

# Beta decay studies of the most exotic nuclei

Giovanna Benzoni

INFN sezione di Milano (Italy)

Outline of the lessons:

**PARTI: General concepts**

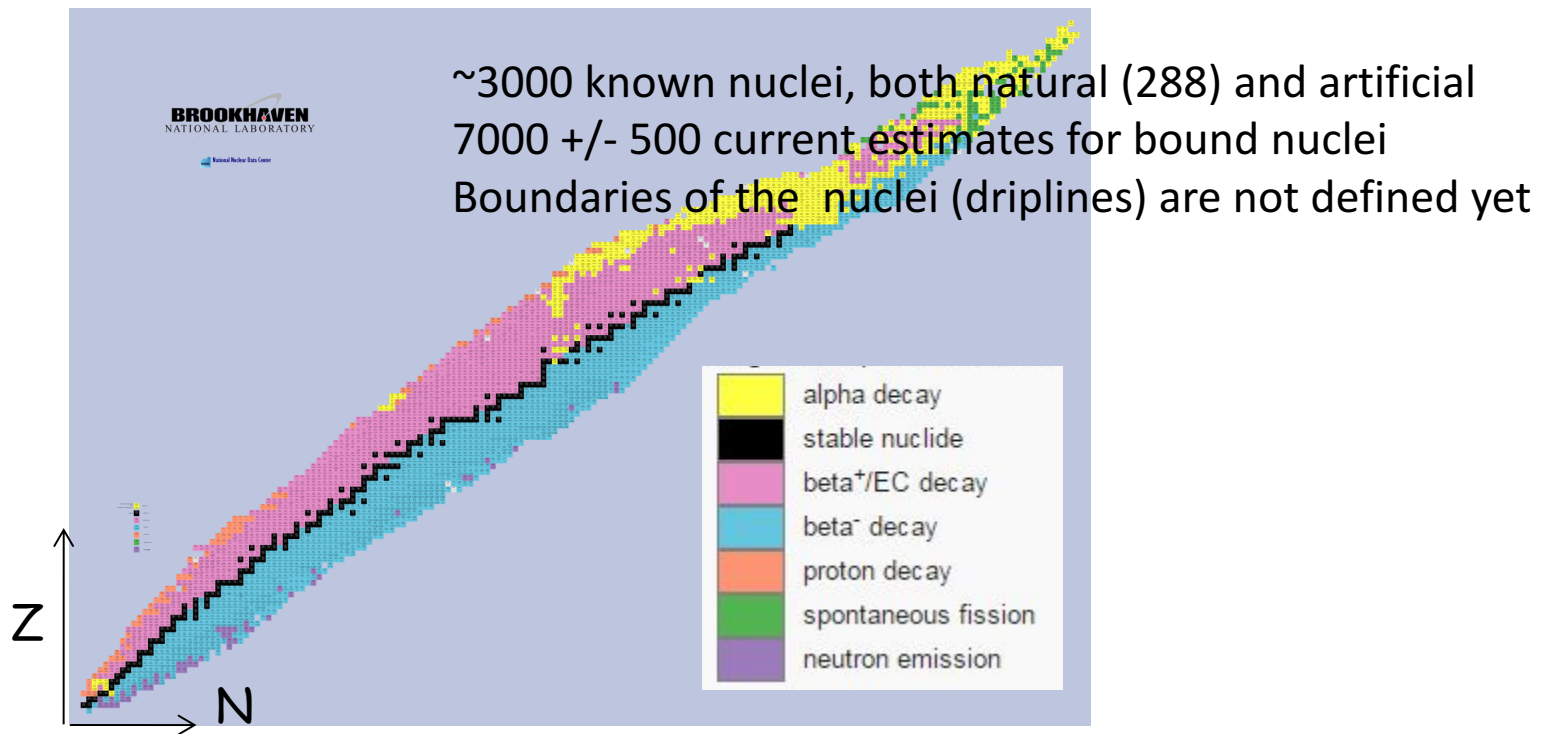
**How to measure  $\beta$  decay in exotic nuclei**

PARTII:  $\beta$  decay gross properties  $T_{1/2}$  and  $P_n$

PARTIII: High resolution vs TAS



# Introduction



$\beta$  decay dominated by weak interaction and its properties directly relate to fundamental questions:

- Properties of the neutrinos → energy spectrum, double- $\beta$  decay, sterile neutrinos
- Abundances of elements in universe → half lives and competing decay modes

and applications:

- PET scanners
- Radioactive sources (science and medicine)
- Homeland security
- Nuclear waste

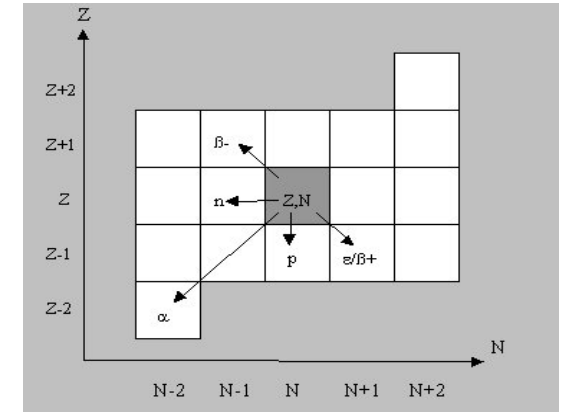
## Decay modes

$\alpha$  decay:  ${}^A_ZX \rightarrow {}^{A-4}_{Z-2}Y + \alpha$

$\beta^-$  decay:  ${}^A_ZX \rightarrow {}^A_{Z+1}Y + e^- + \bar{\nu}_e$

$\beta^+$  decay:  ${}^A_ZX \rightarrow {}^A_{Z-1}Y + e^+ + \nu_e$

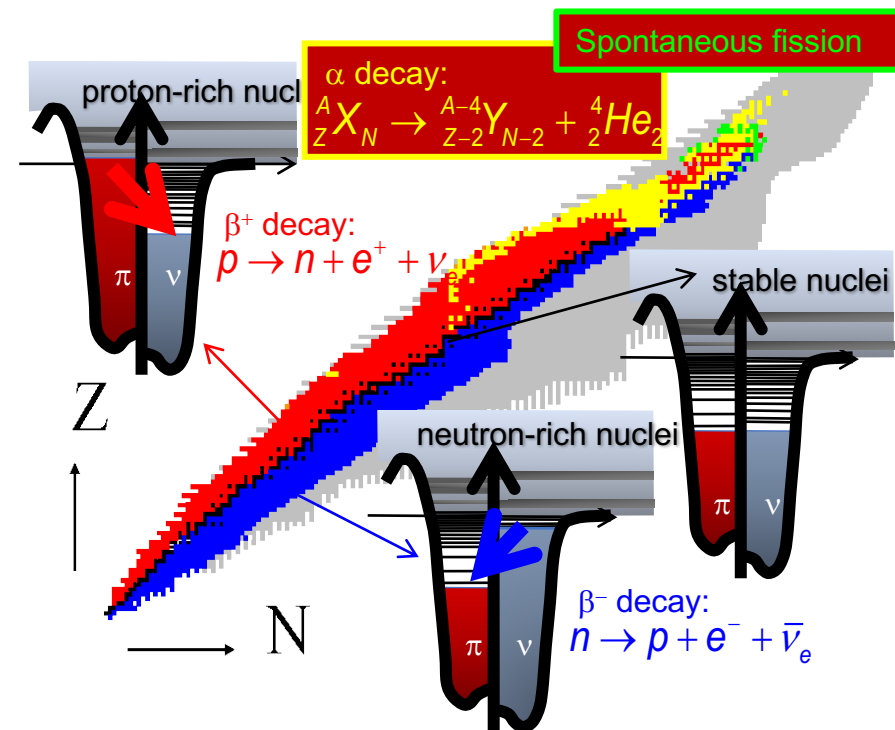
EC:  ${}^A_ZX + e^- \rightarrow {}^A_{Z-1}Y + \nu_e$



Spontaneous Fission

Exotic decay modes: n decay  
p decay  
cluster emission

$\beta$ -mediated decays:  
 $\beta$ -delayed n/p emission  
 $\beta$ -delayed fission

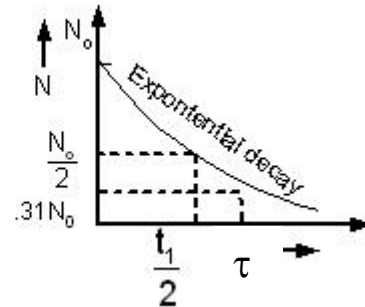


# Introduction

Useful definitions for decays:

Law of radioactive decay:

$$N(t) = N_0 e^{-\lambda t}$$



$\lambda$  decay rate

$\tau = 2\pi/\lambda$  meanlife

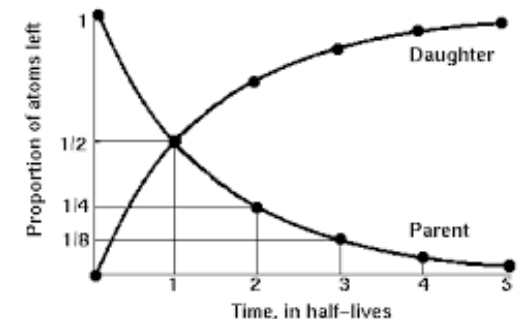
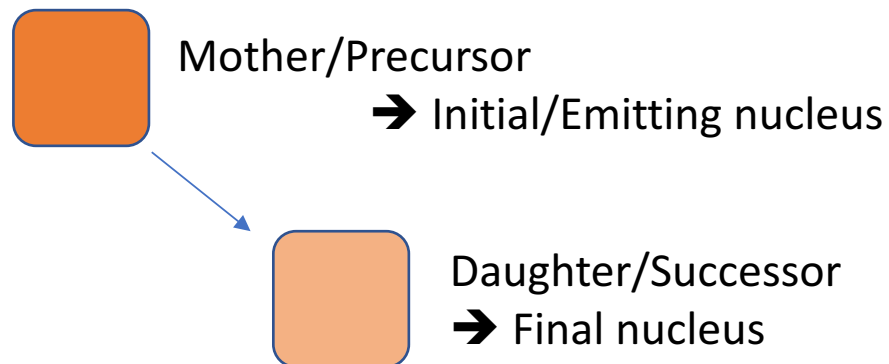
$T_{1/2} = \ln 2 * \tau$  half-life

$$T_{1/2} = \frac{\ln 2}{\lambda} \approx \frac{0.693}{\lambda} \approx 0.693\tau$$

Radioactive half-life      Radioactive decay constant      Mean lifetime

$$Q_{\text{value}} = M_{\text{initial}} - M_{\text{final}}$$

Branching ratio: decay probability of two or more competing processes





**Bateman equations:** mathematical model describing abundances and activities in a **decay chain** as a function of time, based on the **decay rates and initial abundances**. If at a time  $t$ , there are  $N_i(t)$  atoms of the  $i$ -th isotope which decays into the  $i+1$  one with a decay rate  $\lambda_i$ , the amounts of isotopes in the  $k$ -th step of the decay chain evolves as:

$$\begin{aligned}\frac{dN_1(t)}{dt} &= -\lambda_1 N_1(t) \\ \frac{dN_i(t)}{dt} &= -\lambda_i N_i(t) + \lambda_{i-1} N_{i-1}(t) \\ \frac{dN_k(t)}{dt} &= \lambda_{k-1} N_{k-1}(t)\end{aligned}$$

This can be written in an explicit form as:

$$N_n(t) = \prod_{j=1}^{n-1} \lambda_j \sum_{i=1}^n \sum_{j=1}^n \left( \frac{N_i(0) e^{-\lambda_j t}}{\prod_{p=1, p \neq j}^n (\lambda_p - \lambda_j)} \right)$$

This can be extended to cases in which we have many branches

# Introduction

## Secular equilibrium:

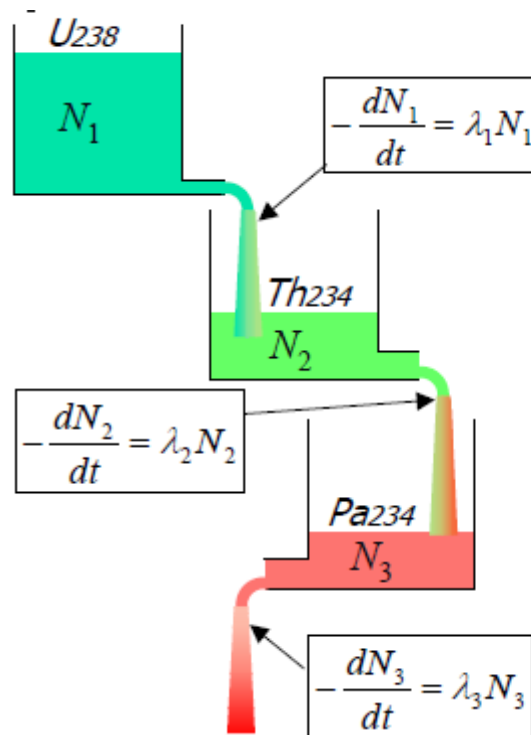
The rate of production and decay of an element is constant

It can be achieved in a decay chain if the daughter B has a much shorter half-life than the mother A

The decay rate of A, which is translated into the production rate of B is constant

The quantity of B accumulates since the number of nuclei which decay in 1 s is equal to that produced in 1 s.

This depends on  $N_A(t=0)$ ,  $\lambda_A$  and  $\lambda_B$



- Each member of the family is a dewar. They are connected and they fill each other in chain
- The rate of emptying ( $-dN/dt$ ) is a function of the level in the dewar  $N$  and the decay rate  $\lambda$
- At equilibrium the rate of evacuation in each dewar is equal
- In nuclear physics the rate of decay is the Activity and depends on the decay constant  $\lambda$
- At equilibrium all activities are equal

# Introduction

## Detecting $\alpha$ $\beta$ $\gamma$ n particles: reminder of basic properties

$\alpha$  particle: massive, charged

- interested in energy, angular distribution
- spectrometers, Si, plastic detectors

$\beta$  particles: light, charged

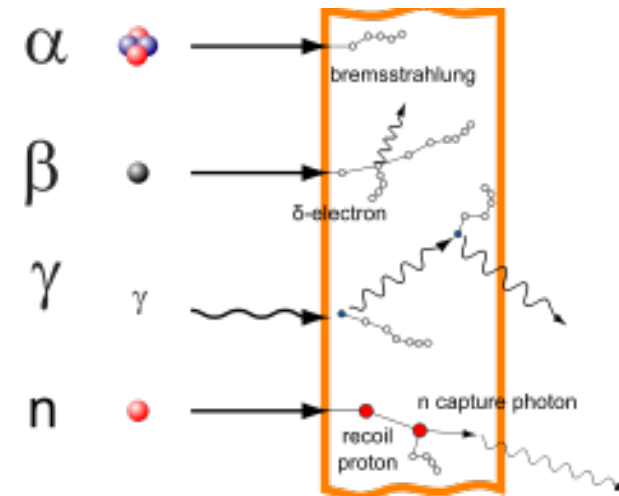
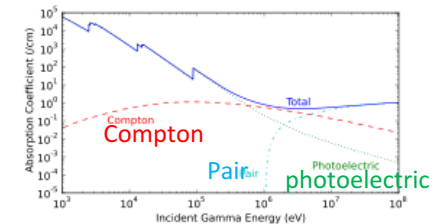
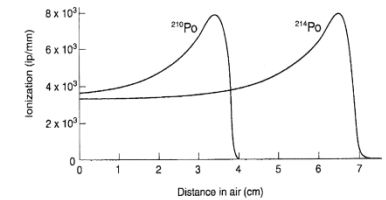
- interested in energy and correlations
- $e^+$  annihilates into two 511-keV  $\gamma$  rays
- spectrometers, Si, plastic detectors

$\gamma$  particles: massless, neutral

- interested in energy and angular distributions
- HPGe detectors, inorganic scintillators

Neutrons: massive, neutral

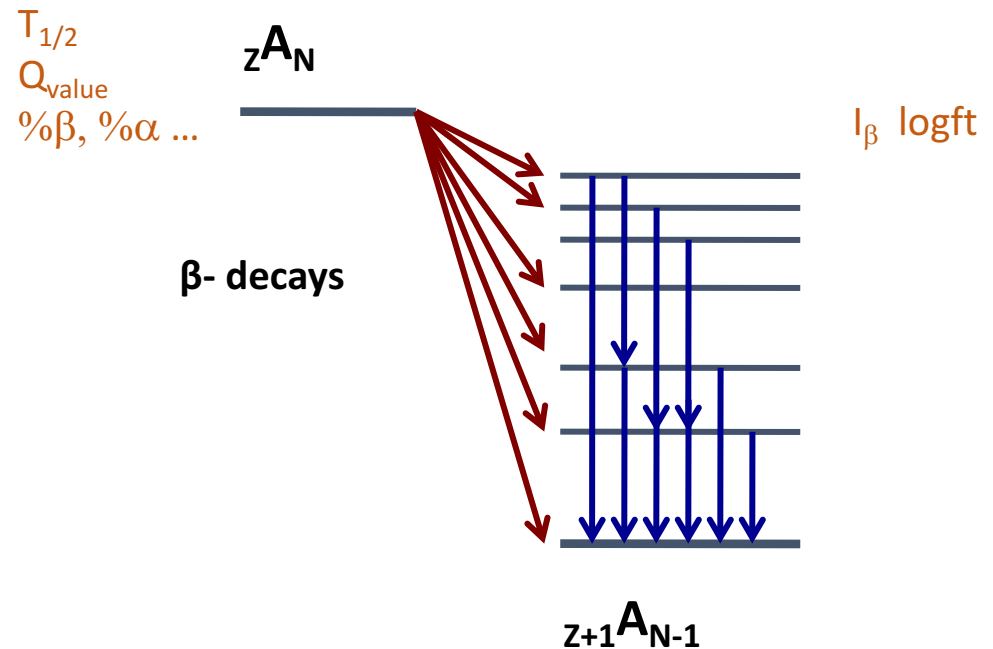
- interested in energy and correlations
- measured indirectly via capture reactions
- liquid scintillators, inorganic scintillators



# Introduction

Example of detailed study of a decay:

- 1) Produce and Identify PARENT nucleus: reactions/ISOL/fragmentation
- 2) Detect the decay: identify the emitted particle  $\alpha/\beta$ 
  - identify competing mechanisms
  - identify decaying state, g.s. or excited isomeric state
  - measure half-life
- 3) Study properties of DAUGHTER nucleus:
  - which levels are populated
  - Branching Ratio btw levels
  - subsequent decays

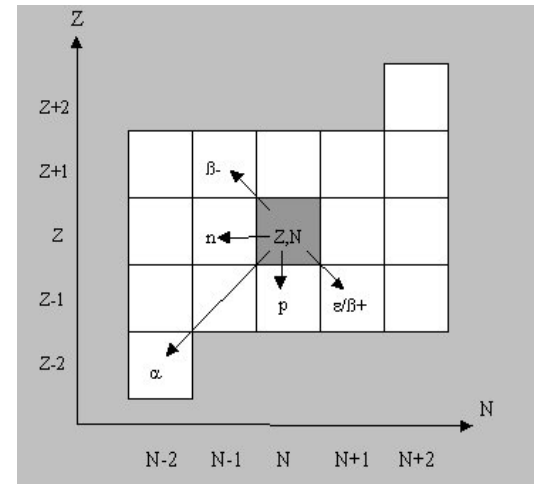


# General Properties of $\beta$ decay

$${}^A_ZX \rightarrow {}^A_{Z+1}Y + e^- + \nu_e \quad \beta^- \text{ decay}$$

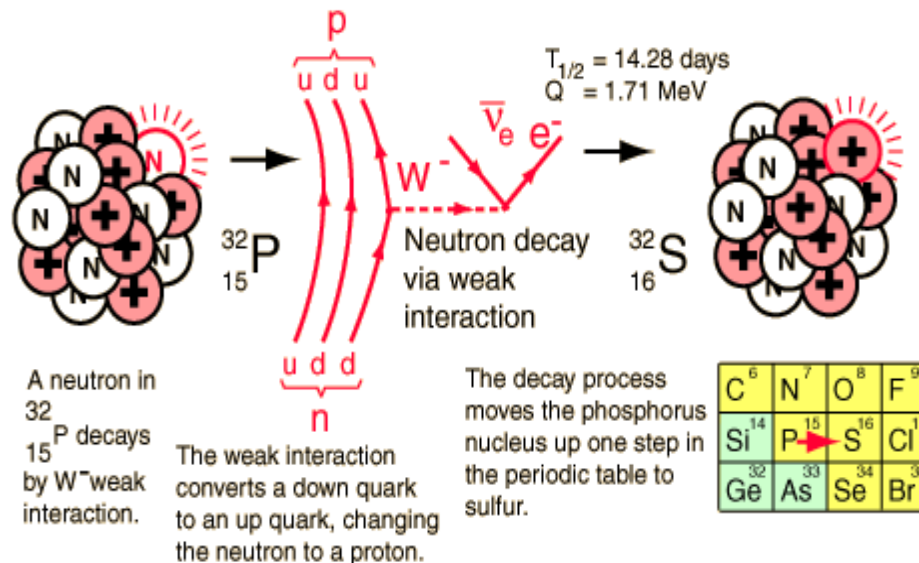
$${}^A_ZX \rightarrow {}^A_{Z-1}Y + e^+ + \nu_e \quad \beta^+ \text{ decay}$$

$${}^A_ZX + e^- \rightarrow {}^A_{Z-1}Y + \nu_e \quad \text{EC}$$



$\beta$  decay is a weak interaction “semi-leptonic” decay

The quark level Feynman diagram for  $\beta^-$  decay is shown here:



# General Properties of $\beta$ decay

- The Q value in  $\beta$  decay is effectively shared between the electron and antineutrino.
- This is the case of gs  $\rightarrow$  gs decay

Q value for  $\beta^-$  decay is

$$Q_{\beta^-} = (M(A, Z) - M(A, Z + 1))c^2$$

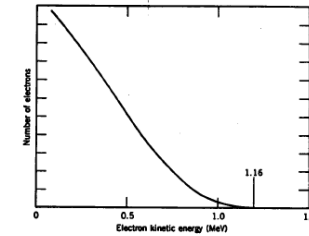


Fig. 9.1 The continuous electron distribution from the  $\beta$  decay of  $^{210}\text{Bi}$  (also 1 RaE in the literature).

$Q_{\beta^-}$  Sum of the energy of the electron (positron) and antineutrino (neutrino)

$$Q_{\beta^-} = T_{M(A, Z+1)^+} + T_e + T_\nu \approx T_e + T_\nu \quad \left( \text{since } T_{M(A, Z+1)^+} < \text{keV} \right)$$

$$Q_{\beta^-} = (T_e)_{\text{max}} = (T_\nu)_{\text{max}} \quad \text{If the other term is null}$$

$$Q_{\beta^+} = M(A, Z) - M(A, Z - 1) - 2m_e c^2$$

For electron capture:

$$Q_{EC} = (M(A, Z) - M'(A, Z - 1))c^2 - B_n$$

$B_n$  binding energy of  $n$  - shell electron

NB: EC and  $\beta^+$  are not always competing:

if  $\beta^+$ , EC is possible, the contrary is not guaranteed

Note these are atomic masses

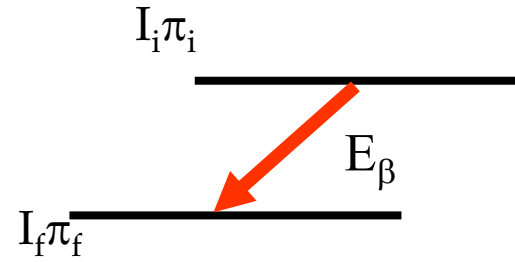
# General Properties of $\beta$ decay

$$I_i = I_f + L_\beta + S_\beta$$

$$L_\beta = l_{\beta-(\beta+)} + l_{\tilde{\nu}(\nu)}$$

$$\pi_i \pi_f = (-1)^{L_\beta}$$

$$S_\beta = S_{\beta-(\beta+)} + S_{\tilde{\nu}(\nu)} = \begin{cases} 0 \\ 1 \end{cases} \begin{matrix} \uparrow\downarrow \\ \downarrow\downarrow \text{ or } \uparrow\uparrow \end{matrix}$$



$L_\beta = n$  defines the degree of **forbiddenness** ( $n$ )

**allowed**

when  $L_\beta = n = 0$

and  $\pi_i \pi_f = +1$

Electron and neutrino do not  
carry angular momentum

**forbidden**

when the angular momentum  
conservation requires that  $L_\beta = n > 0$   
and  $\pi_i \pi_f = (-1)^{L_\beta}$

$$\Delta I = |I_i - I_f| \equiv 0, 1$$

# General Properties of $\beta$ decay

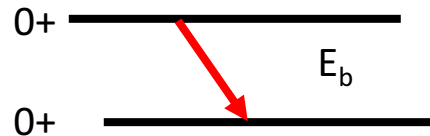
Selection rules for possible decays

Type of transition	Order of forbiddenness	$\Delta I$	$\pi_i \pi_f$
Allowed		0, +1	+1
Forbidden unique	1	$\mp 2$	-1
	2	$\mp 3$	+1
	3	$\mp 4$	-1
	4	$\mp 5$	+1
	.	.	.
Forbidden	1	0, $\mp 1$	-1
	2	$\mp 2$	+1
	3	$\mp 3$	-1
	4	$\mp 4$	+1
	.	.	.

Classification of **allowed decays**

$$(\pi_i \pi_f = +1)$$

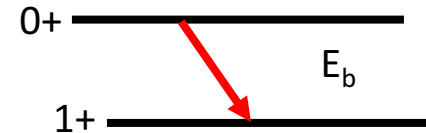
**Fermi**



$$\Delta I = |I_i - I_f| \equiv 0$$

$$L_\beta = 0 \quad S_\beta = 0 \downarrow \uparrow$$

**Gamow-Teller**



$$\Delta I = |I_i - I_f| \equiv 1$$

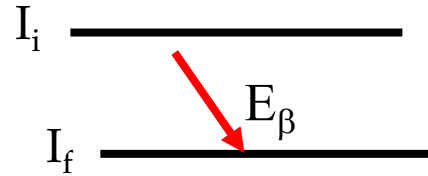
$$L_\beta = 0 \quad S_\beta = 1 \uparrow \uparrow \text{ or } \downarrow \downarrow$$



## Useful empirical rules

The fifth power beta decay rule:

the speed of a  $\beta$  transition increases approximately in proportion to the **fifth power of the total transition energy**



$$\frac{1}{\tau} \propto [(M(Z) - M(Z \pm 1))c^2]^5$$

- ❑ depends on spin and parity changes between the initial and final state
- ❑ additional hindrance due to nuclear structure effects :  
isospin, “l-forbidden”, “K-forbidden”, etc.

## Fermi Golden Rule

- Treat beta decay as a transition that depends upon the strength of coupling between the initial and final states
- Decay constant is given by Fermi's Golden Rule

$$\lambda_{\beta} = \frac{2\pi}{\hbar} |M|^2 \rho(E_f); M = \int \psi_f V \psi_i dv$$

- ➔ Electron and neutrino do not pre-exist in atom but are formed at the time of decay
- ➔ The decay is the result of the interaction btw the nucleon and the field produced by the electron-antineutrino couple ➔ weak interaction
- Perturbation theory can be applied since the interaction is “weak”
- M matrix element which couples the initial and final states
- Rate proportional to the strength of the coupling between the initial and final states factored by the density of final states available to the system electron-antineutrino

# General Properties of $\beta$ decay

$$t \equiv T_{1/2}^{\beta_i} = \frac{T_{1/2}^{\text{exp}}}{P_{\beta_i}} \quad \text{partial half-life of a given } \beta^- (\beta^+, \text{EC}) \text{ decay branch (i)}$$

$$\frac{\ln 2}{T_{1/2}^n} = \frac{g^2}{2\pi^3} \int_1^W p_e W_e (W_0 - W_e)^2 F(Z, W_e) C_n dW_e$$

$$\lambda = \frac{\ln 2}{t_{1/2}} = K |M_{if}|^2 f_o$$

$g$  – weak interaction coupling constant

$p_e$  – momentum of the  $\beta$  particle

$W_e$  – total energy of the  $\beta$  particle

$W_0$  – maximum energy of the  $\beta$  particle

$F(Z, W_e)$  – Fermi function – distortion of the  $\beta$  particle wave function by the nuclear charge

$C_n$  – shape factor

$Z$  – atomic number

$$K = 64\pi^4 m_o^5 c^4 g^2 / h^7$$

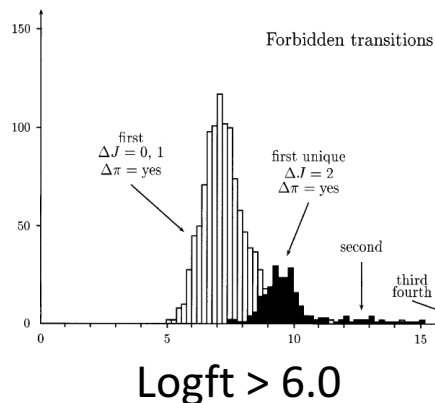
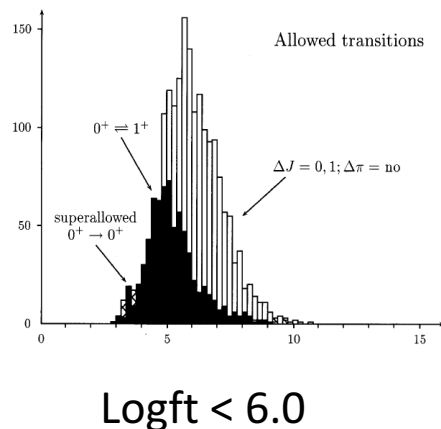
$$f_o = \int_1^{W_o} F(Z, W) W (W^2 - 1)^{1/2} (W_o - W)^2 dW$$

## Comparative Half Lives

- Based on probability of electron energy emission coupled with spectrum and the Coulomb correction  $f_0 t_{1/2}$  is called the comparative half life of a transition

$$f\tau = \frac{2\pi^3 \hbar^7 c^3}{g^2 |M_{if}|^2}$$

- Assumes matrix element is independent of energy (true for allowed transitions)  
Yields  $f\tau$  (or  $f_0 t_{1/2}$ ), comparative half-life may be thought of as the half life corrected for differences in  $Z$  and energy
- ALLOWED transitions second term is independent on nucleus,
- ➔  $f\tau$  has the same value for all allowed transitions
- For forbidden decays  $f\tau$  increases with degree of forbiddenness



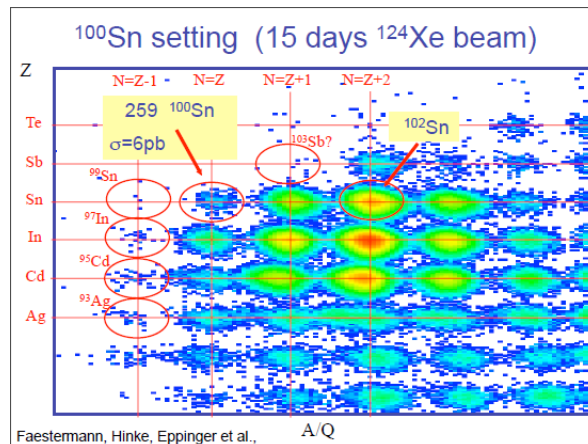
Review Of  $\text{Log}f\tau$  Values  
In  $\beta$  Decay\*

*Nuclear Data Sheets* 84, 487 (1998)  
Article No. DS980015

# Superaligned Gamow–Teller decay of the doubly magic nucleus $^{100}\text{Sn}$

$^{100}\text{Sn}$   $N=Z=50$   
Heaviest  $N=Z$  nucleus

Relativistic Fragmentation reaction  
 $^{124}\text{Xe}$  beam on  $^9\text{Be}$  target



$\sim 260$   $^{100}\text{Sn}$  nuclei produced (0.75/h)  
 $\sim 126$  fully reconstructed decay chains

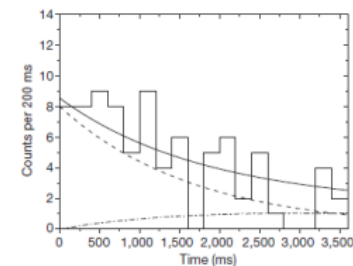
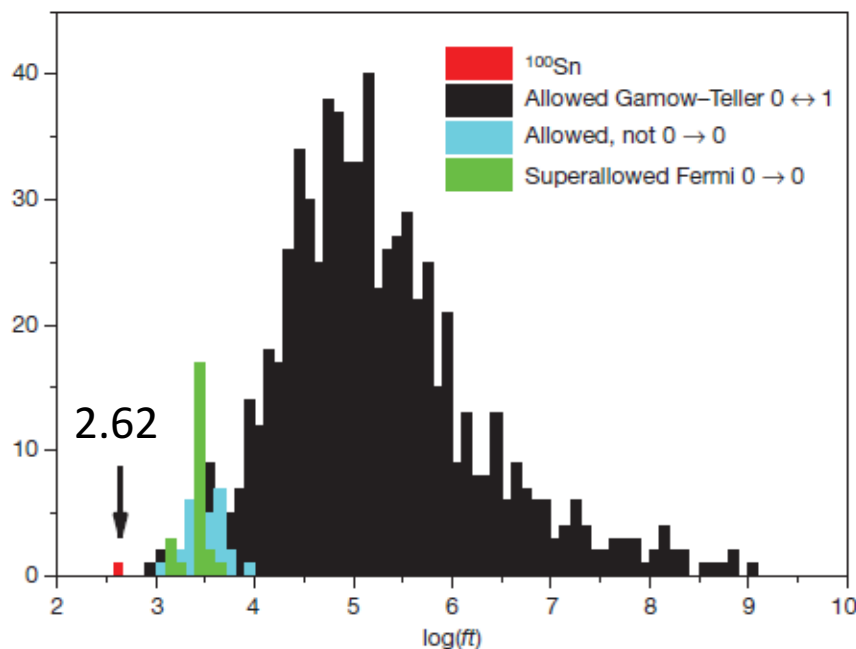


Figure 2 | Time distribution of first decay events. The histogram shows the observed time distribution of all first decay events in the nearest-neighbouring pixels after implantation of  $^{100}\text{Sn}$  nuclei. Decay curves resulting from the MLH analysis are shown individually for  $^{100}\text{Sn}$  (dashed) and its daughter nucleus  $^{100}\text{In}$  (dash-dot). The solid line shows the sum of these decay curves and takes into account a small amount of random background.

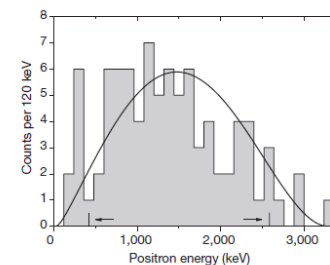


Figure 5 | Distribution of the positron energies emitted in the  $\beta$ -decay of  $^{100}\text{Sn}$ . The spectrum contains only decay events that can be assigned to  $^{100}\text{Sn}$  decays with a probability of at least 75%. The MLH fit was applied to the region between 400 and 2,600 keV, which is indicated with markers. The solid curve illustrates the shape of the best-fitting single-component  $\beta$ -decay phase space function determined by MLH analysis.

## $\beta$ Decay of $^{138}\text{La}$

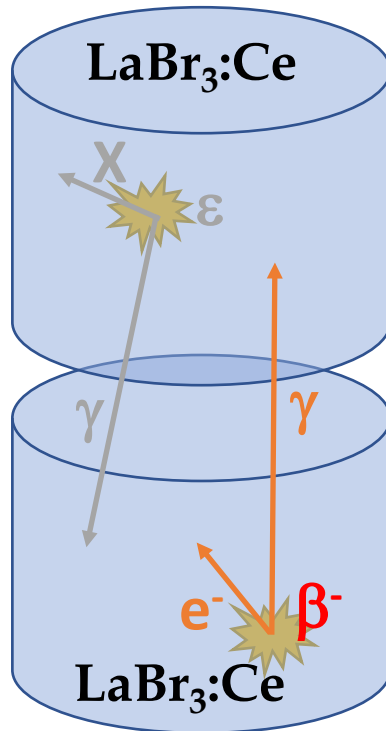
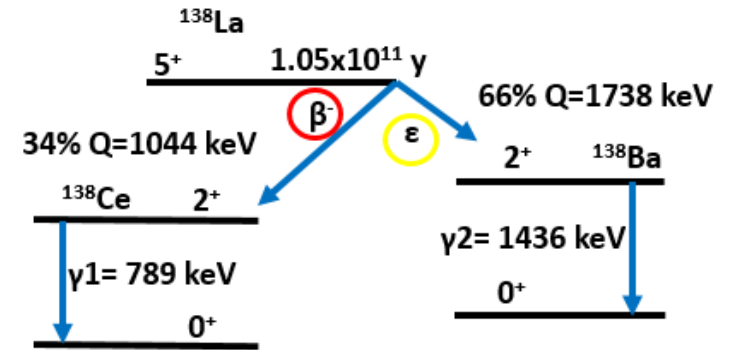
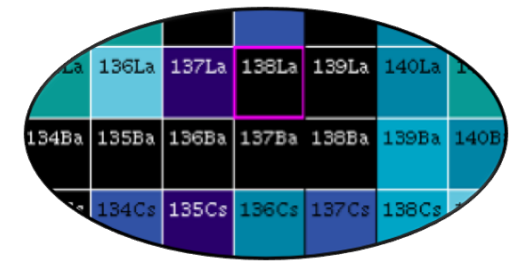
Lanthanum has two isotopes

- $^{139}\text{La}$ , stable
- $^{138}\text{La}$ , **radioactive**:
  - $T_{1/2} = 1.05 \times 10^{11}$  years
  - abundance 0.09%.

$^{138}\text{La}$  has two decay branches

- $\beta^-$  towards  $^{138}\text{Ce}$ , BR= 33.6 %
- $\text{EC}$  towards  $^{138}\text{Ba}$  BR=66 %

Both decays are **second forbidden decays**  
 $5^+ \rightarrow 2^+$

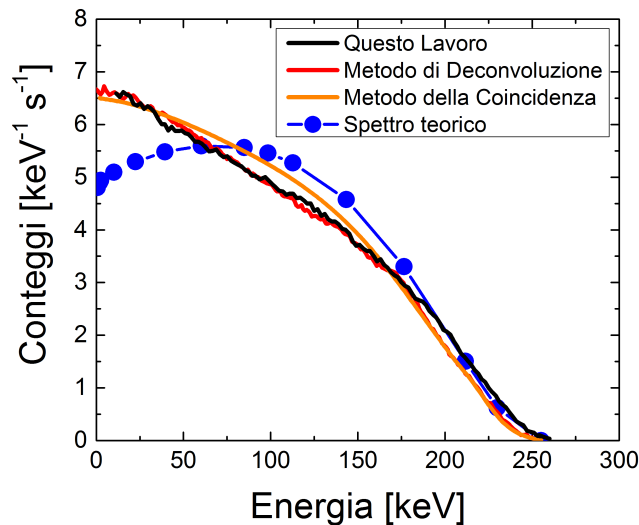
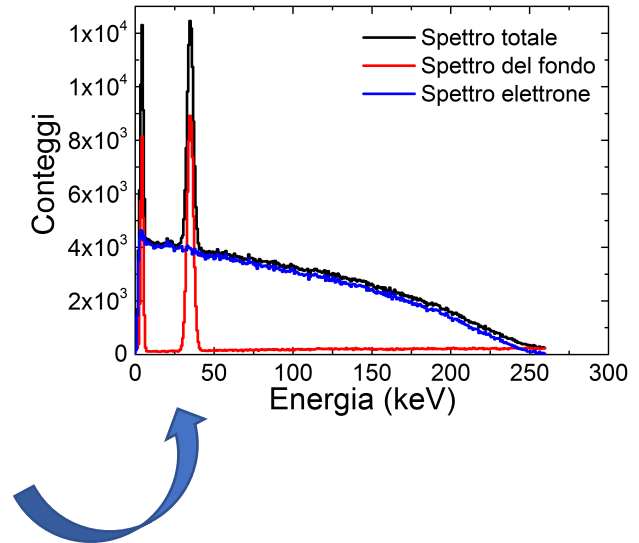
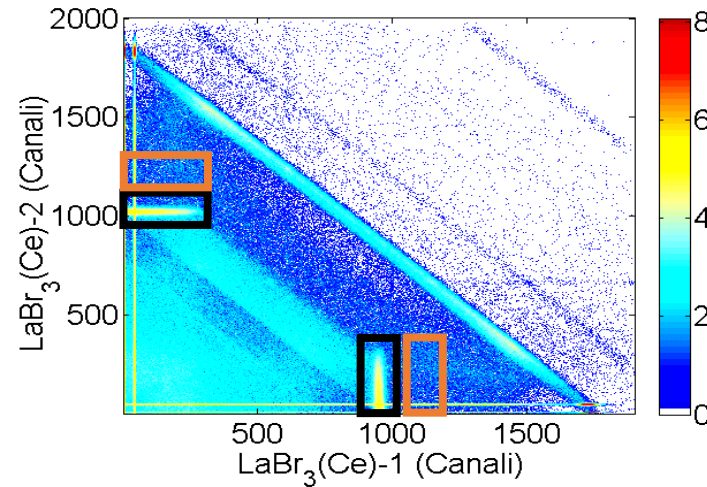


Experiment: measuring the BR and the electron spectra  
 Using  $\text{LaBr}_3(\text{Ce})$  detectors as **source AND as detectors**

- $\beta^-$  decay:  $e^-$  in the det. which  $\beta$  decays, 789 keV  $\gamma$  ray in the other;
- EC: X ray in the det. which decays, 1436 keV  $\gamma$  ray in the other

## Background subtraction to get $e^-$ spectrum:

- Coincidence with 789 keV (black box) **total spectrum**
- **Compton bg** from 1436 keV (red box)
- Bg subtraction cancel peaks at **37.44** and **5.6 keV** in coincidence with 1436 keV.

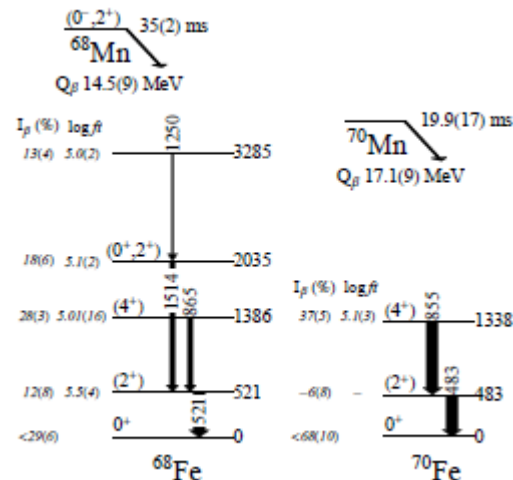
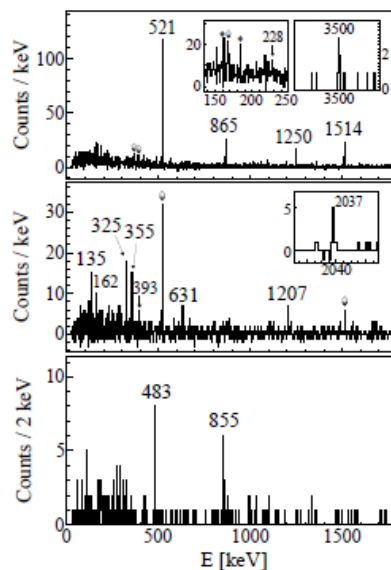
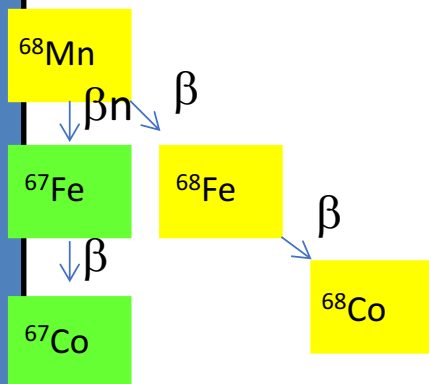
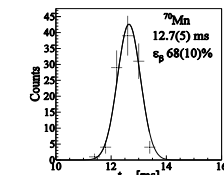
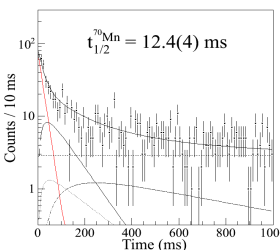
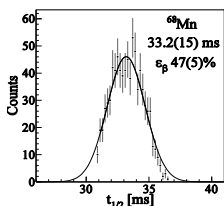
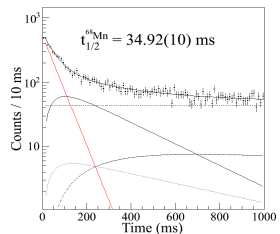
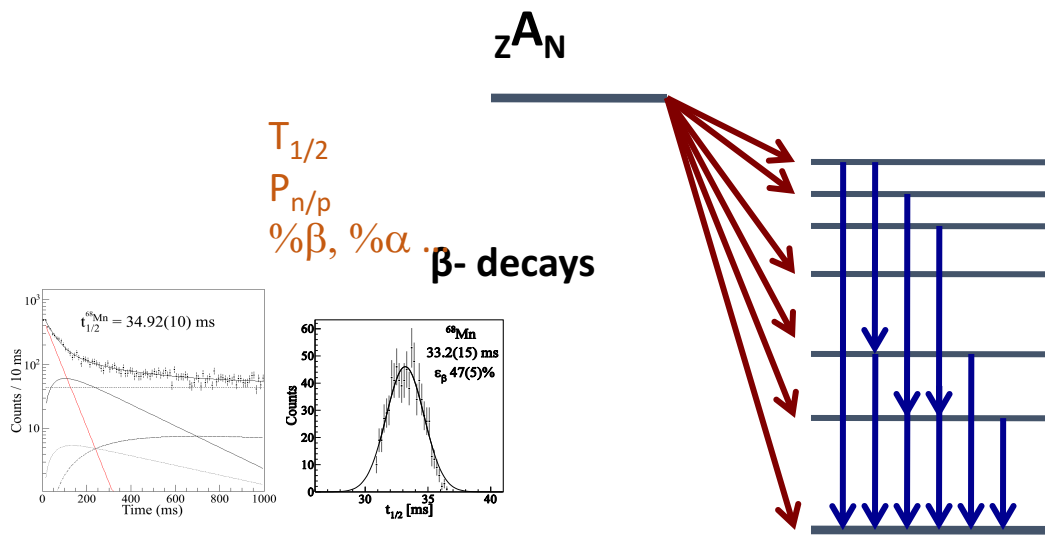


Theory has hard time reproducing low-energy shape of  $\beta$ -electron spectrum

Interaction btw Coulomb field of the nucleus and electron

# General Properties of $\beta$ decay

Quantities that can be extracted in a  $\beta$  decay experiment



Selection rules allow to suggest/define  $I^\pi$  in both mother and daughter

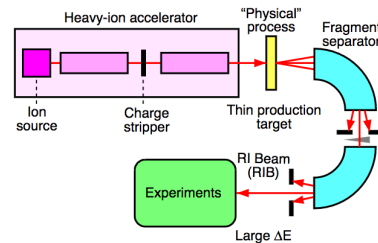
"pandemonium effect" – exotic nuclei –  $\log ft$  is a just upper limit



# Introduction

## Producing radioactive beams i.e. exotic nuclei

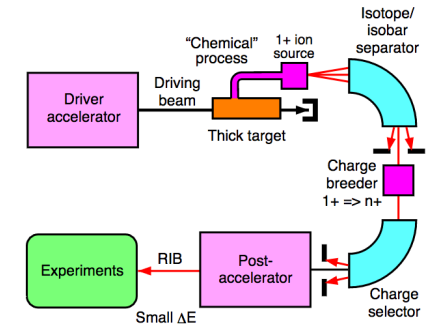
### IN-FLIGHT



**relativistic fragmentation/fission** of heavy nuclei on thin targets

- $> 50 \text{ MeV/u} \rightarrow$  production of cocktail beams of many nuclei
  - Use of spectrometers to transport/separate nuclei of interest  $\rightarrow$  Relatively long decay paths  $\Delta t > 150\text{-}300 \text{ ns}$
  - Nuclei are brought to rest in final focal plane and let decay
- + cocktail beam: many nuclei at once  
 + both short and long-living species  
 + get information already with few ions
- Low cross sections
  - Limitation on rate to distinguish contribution from each species

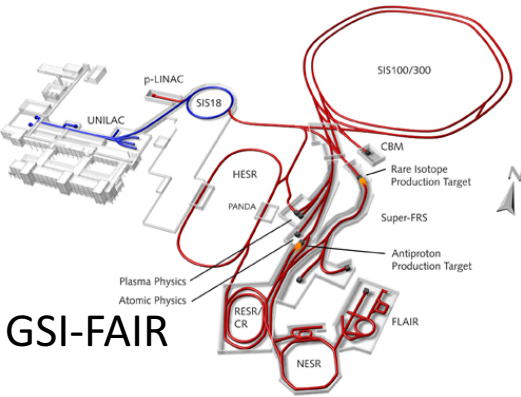
### ISOL method



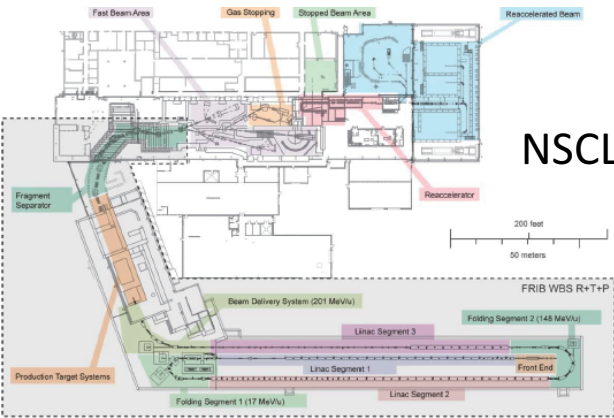
**spallation/fission/fragmentation** on thick targets, followed by chemical/physical processes to extract desired nuclei

- beams produced at very low energies ( $\sim 60 \text{ keV}$ )
  - Mono-isotopic beams sometimes achieved. Impurities due to few contaminant species  $\rightarrow$  usually long-living though
- + high cross section  
 + no need to re-accelerate beams  
 + high rates accepted
- short-living species might not be accessed easily
  - Refractory elements
  - Presence of long-living impurities (isobaric contamination)

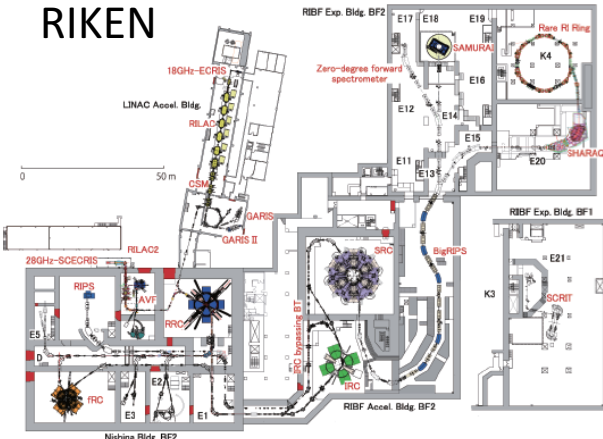
Examples of facilities all over the world...in-flight and isol



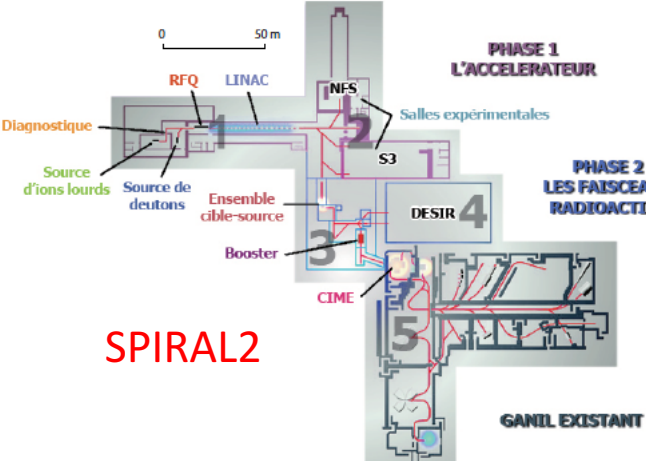
GS-FAIR



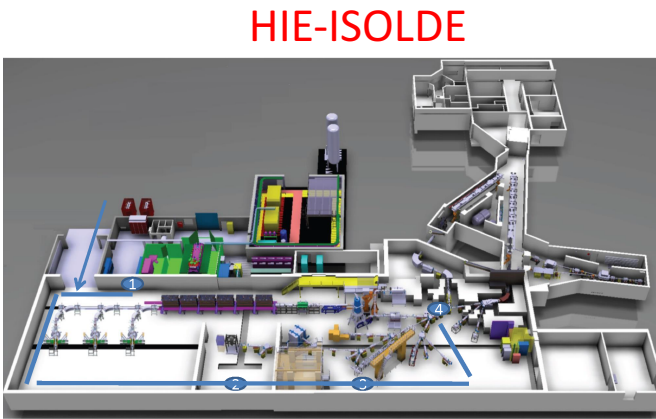
NSCL-FRIB



RIKEN

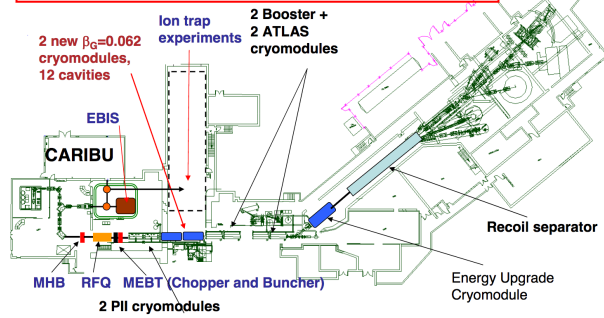


SPIRAL2



HIE-ISOLDE

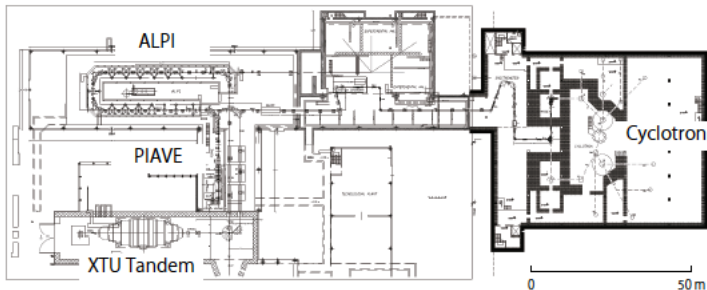
Increase intensity of stable beams from ATLAS by factor >10  
Increase intensity of CARIBU reaccelerated beams by > 5-10  
Increase intensity of in-flight radioactive beams by >100



TRIUMF



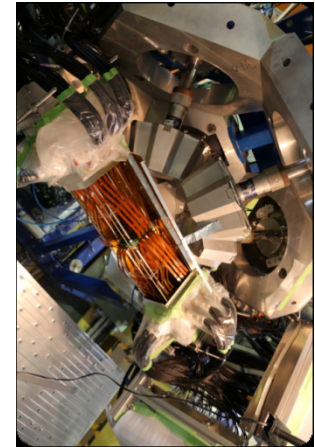
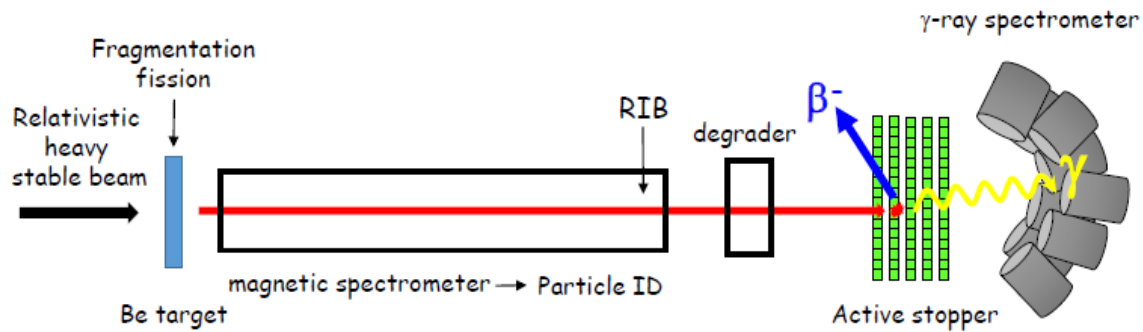
ALTO



S  
P  
S

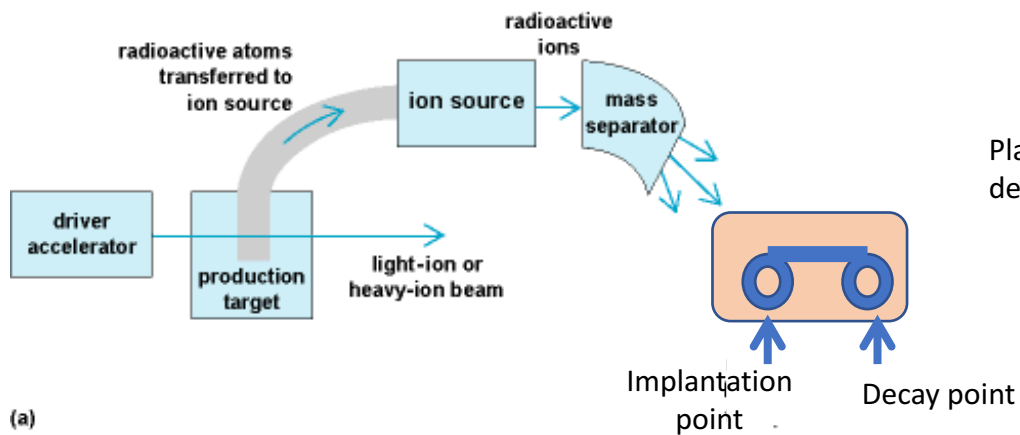
# General Properties of $\beta$ decay

## Measuring $\beta$ decay



Eurica@RIKEN

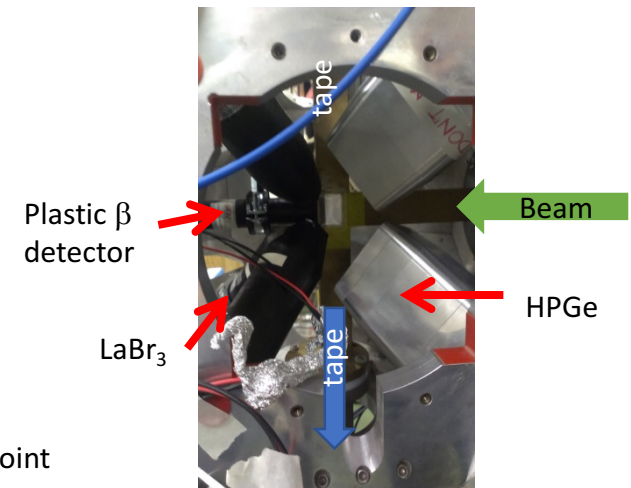
IN-Flight facility



(a)

ISOL facility

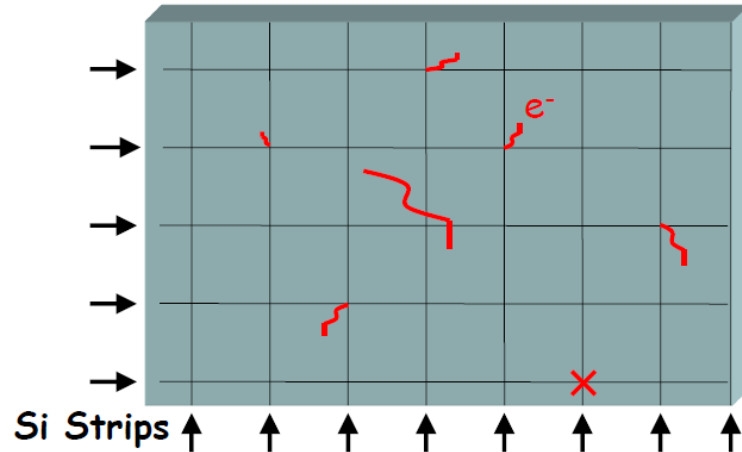
Tape decay stations



IDS @ CERN

# Ion-beta correlation techniques: distinguish implantation and decay within same detector

Focal plane implantation detector sensitive to electron emission



The waiting time between particle implantation and  $\beta$ -particle (or i.c. electron) emission is a measure of the decay half-life. Gamma rays emitted following these decays are detected by the RISING array.

Implantation-decay correlations with large background  
(half lives similar to the implantation rate):

- ✓ implant-decay time correlation: active catcher
- ✓ implant-decay position correlation: granularity
- ✓ implant of several ions: thickness and area
- ✓ energy of the implanted ion and the emitted  $\beta$

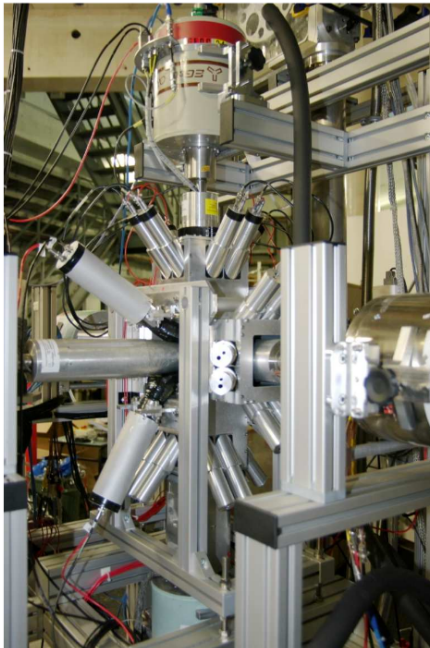
GSI

- \* Dual gain pre-amps on DSSSD to get energies of implanted ion and b-particle
- \* All events time stamped with MHz clock.

RIKEN

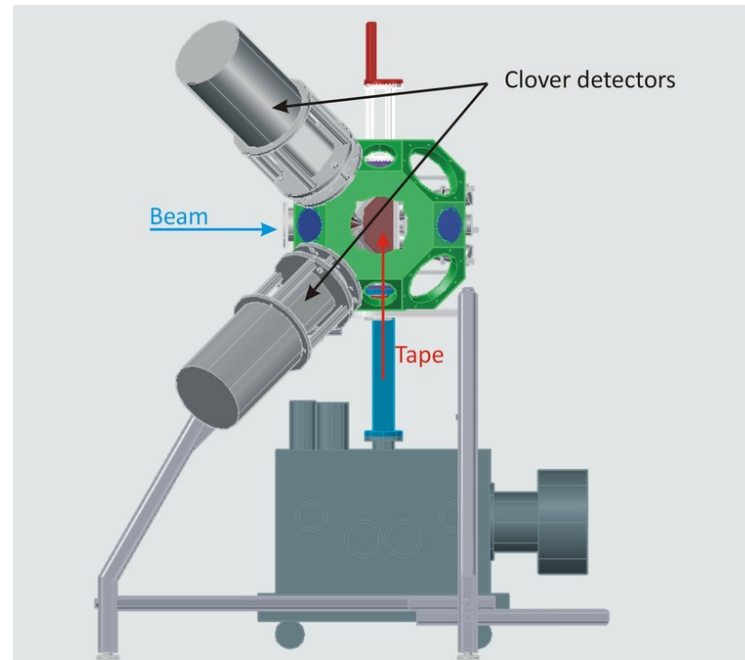
- \* Low gain branch for implanted ions
- \* High-gain branch for  $\beta$  and  $\alpha$  decays

# TAPE Station systems



BEDO@ALTO

## Principle of operation

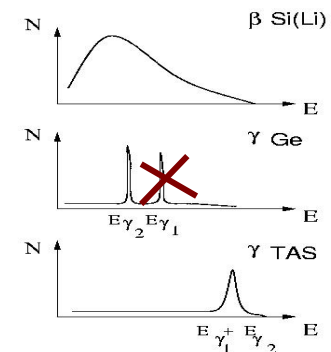
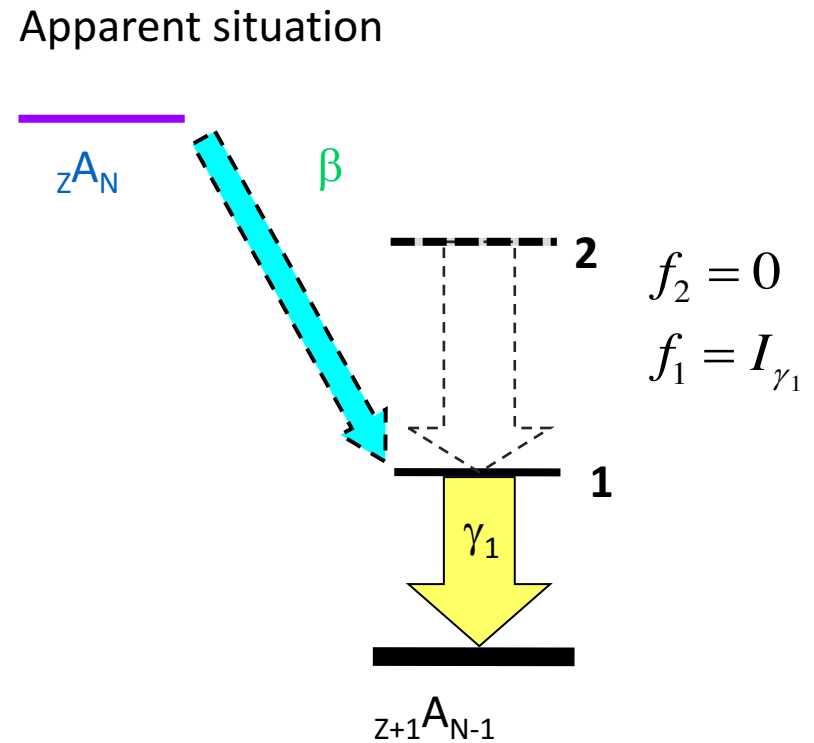
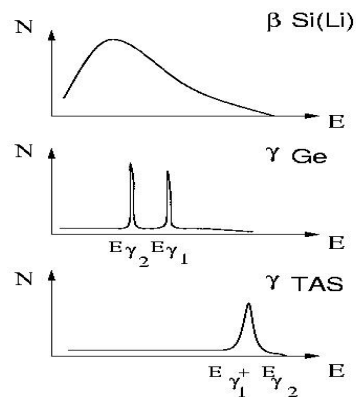
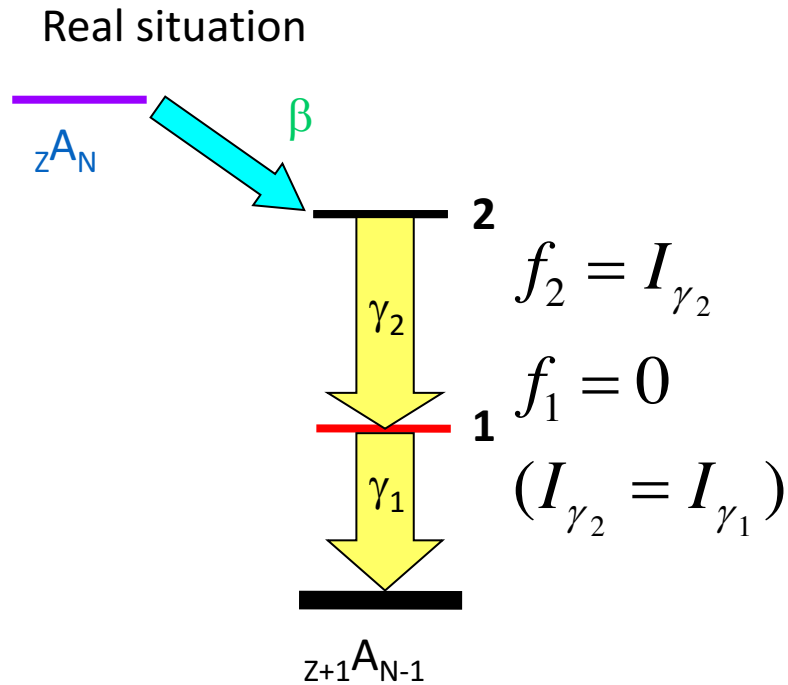


IDS@ISOLDE

- Implantation and decay in the same measuring point
- Equipped with egs. Ge detectors and Plastic (or Si detectors) for  $\beta$  particles
- Trigger given by proton arrival and  $\beta$  signal
- Long-living activity is removed by moving away the tape
- NB: no direct signal of implantation

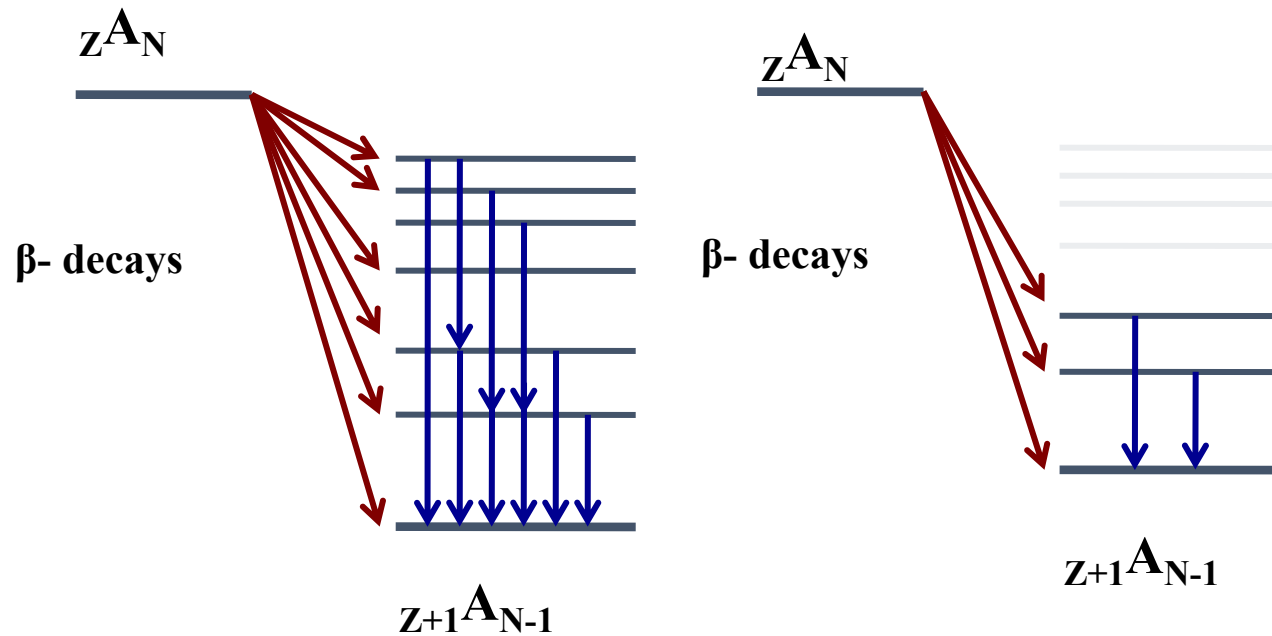
# General Properties of $\beta$ decay

## Pandemonium effect



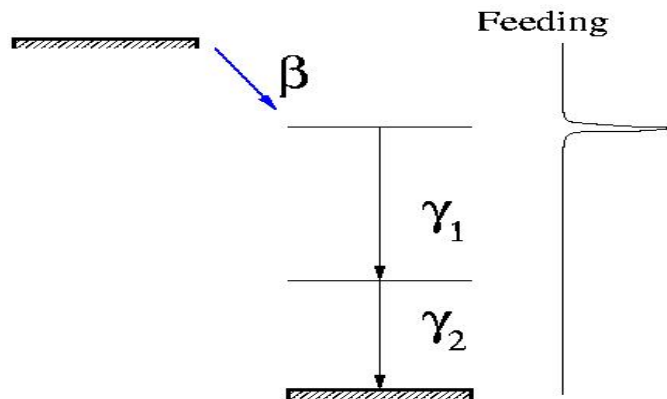


# General Properties of $\beta$ decay



- HPGe detectors are conventionally used to construct the level scheme populated in the decay
  - ➔ Higher  $Q_{\text{value}}$  higher possibility of missing feeding owing to high-energy  $\gamma$  ray emission or emission of large number of  $\gamma$  rays
  - From the  $\gamma$  intensity balance we deduce the  $\beta$ -feeding
- Pandemonium effect implies
- ➔ Wrong definition of gamma feeding and branching ratios  $I_\beta$  and  $\log ft$

# Total Absorption Gamma Spectrometer measurements



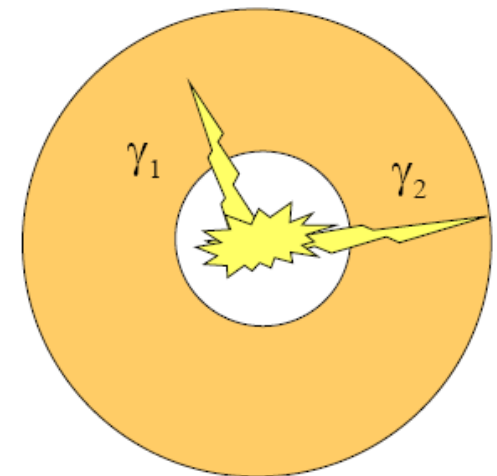
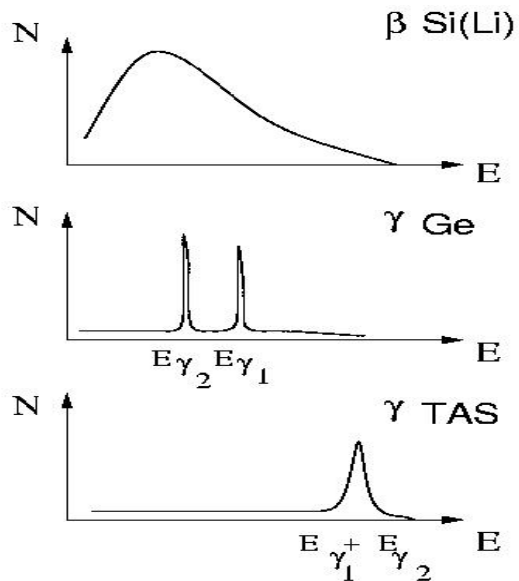
Since the gamma detection is the only reasonable way to solve the problem, we need a highly efficient device:

## A TOTAL ABSORTION SPECTROMETER

Instead of detecting the individual gamma rays we sum the energy deposited by the gamma cascades in the detector.

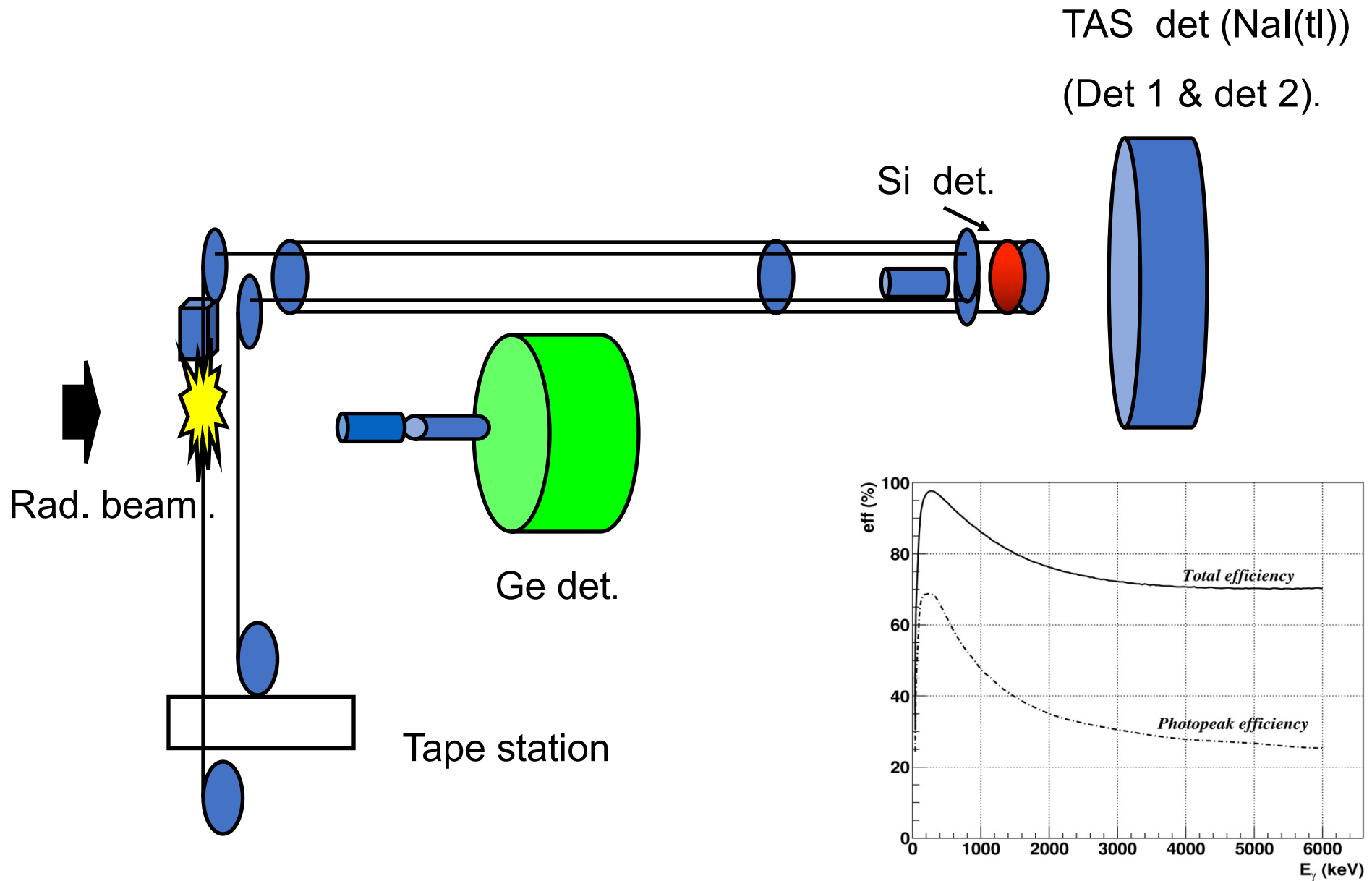
A TAS is a **calorimeter**!

Big crystal,  $4\pi$  (BaF<sub>2</sub>/NaI/HPGe)

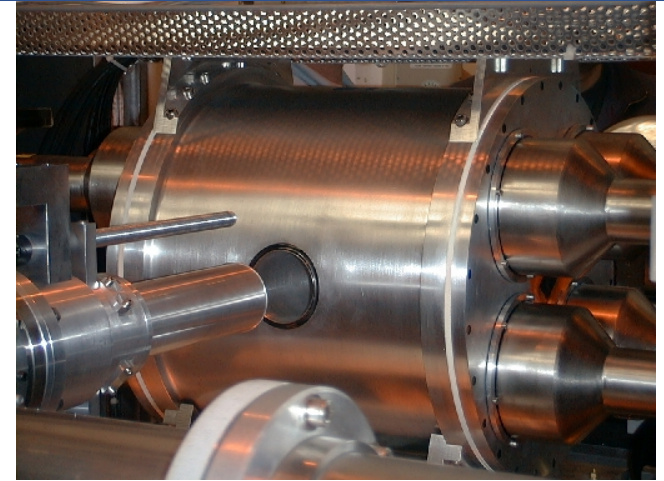
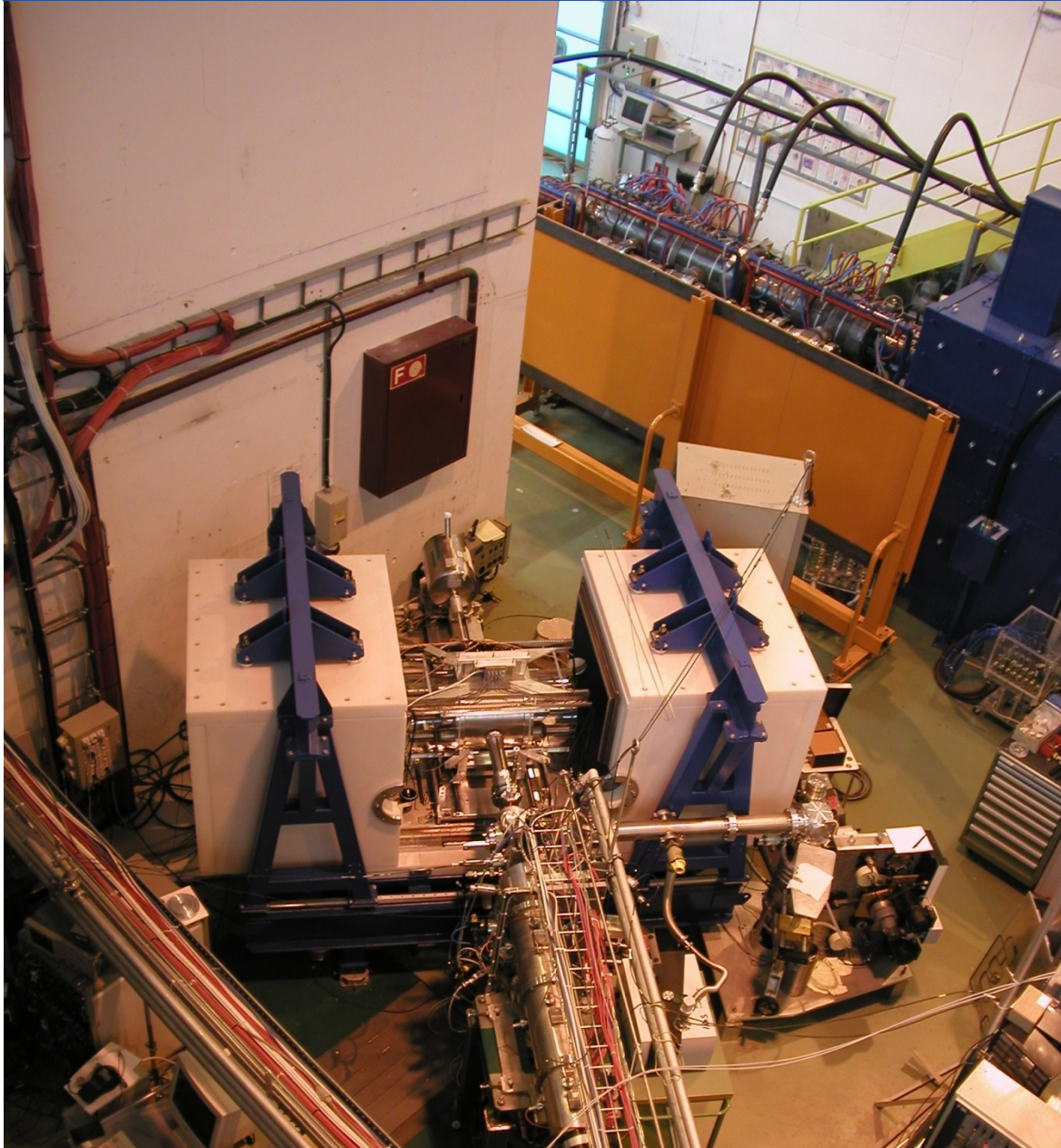




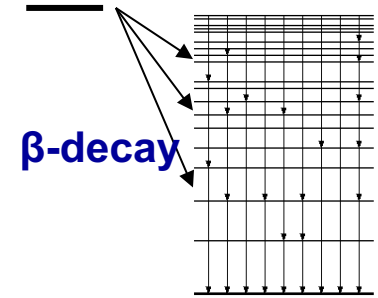
# TAS experimental setup



# Lucrecia: the TAS at ISOLDE (CERN) (Madrid-Strasbourg-Surrey-Valencia)



- A large NaI cylindrical crystal 38 cm Ø, 38cm length
- An X-ray detector (Ge)
- A  $\beta$  detector
- Possibility of collection point inside the crystal



$$d_i = \sum_j R_{ij} f_j \quad \text{or} \quad \mathbf{d} = \mathbf{R} \cdot \mathbf{f}$$

$\mathbf{R}$  is the response function of the spectrometer,  $R_{ij}$  means the probability that feeding at a level  $j$  gives counts in data channel  $i$  of the spectrum

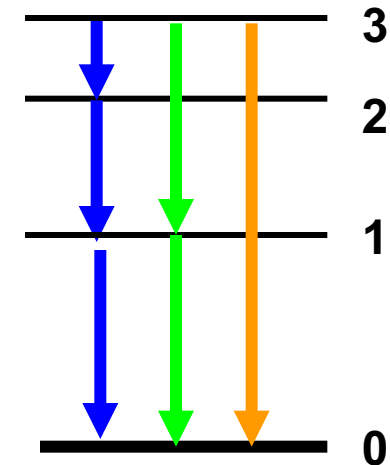
The response matrix  $\mathbf{R}$  can be constructed by recursive convolution:

$$\mathbf{R}_j = \sum_{k=0}^{j-1} b_{jk} \mathbf{g}_{jk} \otimes \mathbf{R}_k$$

$\mathbf{g}_{jk}$ :  $\gamma$ -response for  $j \rightarrow k$  transition

$\mathbf{R}_k$ : response for level  $k$

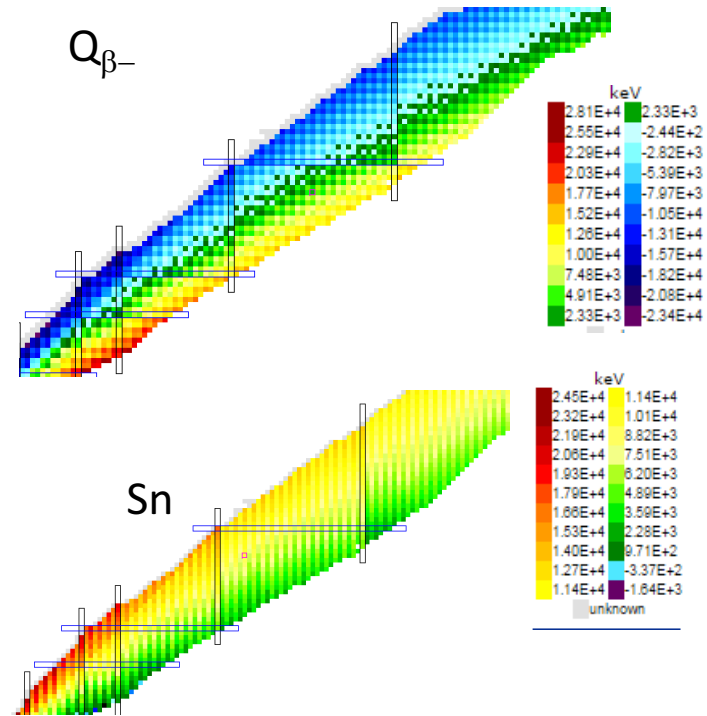
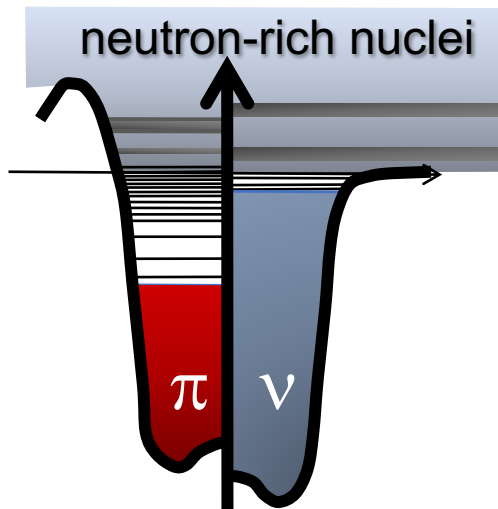
$b_{jk}$ : branching ratio for  $j \rightarrow k$  transition



# $\beta$ Decay in Exotic Nuclei

Going away from stability

- Shortening of half-lives
- Increasing  $Q_{\text{value}}$
- Decreasing  $S_n$  (on n-rich side)
- Higher possibilities of competing mechanisms
- GT unfavoured owing to large mismatch in wave functions btw mother/daughter
- Increasing role of forbidden decays



- Pandemonium effect

