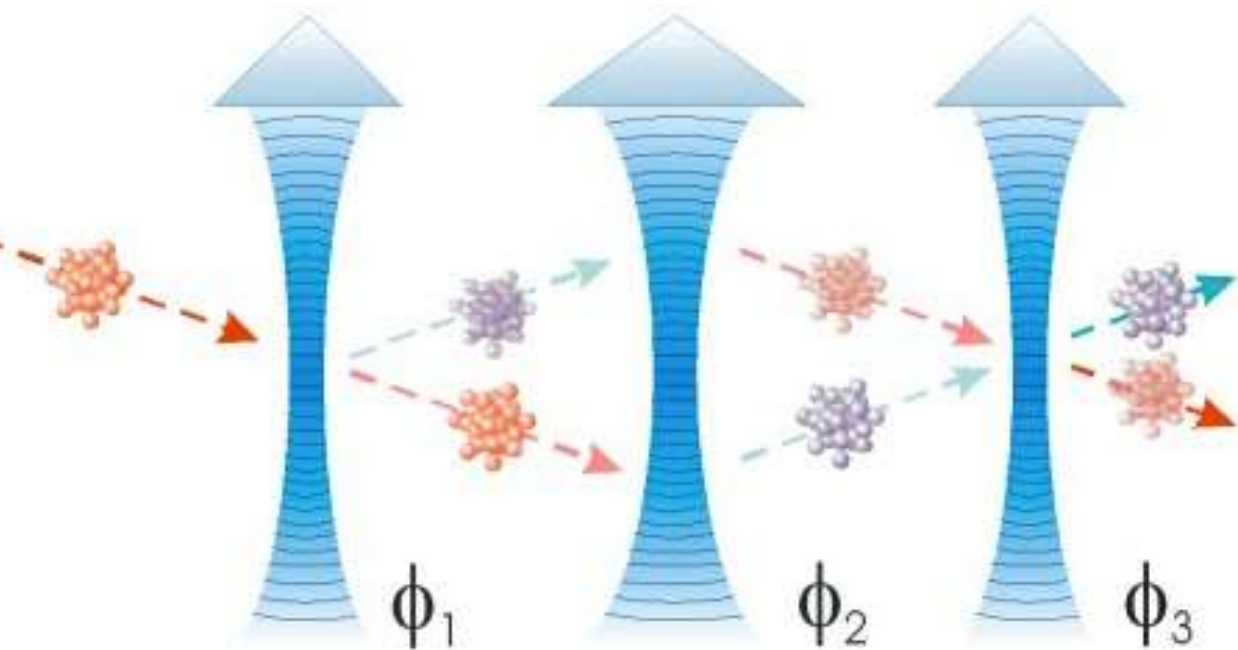
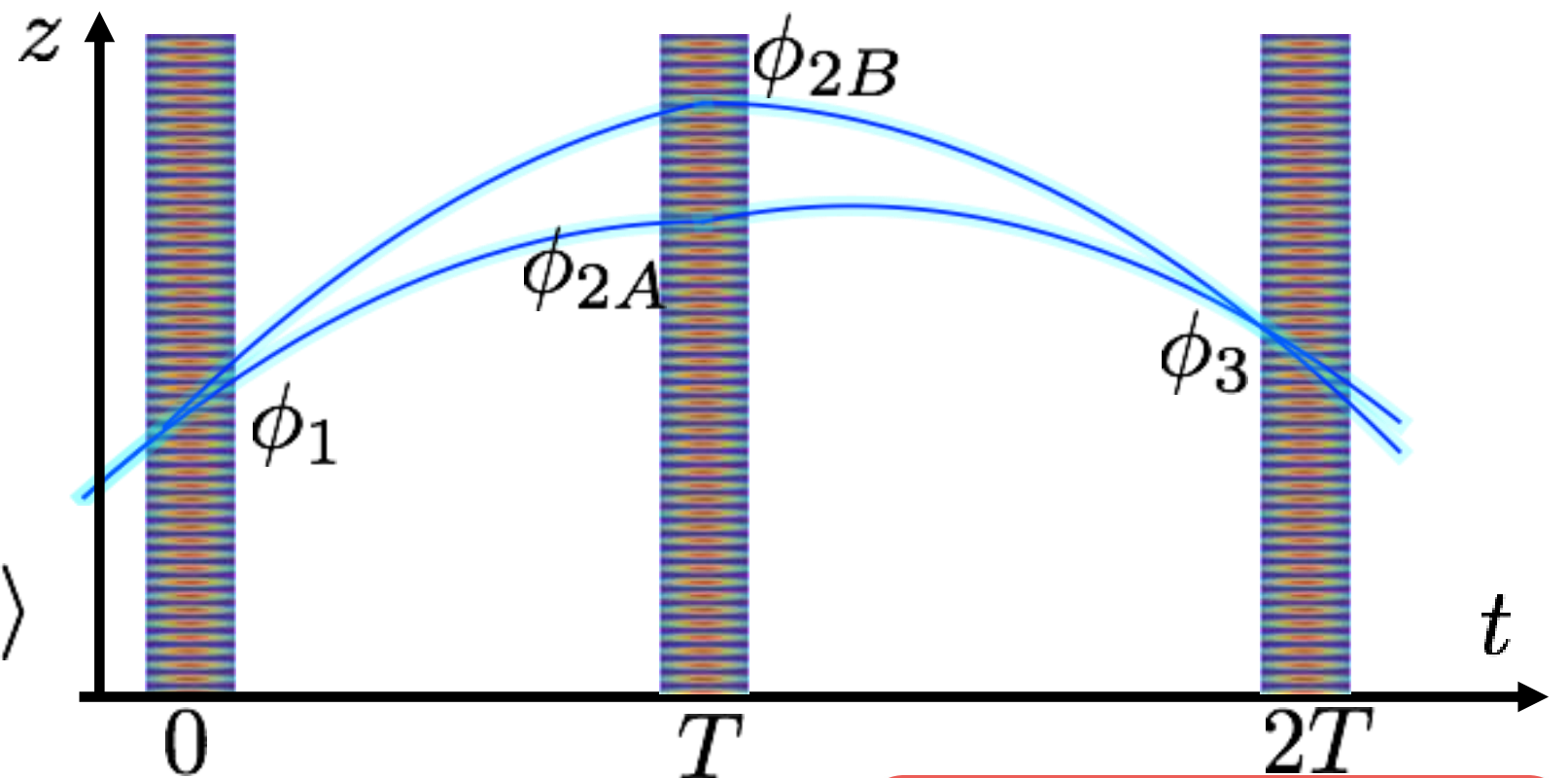
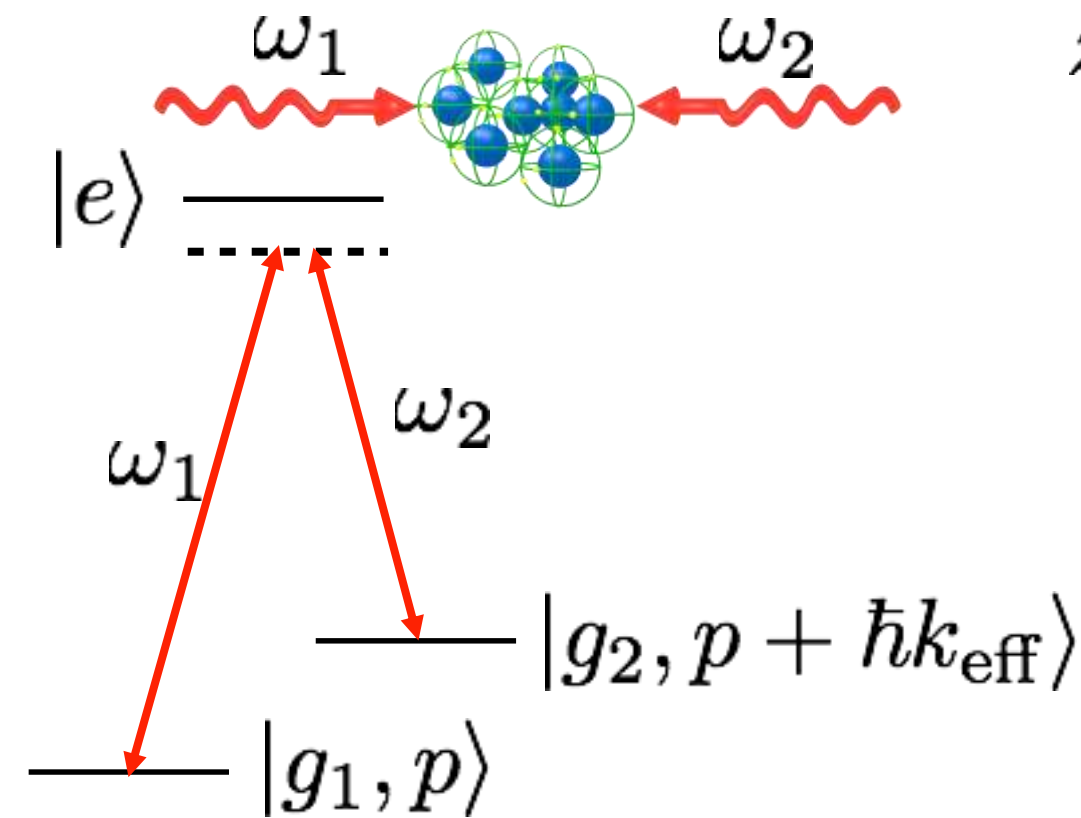
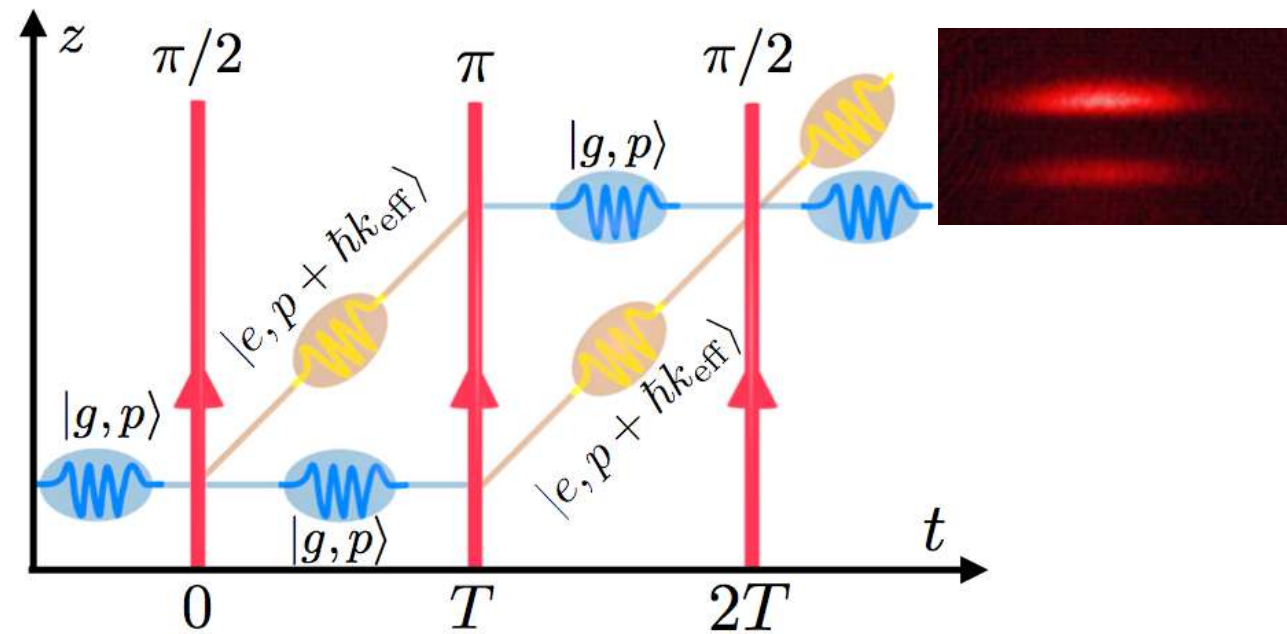
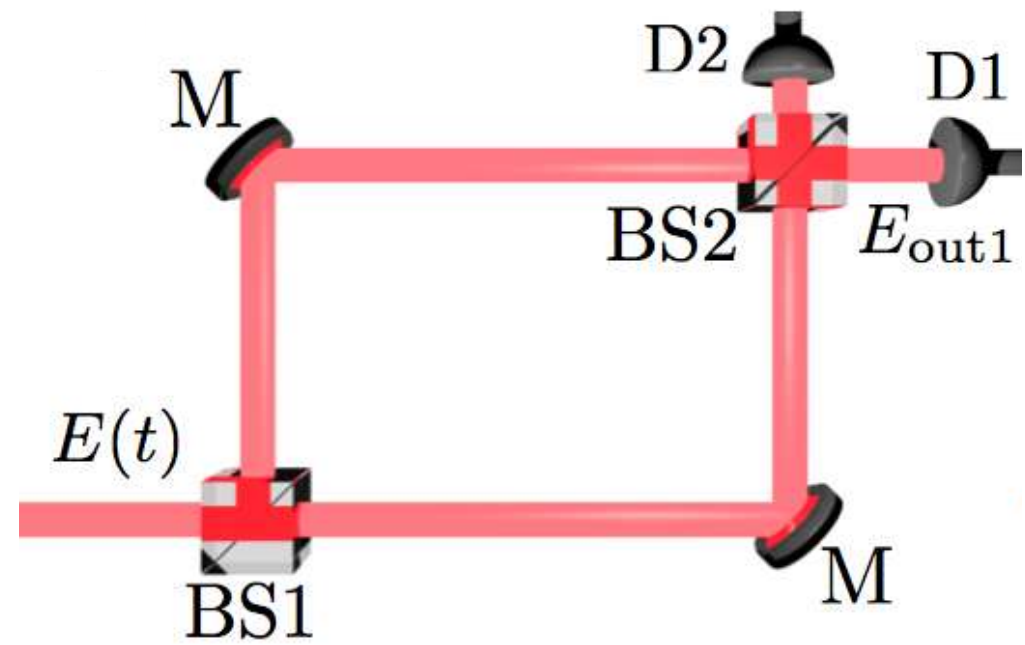


Testing gravity with atomic quantum sensors

L. Salvi, L. Cacciapuoti, G. D'Amico, M. Fattori, L. Hu, M. Jain, G. Rosi, E. Wang, N. Poli, G. M. Tino
September 20, 2018

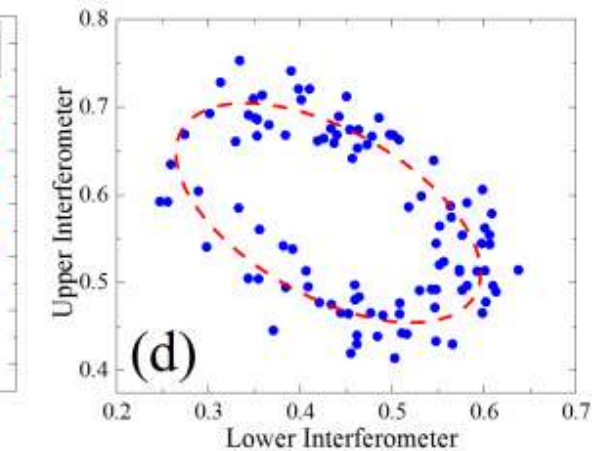
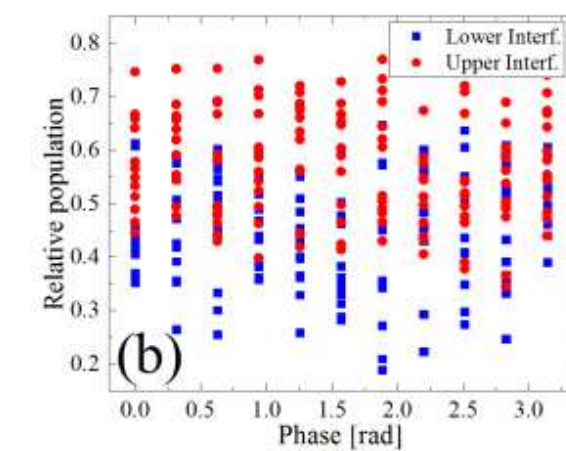
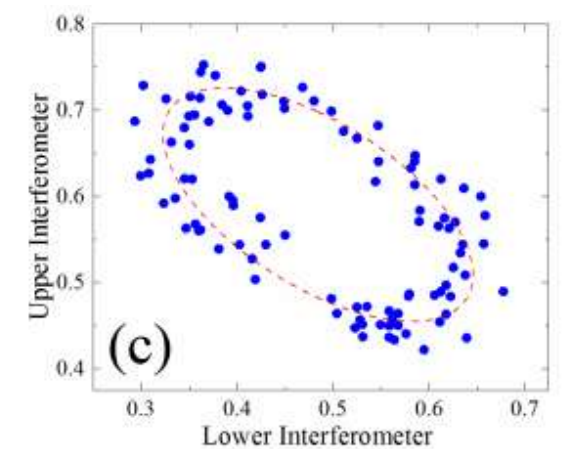
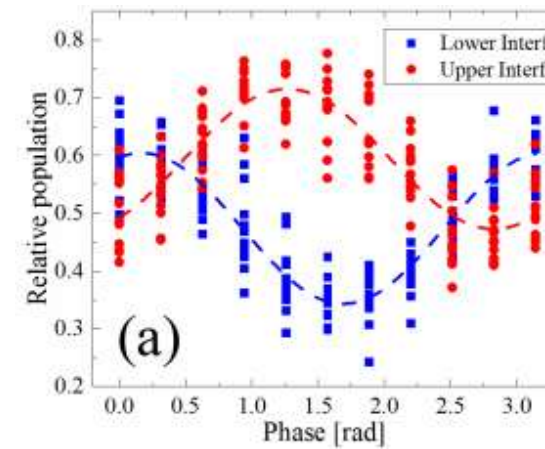
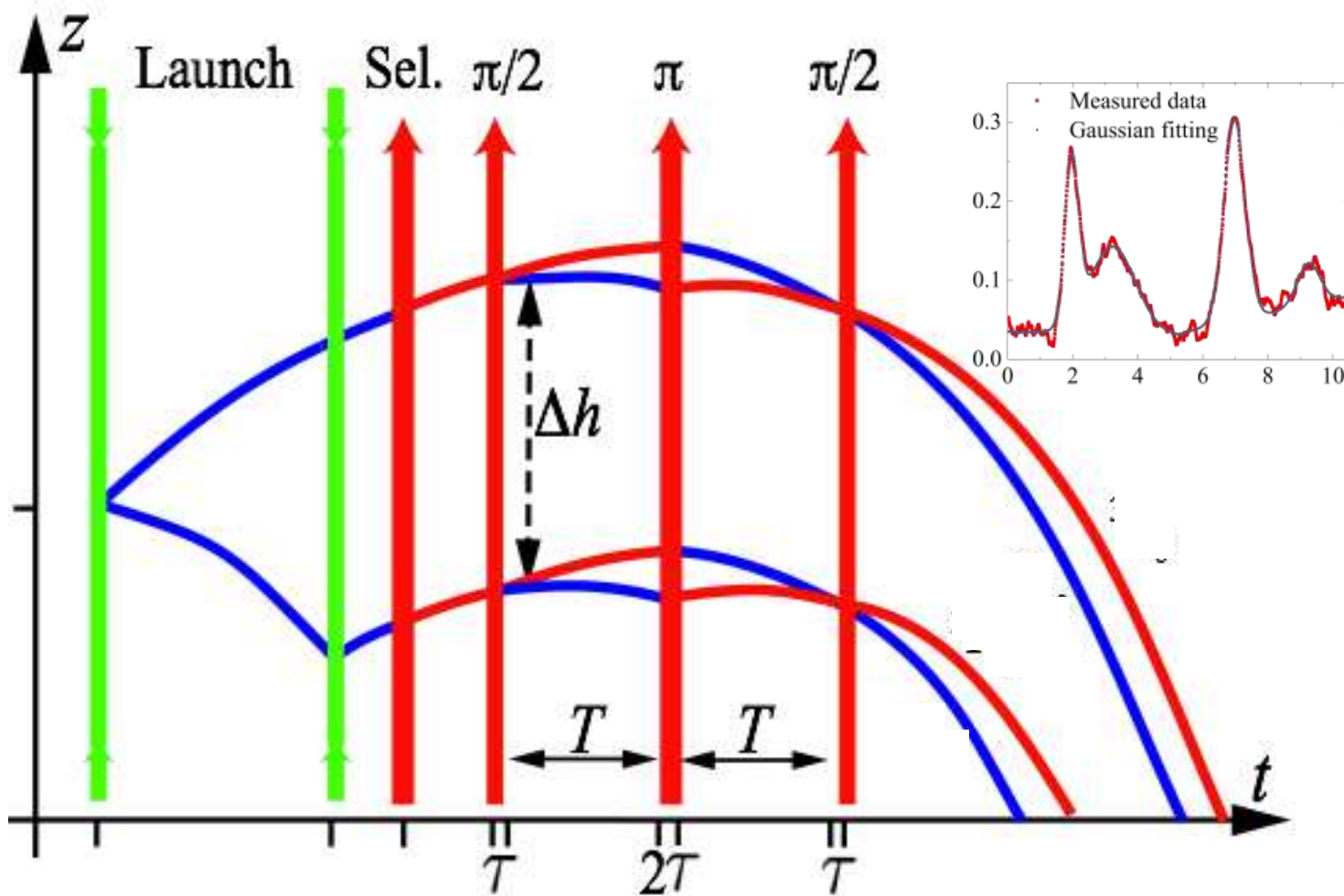
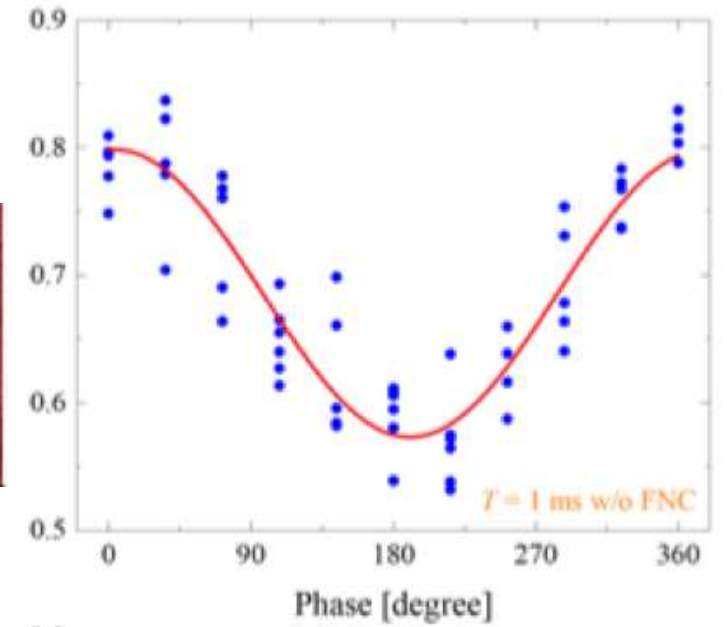
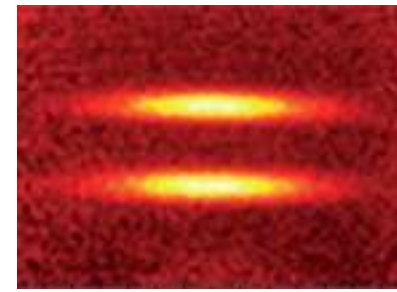
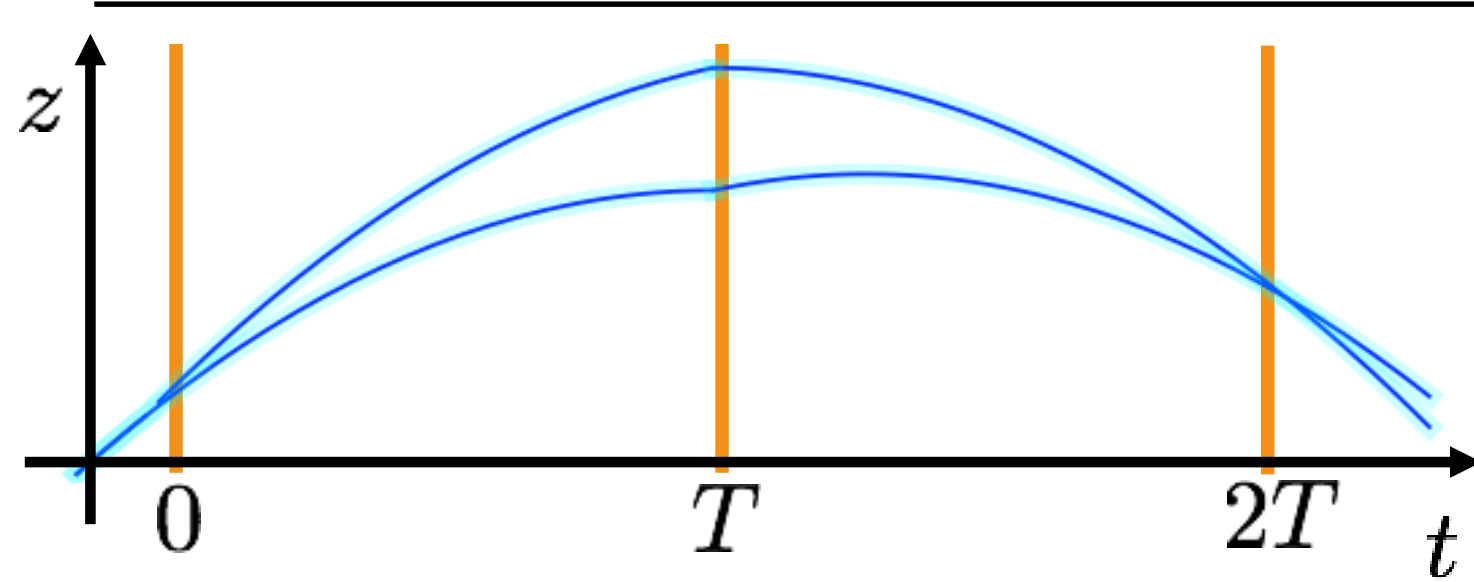


Atom Interferometry and gravity measurements



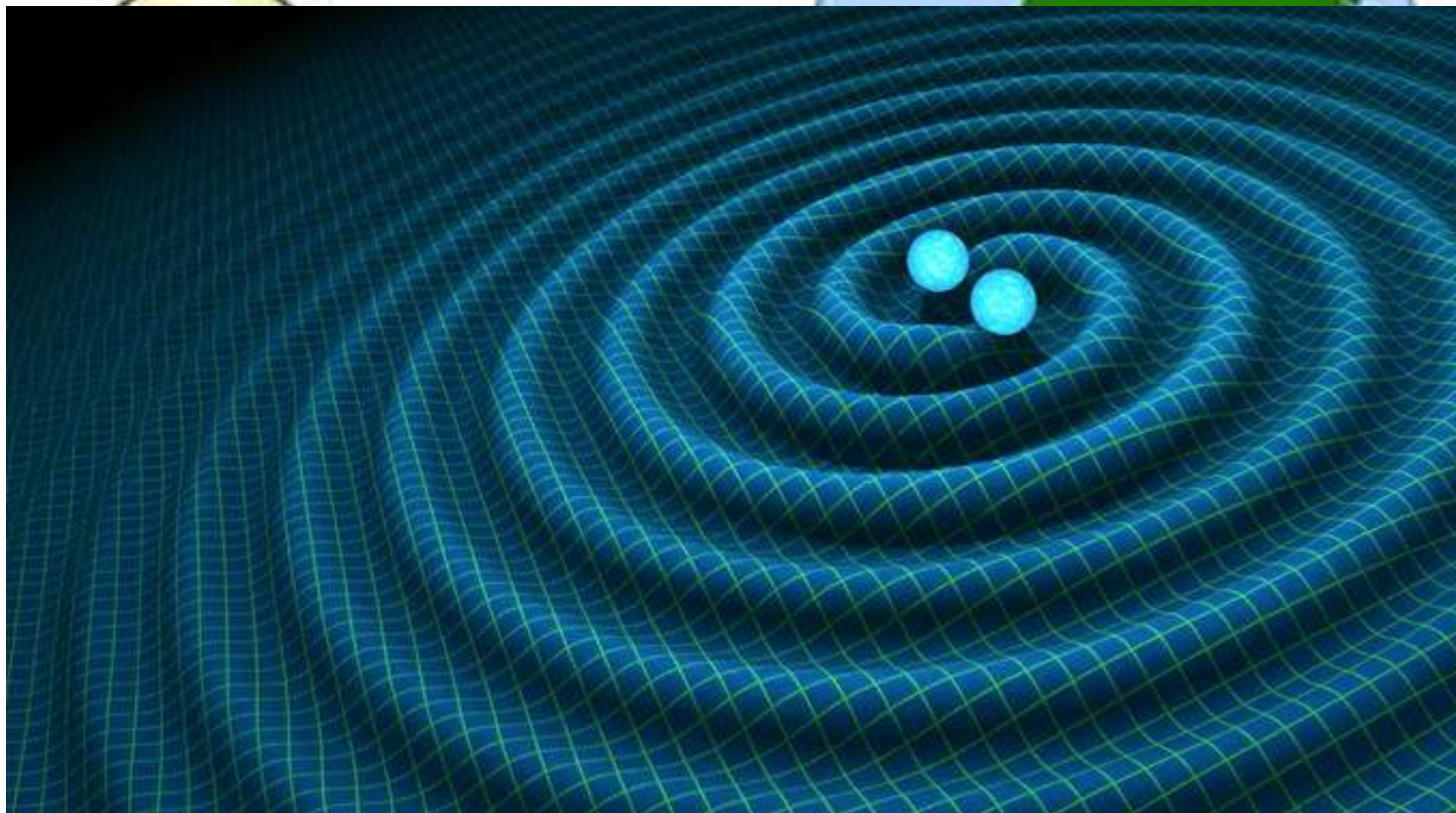
$$\Delta\phi = (\phi_1 - \phi_{2A}) - (\phi_{2B} - \phi_3) \longrightarrow \Delta\phi = k_{\text{eff}} g T^2$$

Gravity and gradient measurements



Tests of gravity with Atom Interferometry

- Equivalence principle tests at the quantum level
- Measurement of the Newtonian constant G
- Gravity differences at large distances



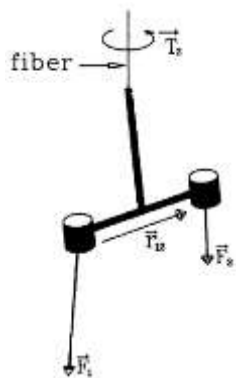
Test of the Weak Equivalence Principle

Einstein Equivalence Principle:

- **Universality of free fall or Weak Equivalence Principle** →
- Local position invariance
- Local Lorentz invariance

The trajectory of a chargeless body is independent of its internal structure and composition

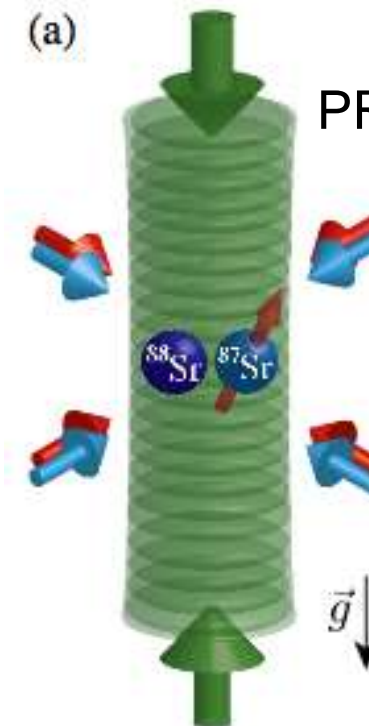
$$\eta = 2 \left| \frac{a_A - a_B}{a_A + a_B} \right| = 2 \left| \frac{(m_i/m_g)_A - (m_i/m_g)_B}{(m_i/m_g)_A + (m_i/m_g)_B} \right|$$



Wagner et al.
CQG **29** 18 (2012)



Touboul et al.
PRL **119** 231101 (2017)



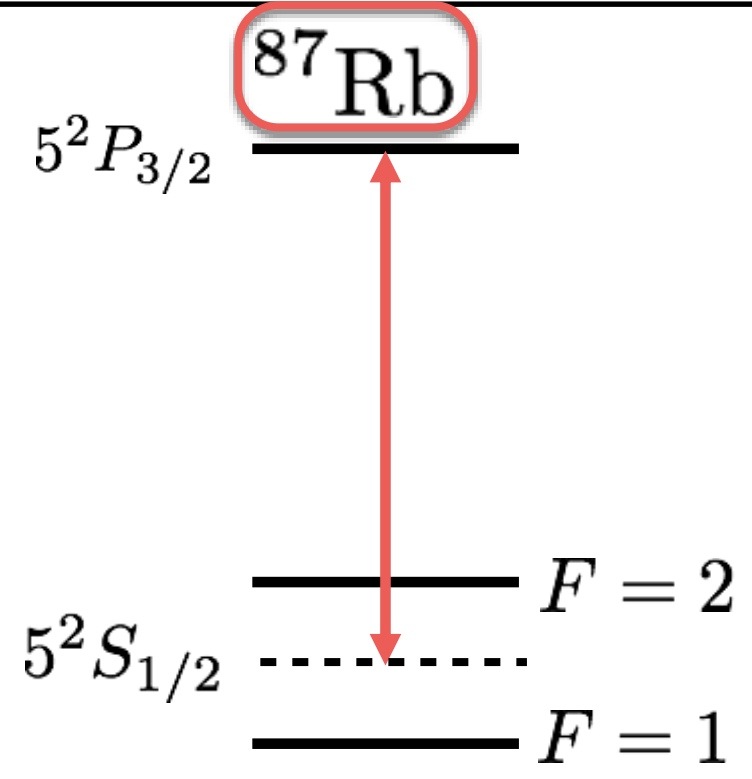
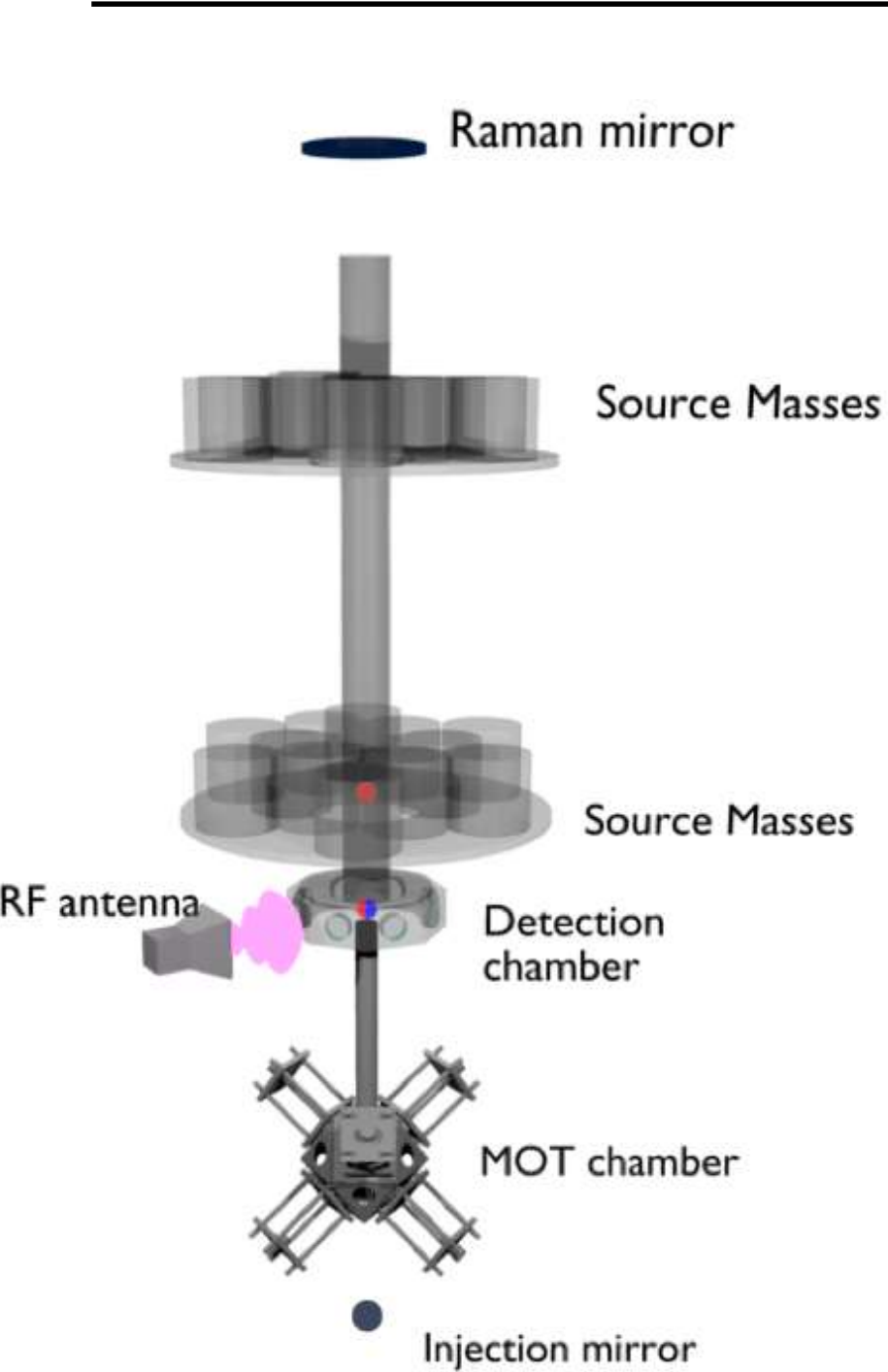
Tarallo et al.
PRL **113** 023005
(2014)

Quantum formulation of the Weak Equivalence Principle for two-level systems

$$\hat{M}_g \hat{M}_i^{-1} = \begin{pmatrix} r_1 & r \\ r^* & r_2 \end{pmatrix}$$

The off-diagonal elements can only be tested through coherent superpositions

Quantum test of WEP



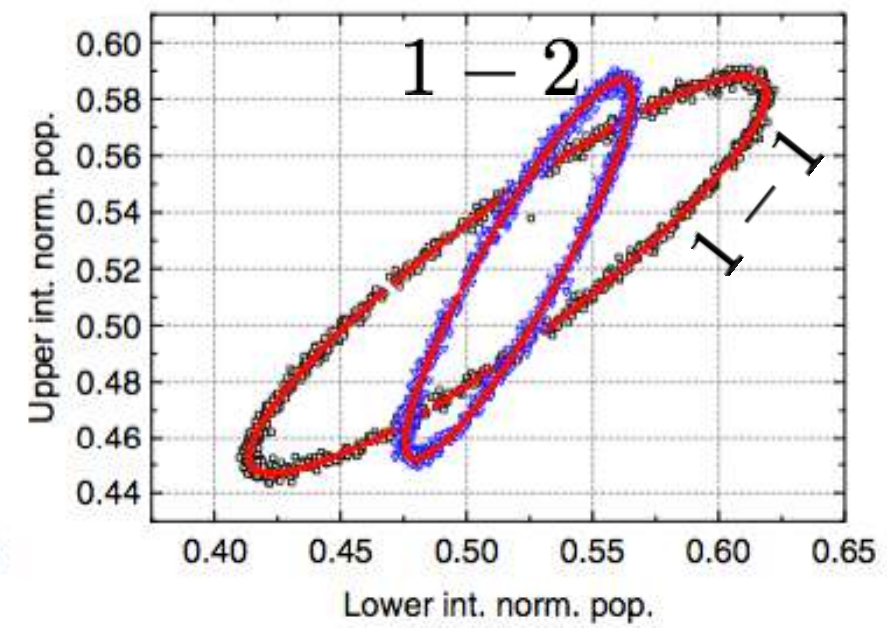
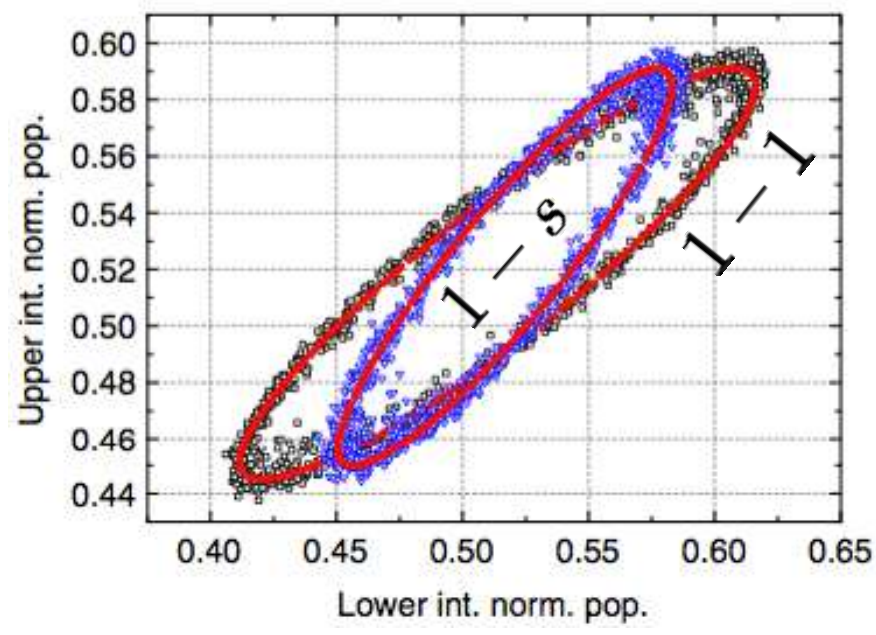
$$a_1 = g \langle 1 | \hat{M}_g \hat{M}_i^{-1} | 1 \rangle = gr_1$$

$$a_2 = g \langle 2 | \hat{M}_g \hat{M}_i^{-1} | 2 \rangle = gr_2$$

$$|s\rangle = \frac{1}{\sqrt{2}} (|1\rangle + e^{i\gamma} |2\rangle)$$

$$a_s = \langle s | \hat{M}_g \hat{M}_i^{-1} | s \rangle$$

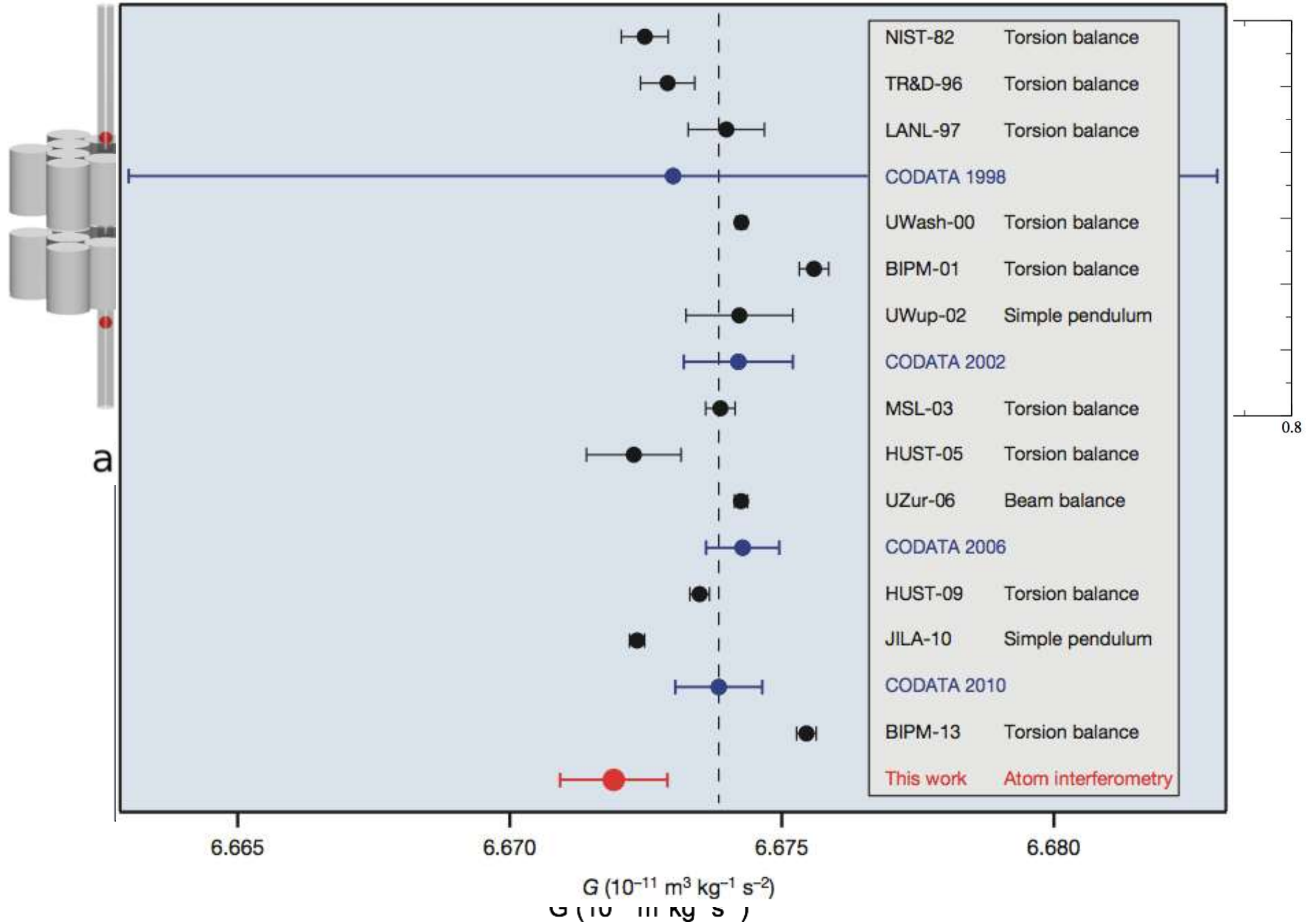
$$= g \left[\frac{r_1 + r_2}{2} + |r| \cos(\phi_r + \gamma) \right]$$



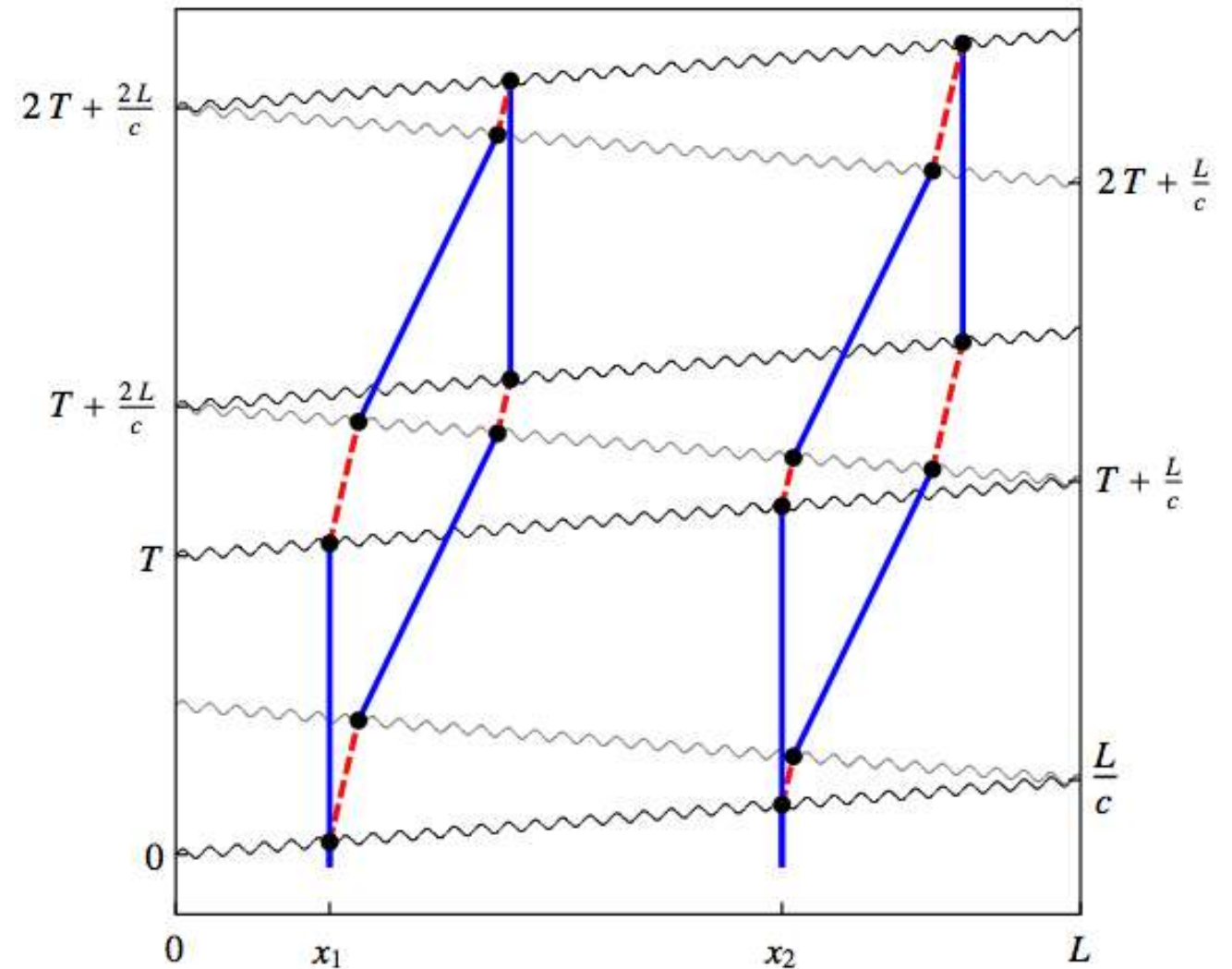
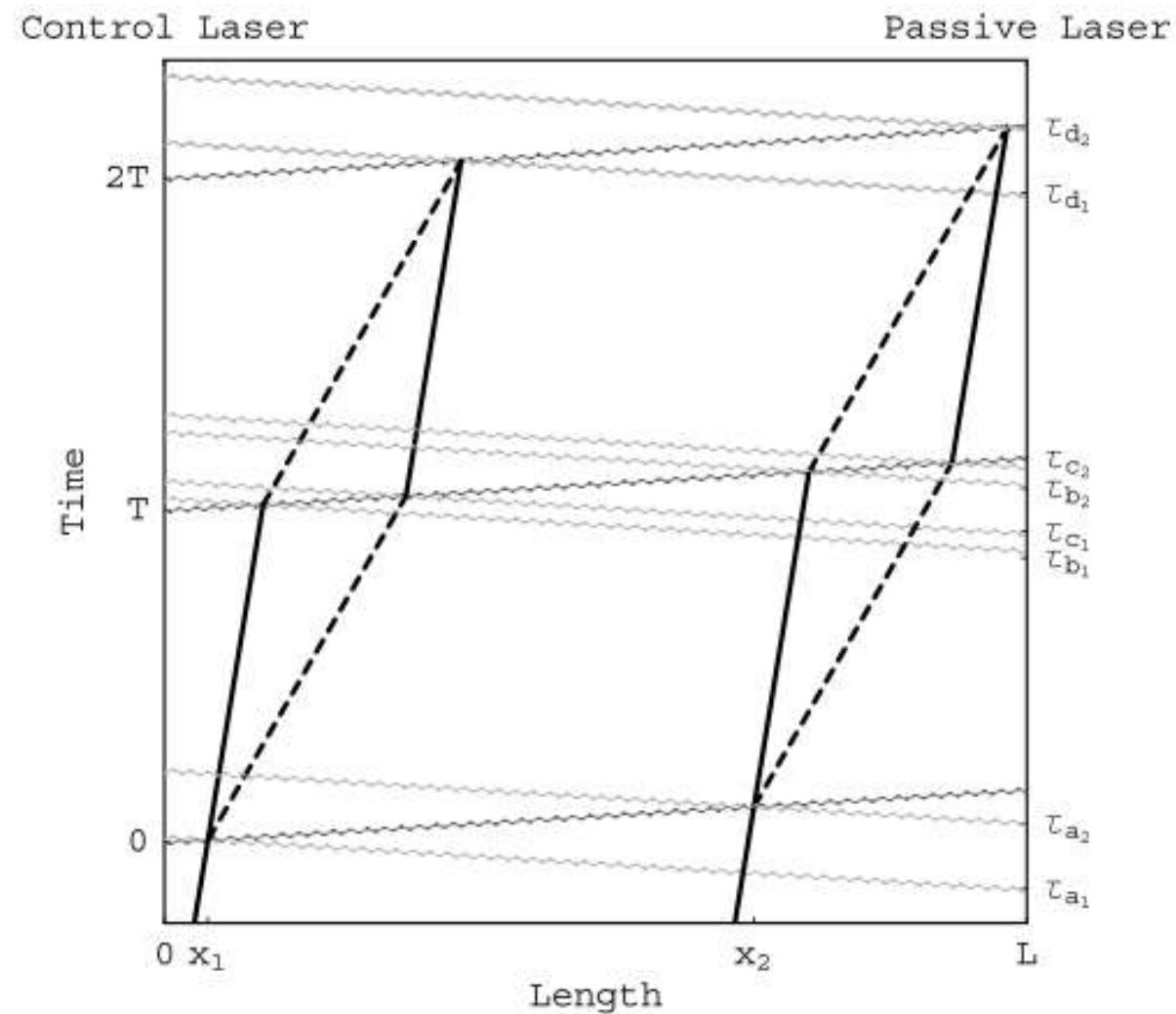
$$\eta_{1-2} = (1.0 \pm 1.4) \times 10^{-9} \quad \eta_{1-s} = (3.3 \pm 2.9) \times 10^{-9} \quad |r| \leq 5 \times 10^{-8}$$

Rosi et al. Nature Communications **8**, 15529 (2017)

Measurement of the Newtonian constant G



Long-baseline atom interferometry



Detector bandwidth limit:

$$\Delta\nu_{\max} \sim c/L$$

Absence of phase noise from the laser:

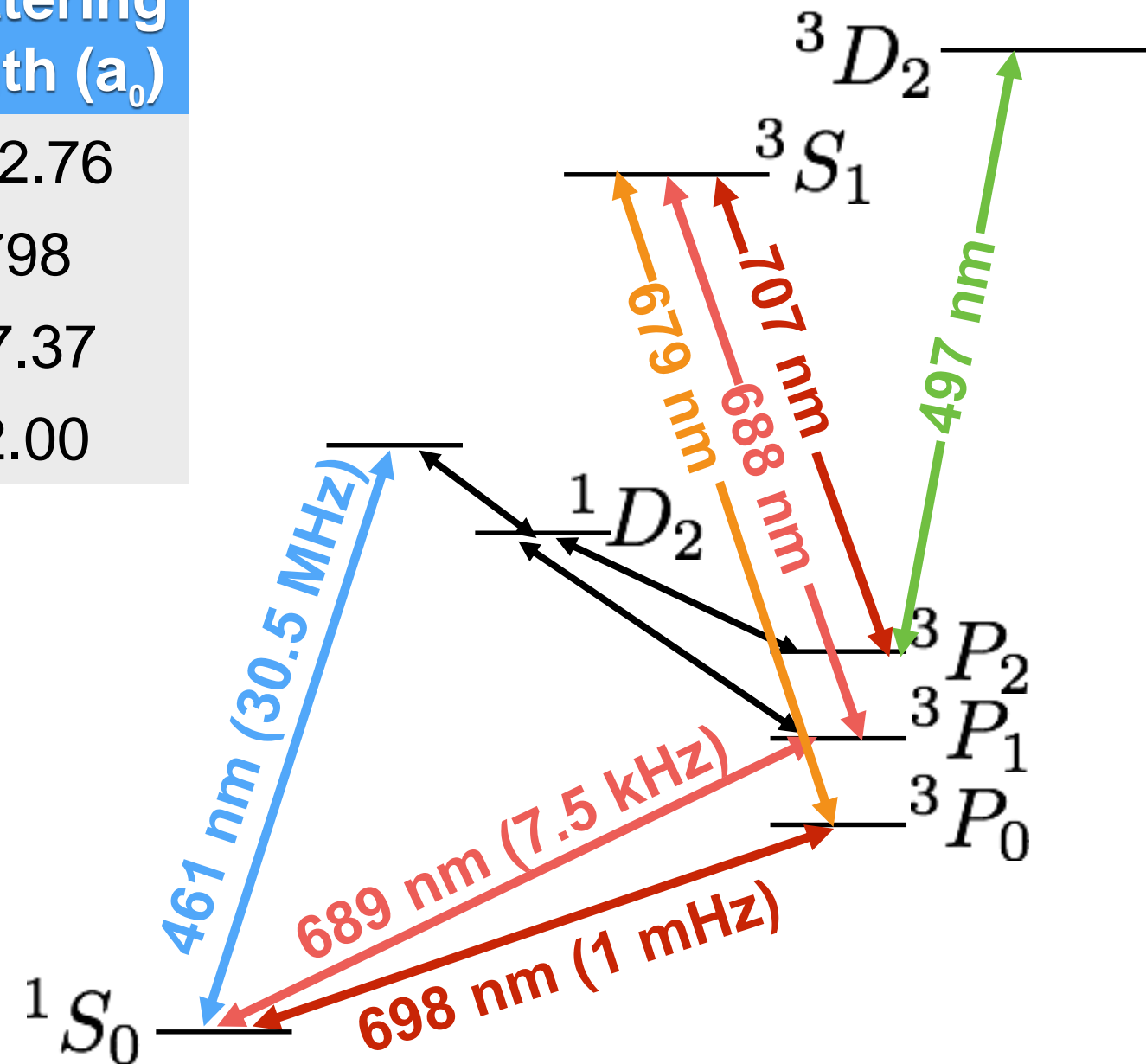
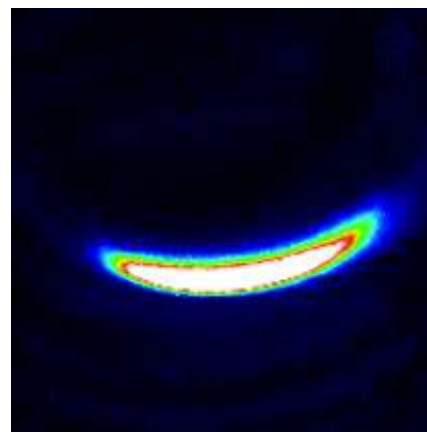
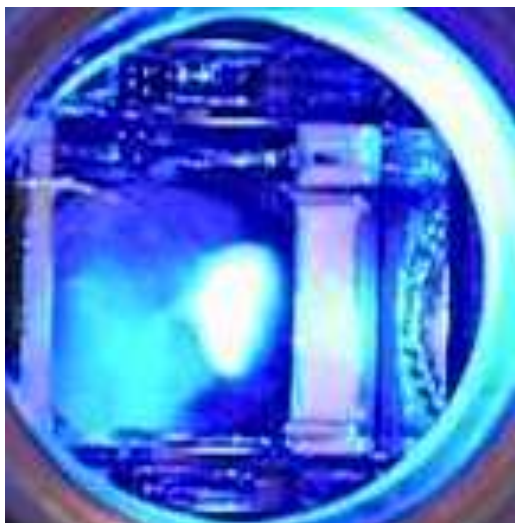
single-photon transition required

- S. Dimopoulos et al., *Phys. Lett. B* **678**, 37-40 (2009)
- P. W. Graham et al., *Phys. Rev. Lett.* **110**, 171102 (2013)

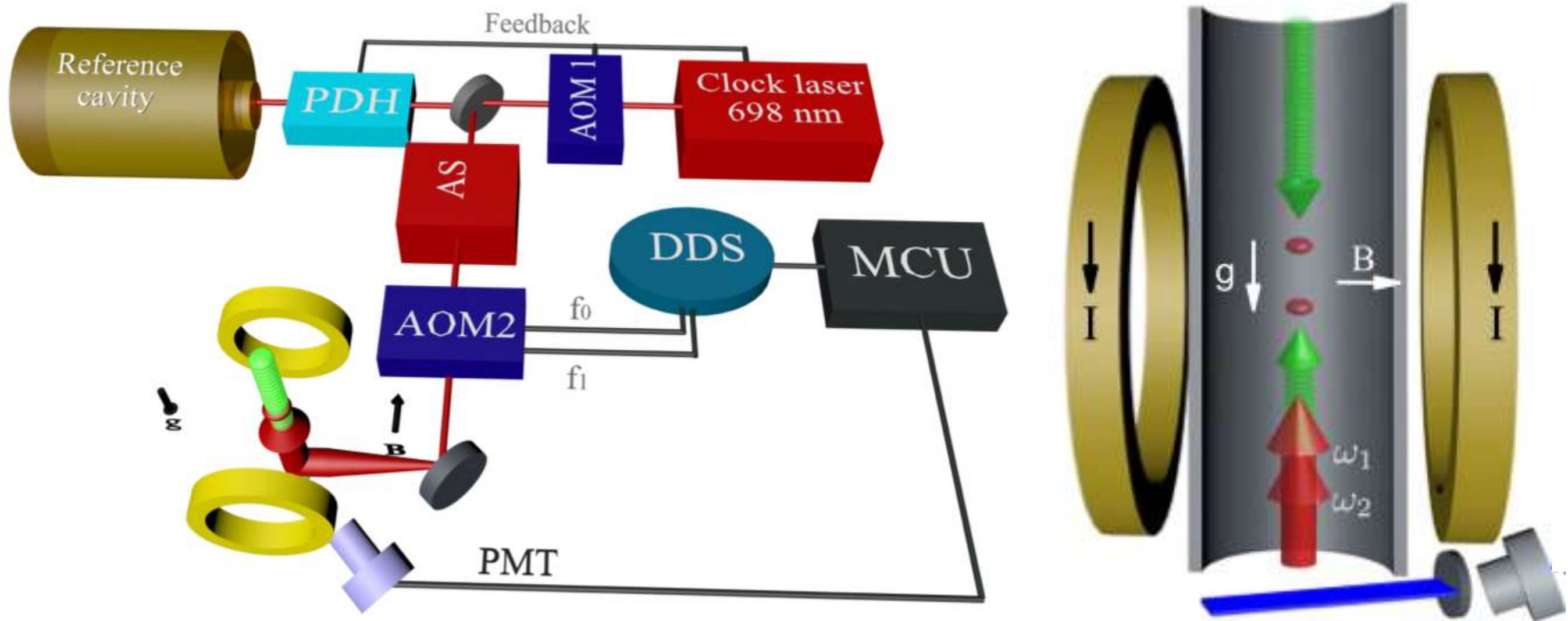
Requirements for the AI based on a single-photon transition

- Lifetime of the excited state longer than the interferometer time
- Ultra-stable laser to address the transition
- High-power laser to attain good transfer efficiency

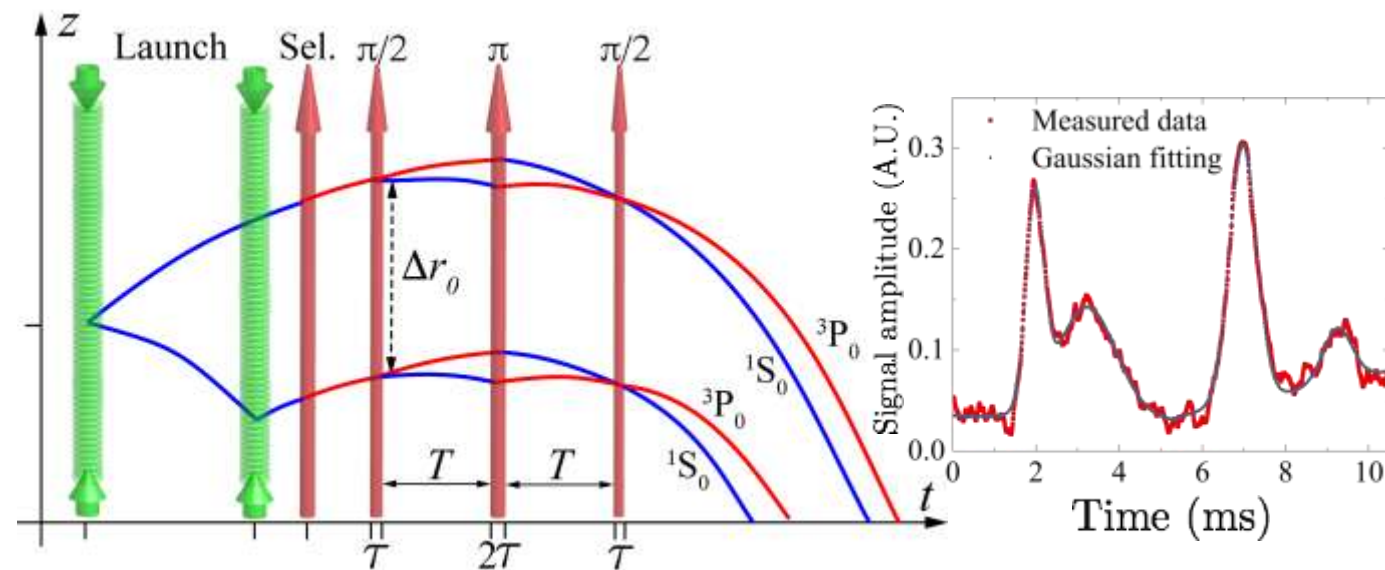
Atomic mass	Abundance	Nuclear spin	Scattering length (a_0)
84	0.56%	0	122.76
86	9.86%	0	798
87	7.00%	9/2	97.37
88	82.58%	0	-2.00



Experiment setup for proof of principle



Gradiometer setup and signal



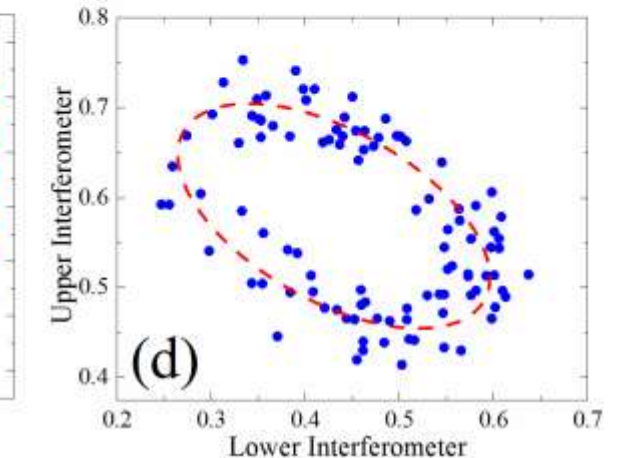
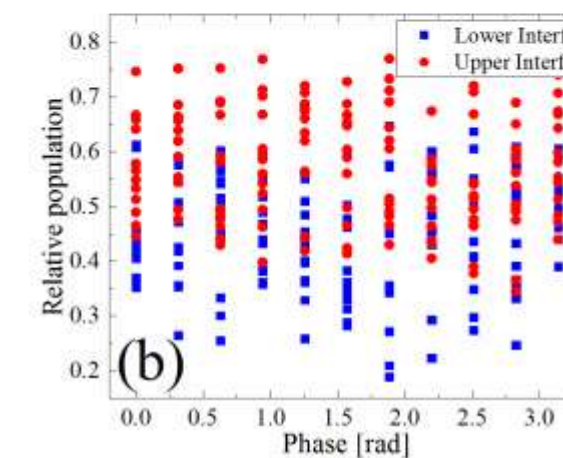
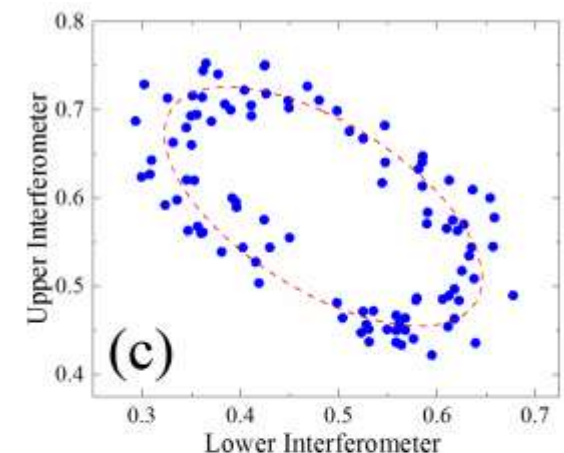
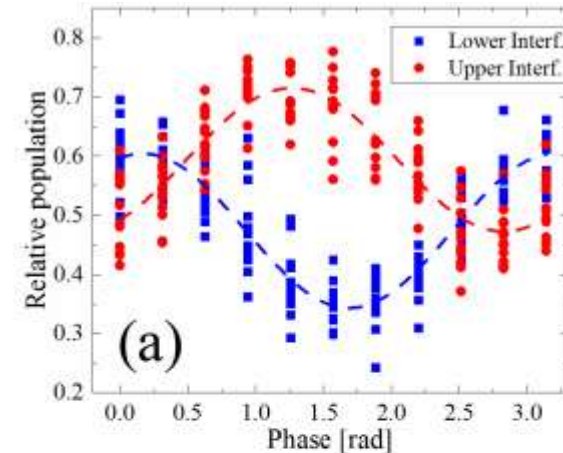
$$\Delta r_0 = 1.9 \text{ mm}$$

$$\Delta v_0 = 3.4 \text{ cm/s}$$

Because of the velocity difference, a single laser frequency cannot interact with both probes.

With two frequencies:

- both clouds addressed
- tunable artificial phase shift



Attained phase noise 5 times greater than the atom shot noise due to limited detection efficiency

Conclusions and prospects

- Quantum test of the Equivalence principle
- Accurate measurement of G with a completely new method, **attained the atom shot noise level**
- New type of interferometer operating on a single-photon transition

Future work

- Implementation of entangled (squeezed states) to overcome the shot noise
- Quantum test of WEP with superpositions of states separated by the optical clock transition

Thank you for your attention