

\hbar



Out-of-equilibrium density dynamics of a quenched fermionic system

S. Porta¹, F. M. Gambetta^{1,2}, F. Cavaliere^{1,2},
N. Traverso Ziani³, M. Sasseti^{1,2}

¹*Dipartimento di Fisica, Università degli Studi di Genova*

²*SPIN-CNR Genova*

³*Institute for Theoretical Physics and Astrophysics, University of Würzburg*

Reference: [Phys. Rev. B 94, 085122 \(2016\)](#)

— Padova, 102° Congresso Nazionale SIF —
26-30 Settembre 2016

Introduction

thermalization &
spreading of information

coherent quantum
technologies

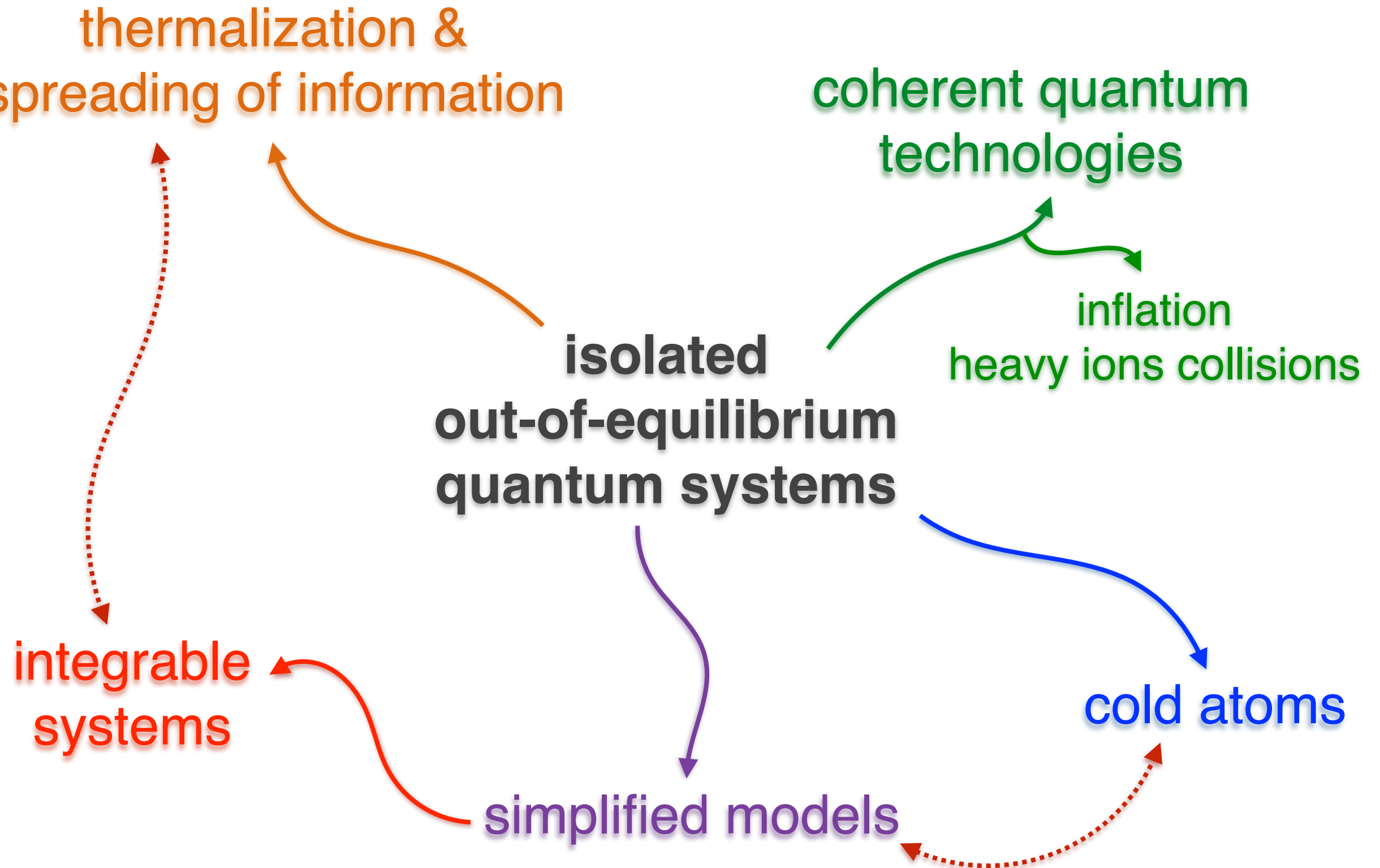
inflation
heavy ions collisions

**isolated
out-of-equilibrium
quantum systems**

integrable
systems

cold atoms

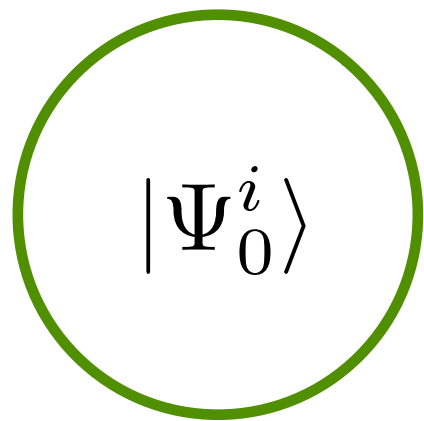
simplified models



Quantum quenches


Definition: a **change in time** of the **parameter(s)** that governs the dynamics of an **isolated** quantum system.

- ▶ sudden, adiabatic, or more general: linear ramps,...
- ▶ **parameter(s)**: coupling constants, external fields, confinement,...
- ▶ **isolated** quantum system: no coupling with the environment → unitary time-evolution

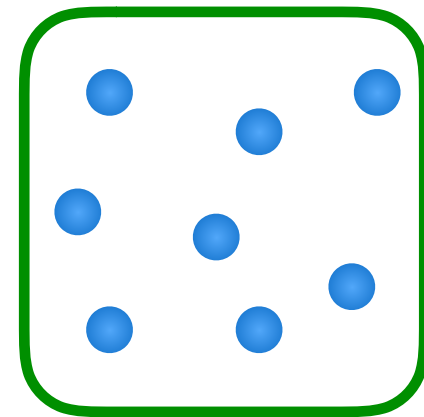


$$H_i = H(\{\mathbf{g}_i\})$$

Quench!



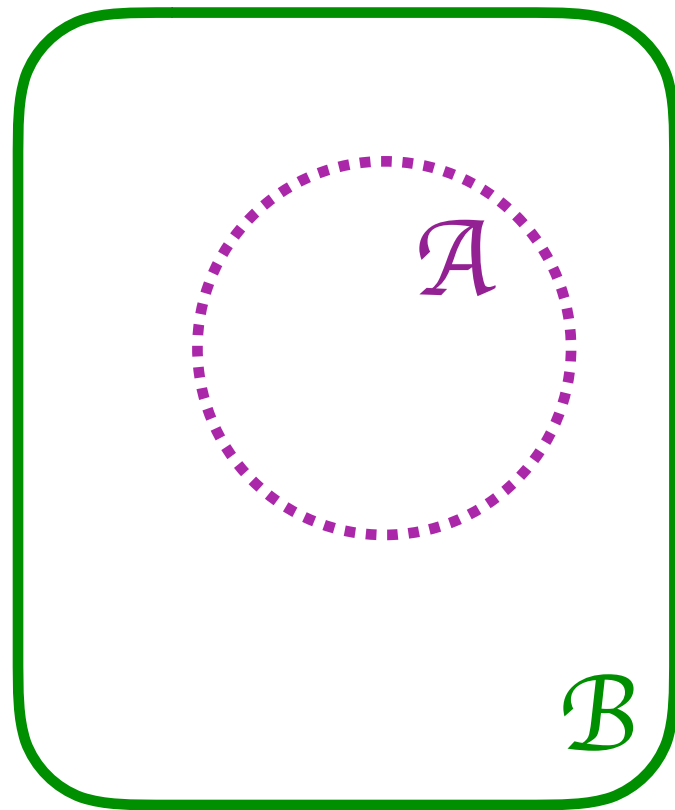
$$H(t) = H(\{\mathbf{g}(t)\})$$



$$H_f = H(\{\mathbf{g}_f\})$$

Thermalization in isolated quantum systems

Does an isolated quantum system thermalize?



- ▶ **thermal** density matrix $\rho^{\text{th}} = \frac{e^{-\beta H}}{Z}$, with β fixed by $\langle \Psi_0^i | H_f | \Psi_0^i \rangle = \text{Tr}[\rho^{\text{th}} H_f]$
- ▶ reduced density matrix $\rho_A(\infty) = \lim_{t \rightarrow \infty} \text{Tr}_B \rho(t)$
- ▶ thermalization if $\rho_A(\infty) = \rho_A^{\text{th}}$

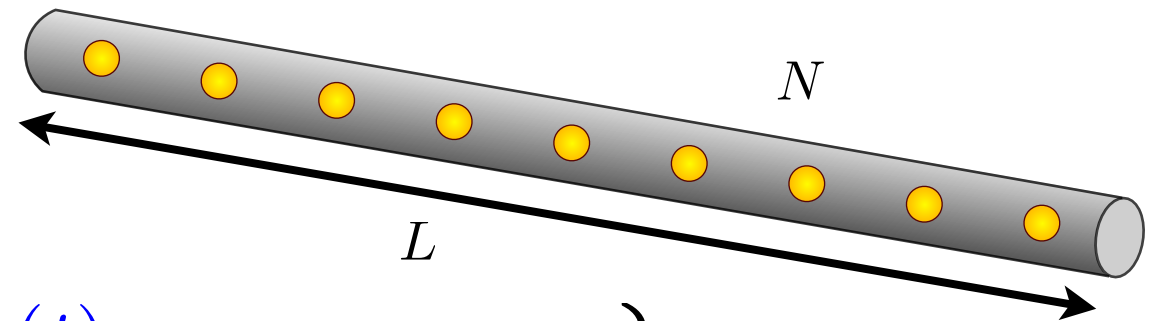
Polkovnikov et al., RMP (2011)

What happens in integrable systems?

- ▶ Generalized Gibbs Ensemble (**GGE**) $\rho^{\text{GGE}} = \frac{e^{-\sum_m \lambda_m I_m}}{Z^{\text{GGE}}}$, with λ_m fixed by $\langle \Psi^i(t) | I_m | \Psi^i(t) \rangle = \langle \Psi_0^i | I_m | \Psi_0^i \rangle = \text{Tr}[\rho^{\text{GGE}} I_m]$
- ▶ GGE “thermalization” $\rho_A(\infty) = \rho_A^{\text{GGE}}$

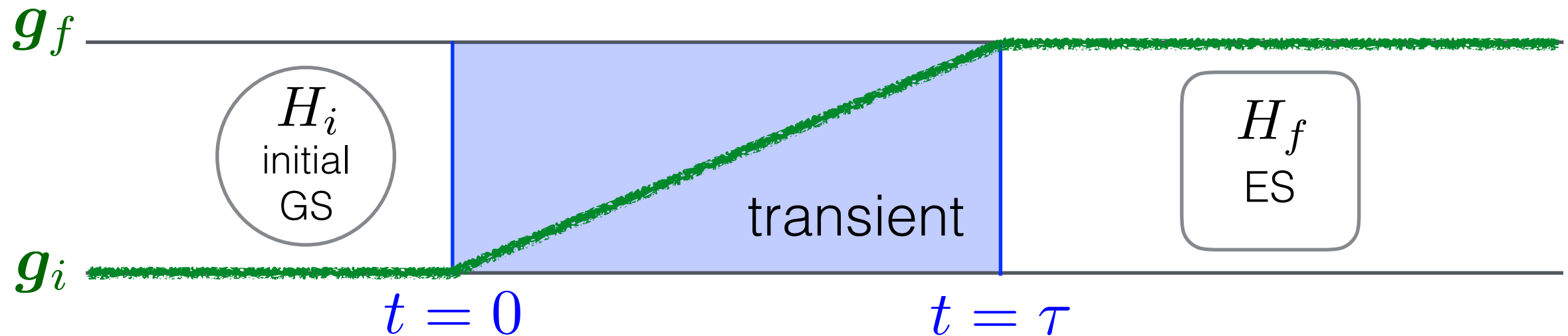
Rigol et al., PRL (2007) 4 / 11

Quench in a open boundaries LL



► (bosonized) Hamiltonian

$$H(t) = \sum_{q>0} q \left\{ [v_F + g_4(t)] b_q^\dagger b_q - \frac{g_2(t)}{2} (b_q b_q + b_q^\dagger b_q^\dagger) \right\} + \mathcal{E}(N)$$



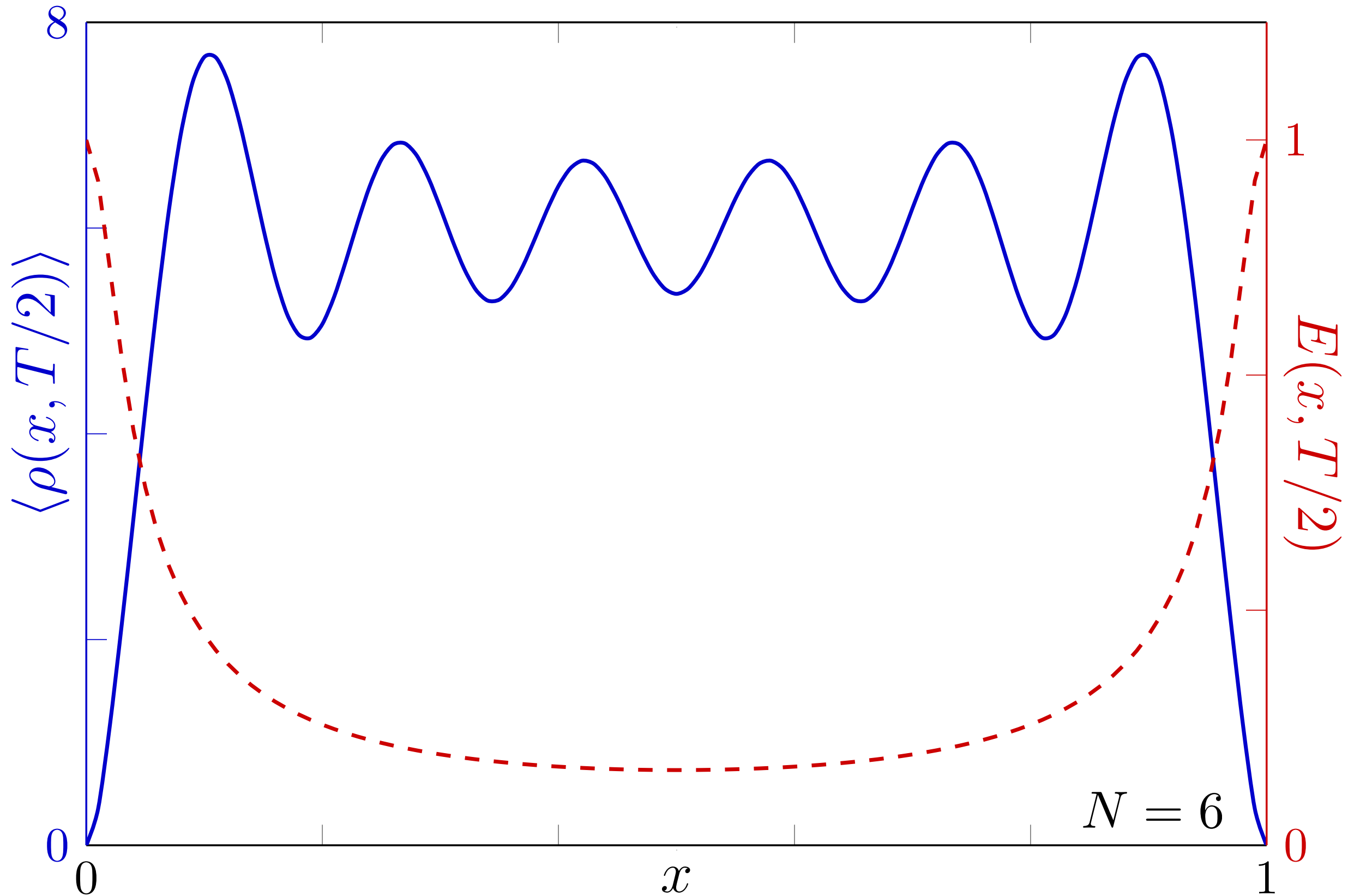
$$|\Psi\rangle = |GS_i\rangle \longrightarrow |\Psi(t)\rangle = U(t, 0)|\Psi\rangle \longrightarrow |\Psi(t)\rangle = e^{-iH_f(t-\tau)}|\Psi(\tau)\rangle$$

Cazalilla, PRL (2006); Dóra et. al, PRL (2011)

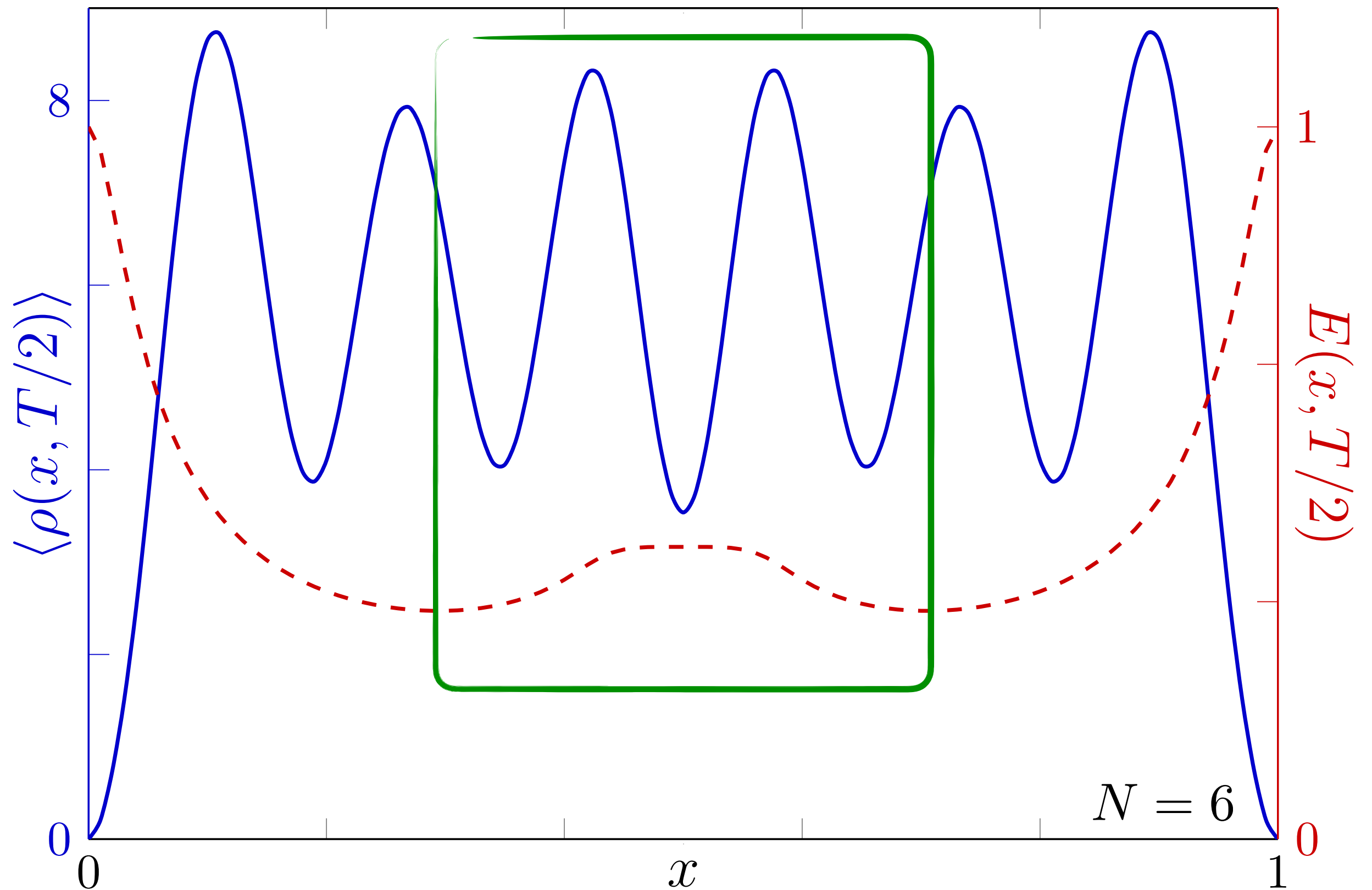
► Average particle density ($\hat{\rho}(x, t) = \hat{\Psi}^\dagger(x, t)\hat{\Psi}(x, t)$)

$$\langle \rho(x, t) \rangle_i = \frac{N}{L} \{1 - E(x, t) \cos [2k_F x - 2f(x)]\}$$

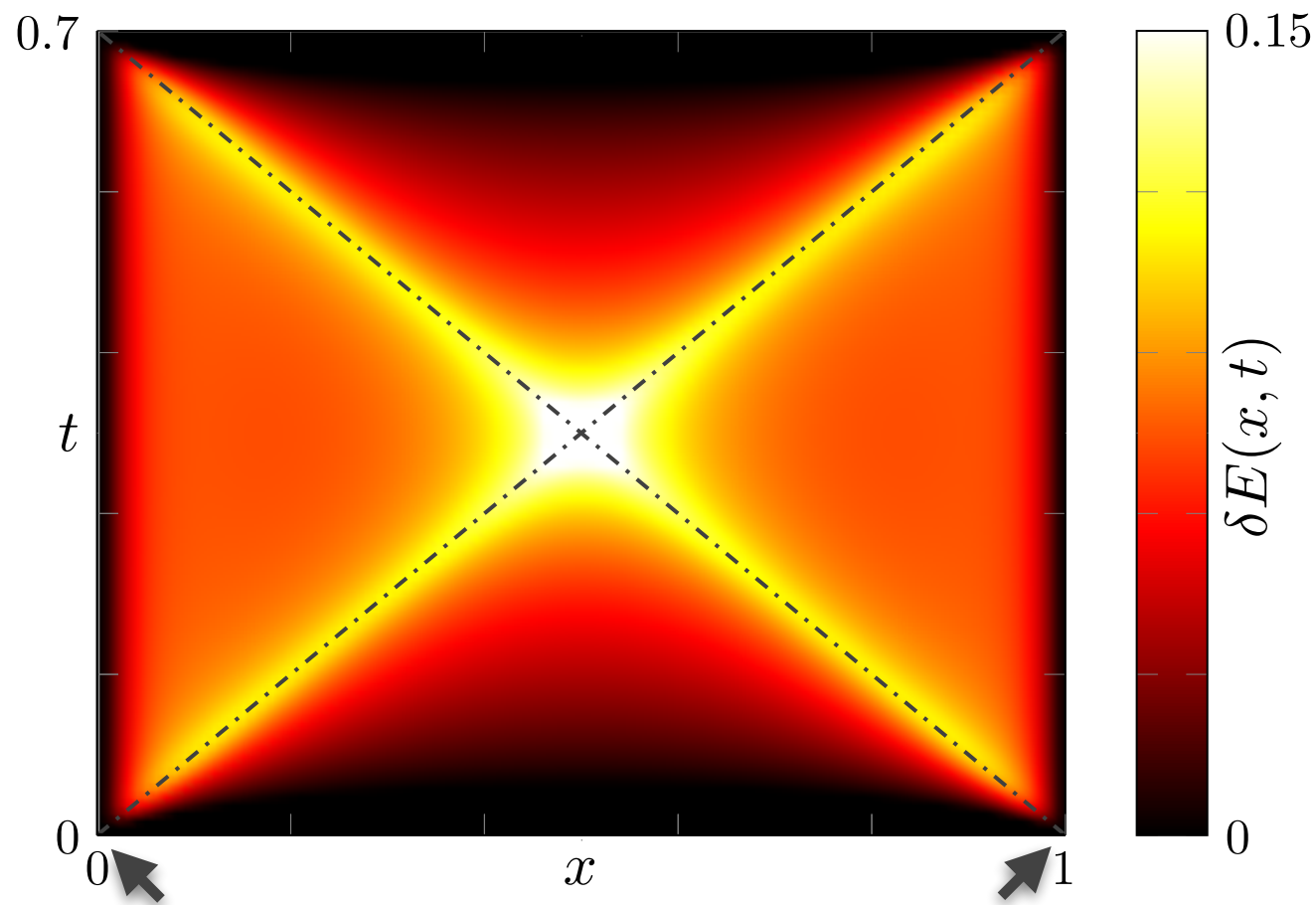
Density in a standard 1D LL with OBC



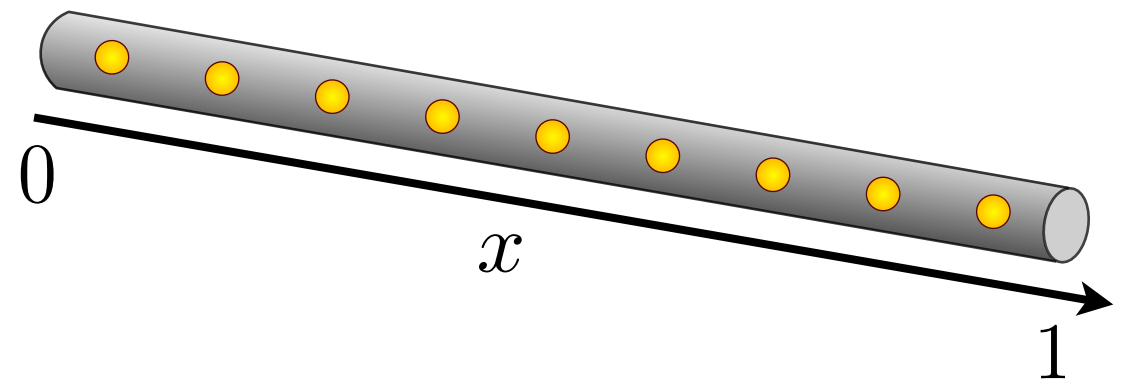
Density in a quenched 1D LL with OBC



Light cones dynamics – sudden quench



► Focus on
 $\delta E(x, t) = E(x, t) - E(x, 0)$
 over one period.



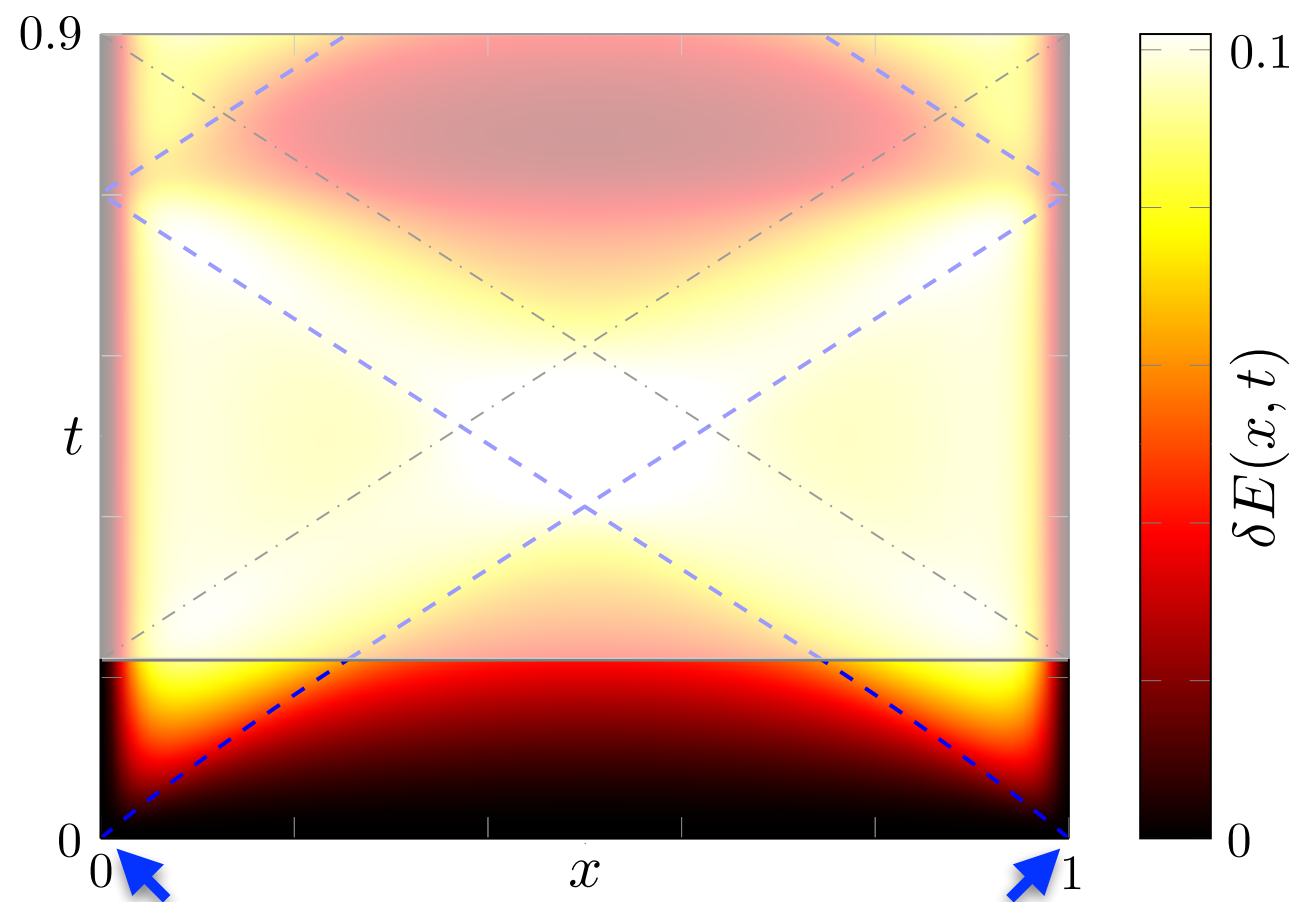
► **Sudden quench:** a **LC** perturbation emerges at the boundaries and moves *ballistically* at velocity v_f

$$H_f = \sum_{q>0} q v_f \beta_{q,f}^\dagger \beta_{q,f}$$

► Analytically $E(x, t) = E_{\text{sq}}^{(\text{GGE})}(x) f_{\text{sq}}(t) \mathcal{C}_{\text{sq}}(x, v_f t)$

$$\mathcal{C}_{\text{sq}}(x, y) = \mathcal{C}_{\text{sq}}^R(x - y) \mathcal{C}_{\text{sq}}^L(x + y)$$

Light cones dynamics – finite duration quench



► Adiabatic quench
 $(\tau \gg \tau_{\text{ad}} \sim L|\eta|/v_i)$

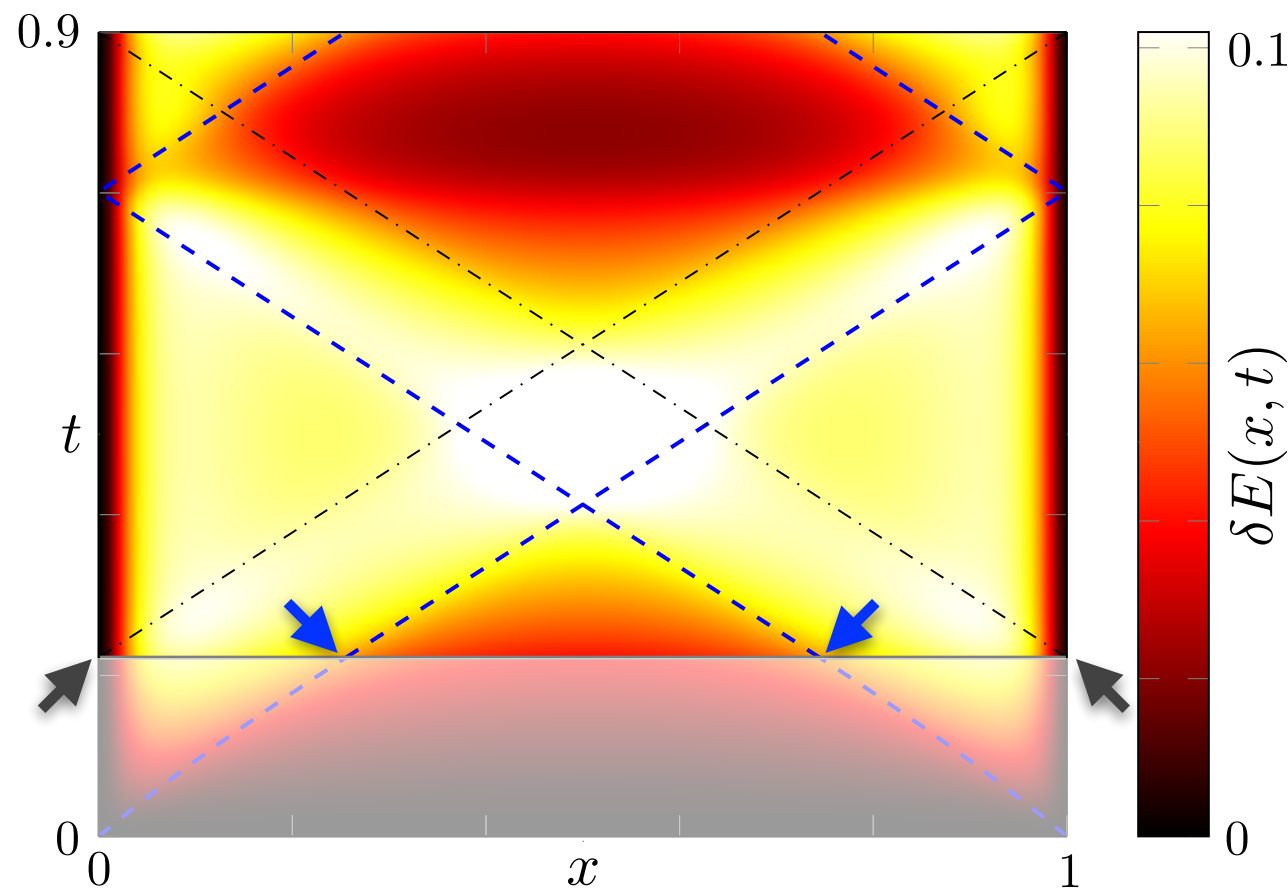
$$E(x, t) = G(x, t) \mathcal{C}_{\text{ad}}(x, \ell(t))$$

$$\mathcal{C}_{\text{ad}}(x, y) = \mathcal{C}_{\text{ad}}^R(x - y) + \mathcal{C}_{\text{ad}}^L(x + y)$$

► On-ramp: **LC1** emerges at the boundaries

$$H(\bar{t}) = \sum_{q>0} q \bar{v}(\bar{t}) \beta_{q,\bar{t}}^\dagger \beta_{q,\bar{t}} \longrightarrow \bar{v}(t) = v_i \sqrt{1 + \frac{\eta t}{\tau}}$$

Light cones dynamics – finite duration quench



► Adiabatic quench
 $(\tau \gg \tau_{\text{ad}} \sim L|\eta|/v_i)$

$$E(x, t) = E_{\text{ad}}^{(\text{GGE})}(x) [f_{\text{ad}}(t) + A_1 \mathcal{C}_{\text{ad}}(x, v_f(t - \tau) + d) - A_2 \mathcal{C}_{\text{ad}}(x, v_f(t - \tau))]$$

$$\mathcal{C}_{\text{ad}}(x, y) = \mathcal{C}_{\text{ad}}^R(x - y) + \mathcal{C}_{\text{ad}}^L(x + y)$$

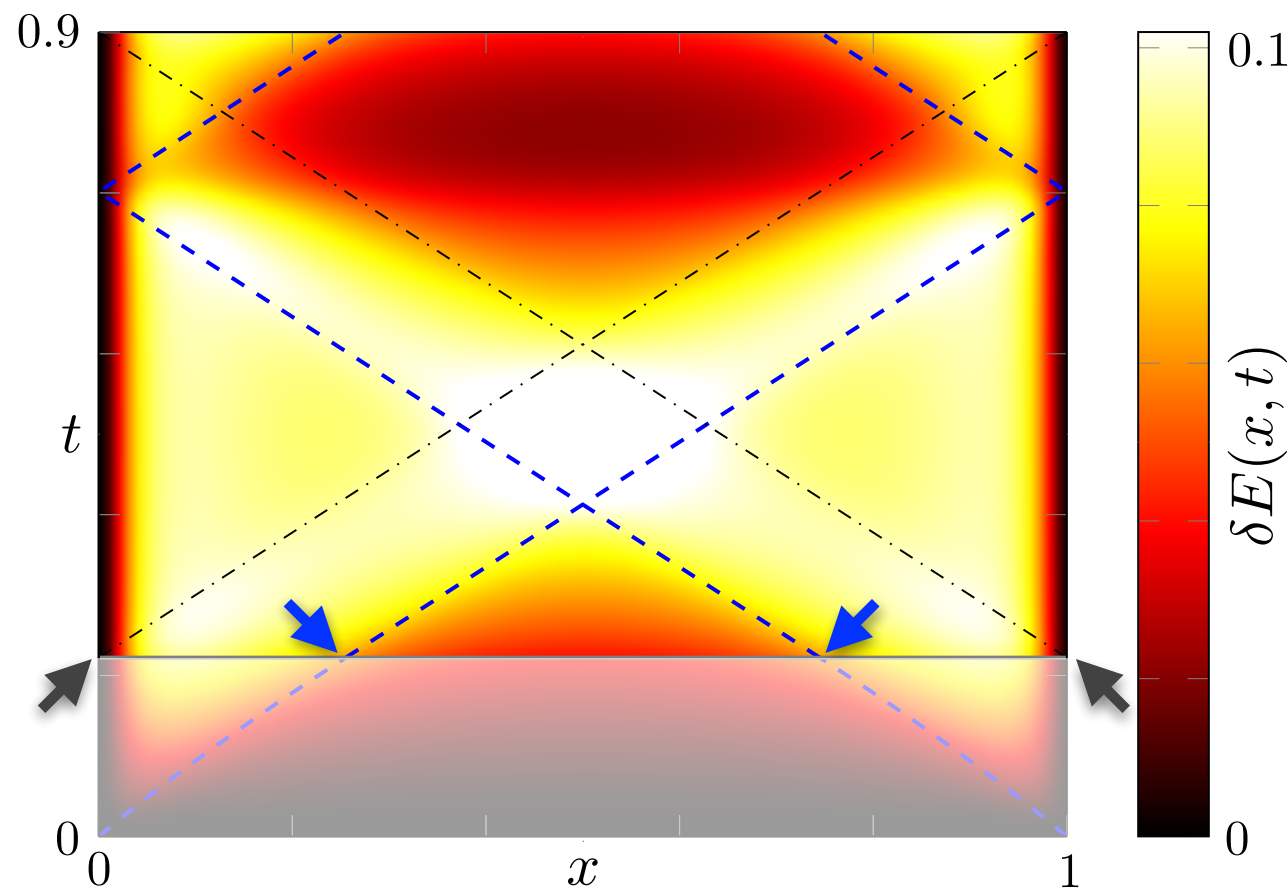
► On-ramp: **LC1** emerges at the boundaries

$$H(\bar{t}) = \sum_{q>0} q \bar{v}(\bar{t}) \beta_{q,\bar{t}}^\dagger \beta_{q,\bar{t}} \longrightarrow \bar{v}(t) = v_i \sqrt{1 + \frac{\eta t}{\tau}}$$

► Post-quench: **LC2** emerges at the boundaries, **LC1** goes on with velocity v_f

$$H_f = H(t \geq \tau) = \sum_{q>0} q v_f \beta_{q,f}^\dagger \beta_{q,f}$$

Light cones dynamics – finite duration quench



► Adiabatic quench
 $(\tau \gg \tau_{\text{ad}} \sim L|\eta|/v_i)$

$$E(x, t) = E_{\text{ad}}^{(\text{GGE})}(x) [f_{\text{ad}}(t) + A_1 \mathcal{C}_{\text{ad}}(x, v_f(t - \tau) + d) - A_2 \mathcal{C}_{\text{ad}}(x, v_f(t - \tau))]$$

$$\mathcal{C}_{\text{ad}}(x, y) = \mathcal{C}_{\text{ad}}^R(x - y) + \mathcal{C}_{\text{ad}}^L(x + y)$$

► On-ramp: **LC1** emerges at the boundaries

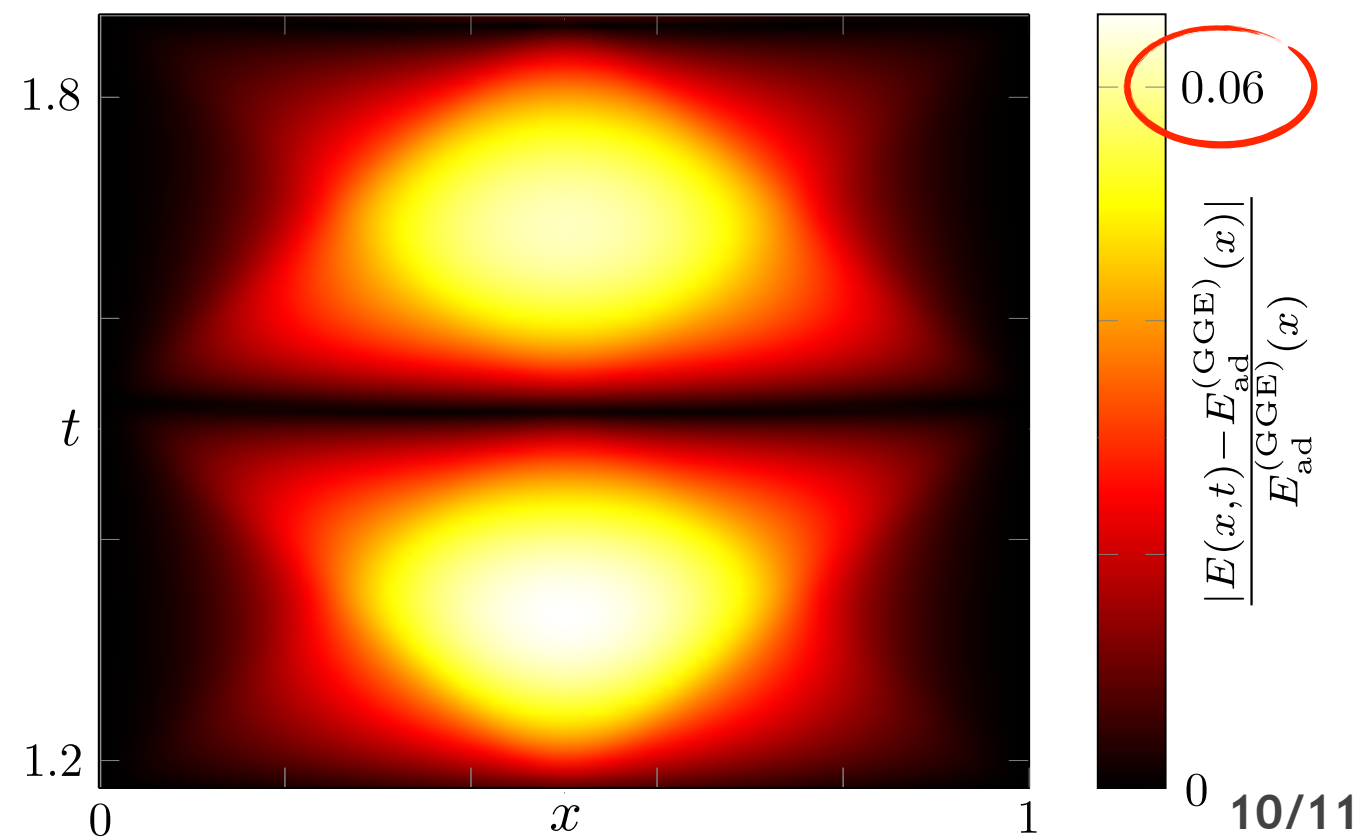
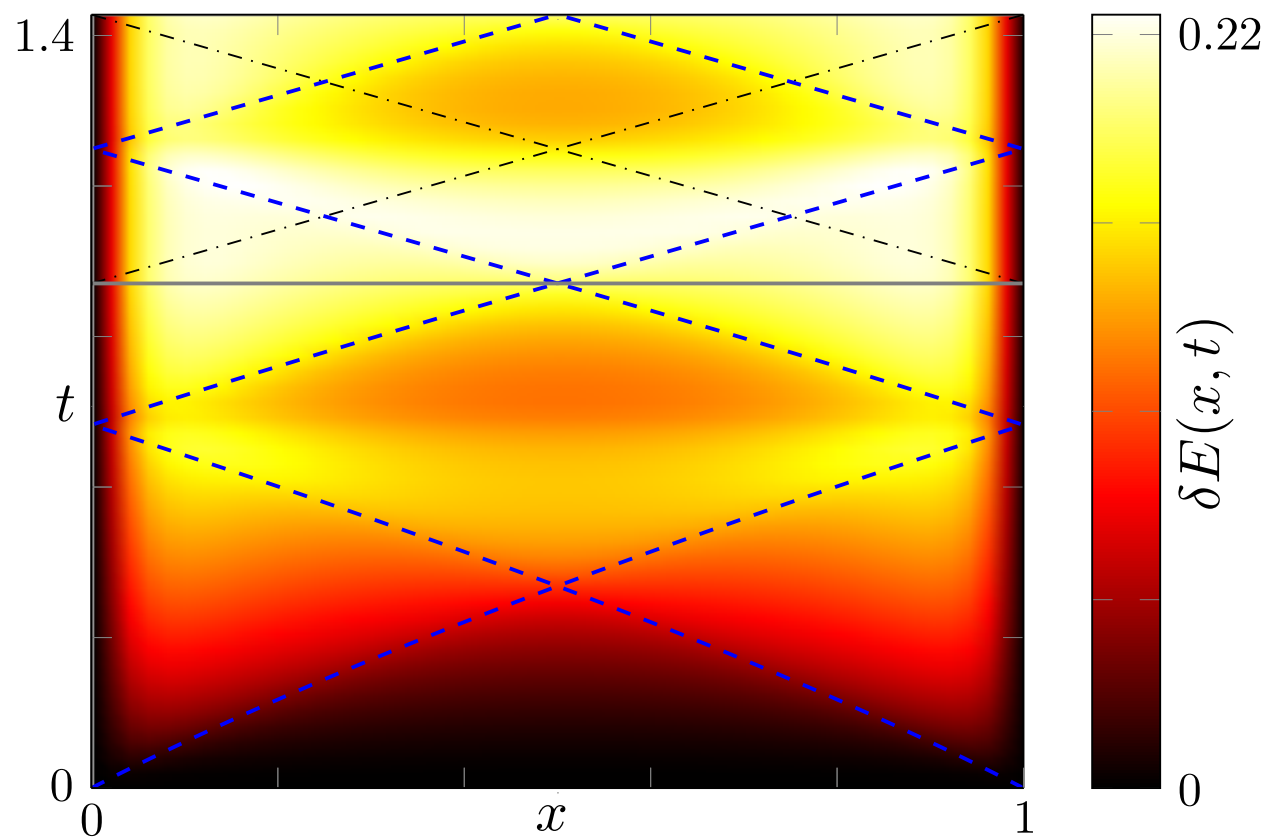
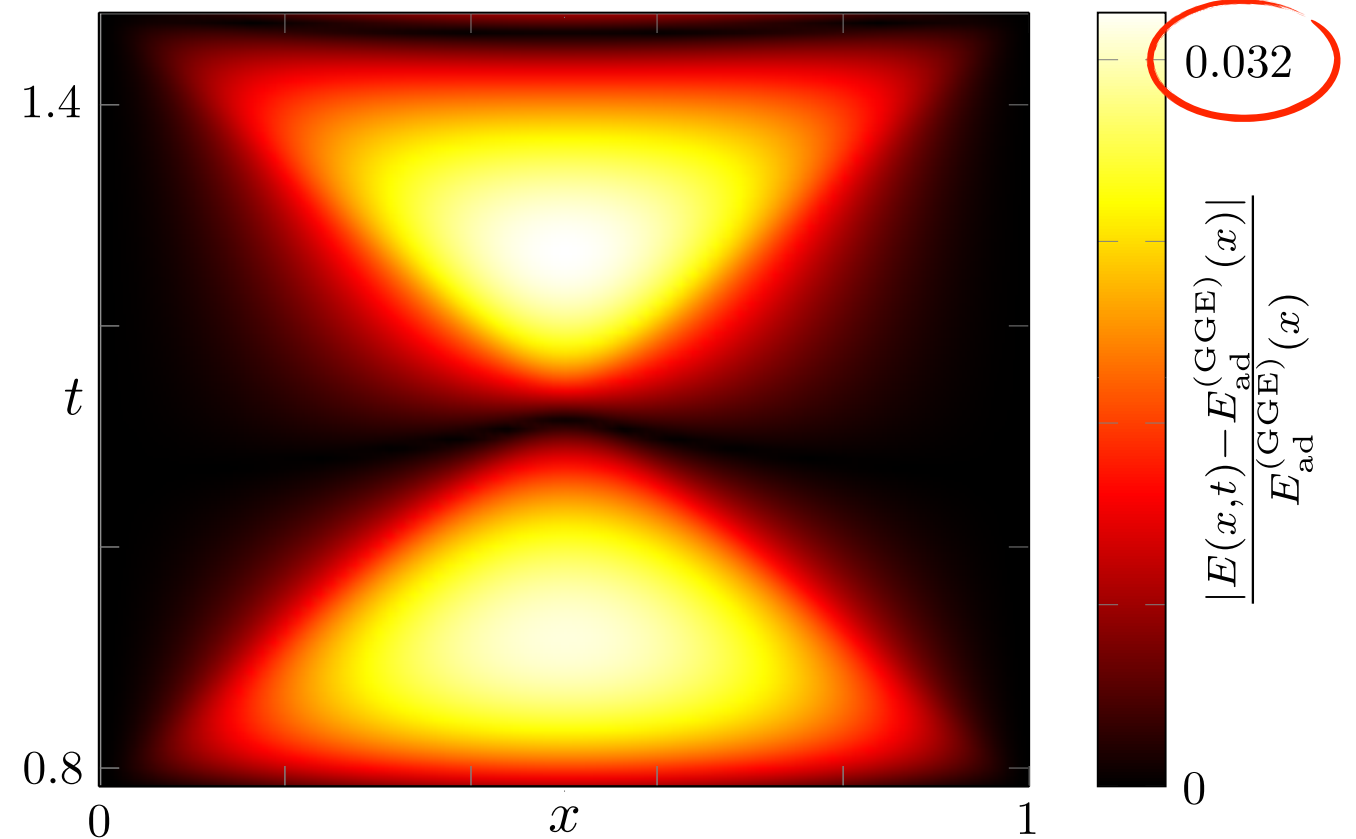
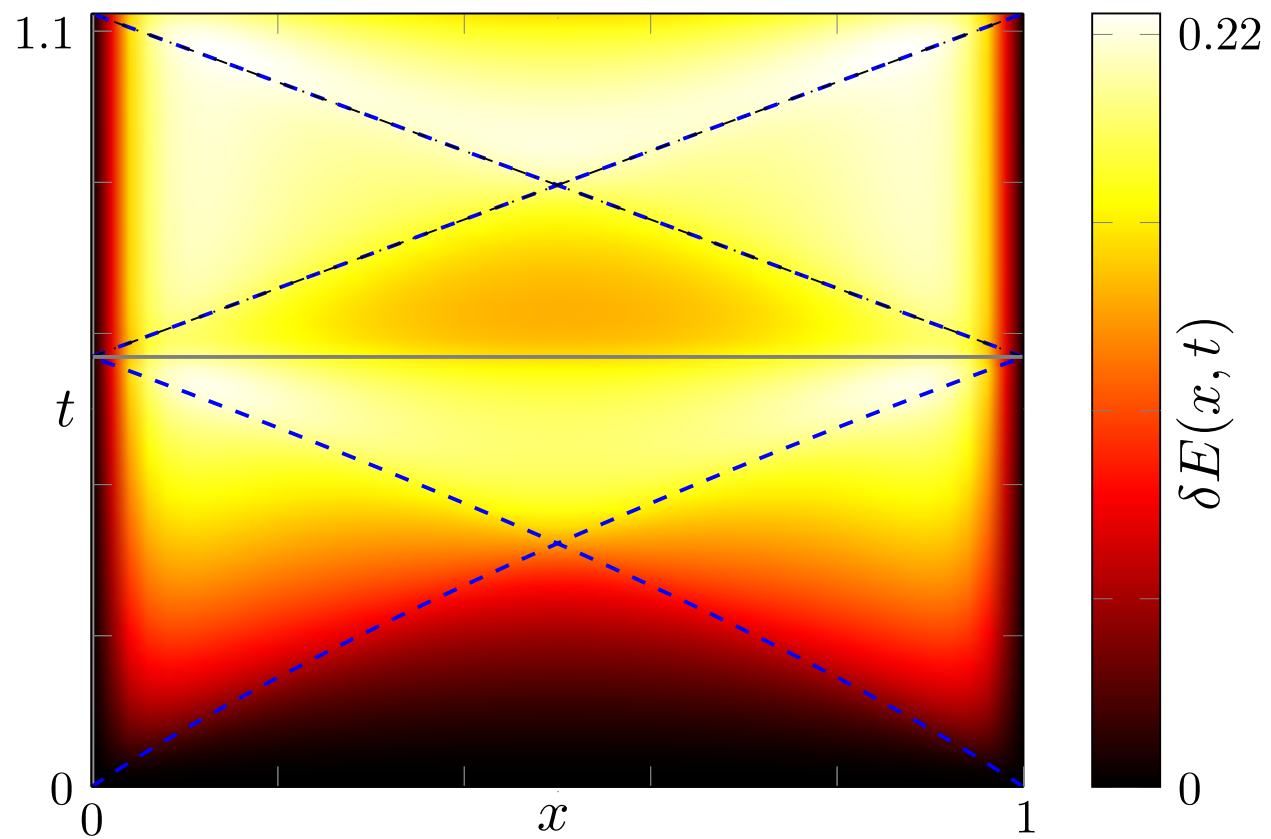
$$H(\bar{t}) = \sum_{q>0} q \bar{v}(\bar{t}) \beta_{q,\bar{t}}^\dagger \beta_{q,\bar{t}} \longrightarrow \bar{v}(t) = v_i \sqrt{1 + \frac{\eta t}{\tau}}$$

► Post-quench: **LC2** emerges at the boundaries, **LC1** goes on with velocity v_f

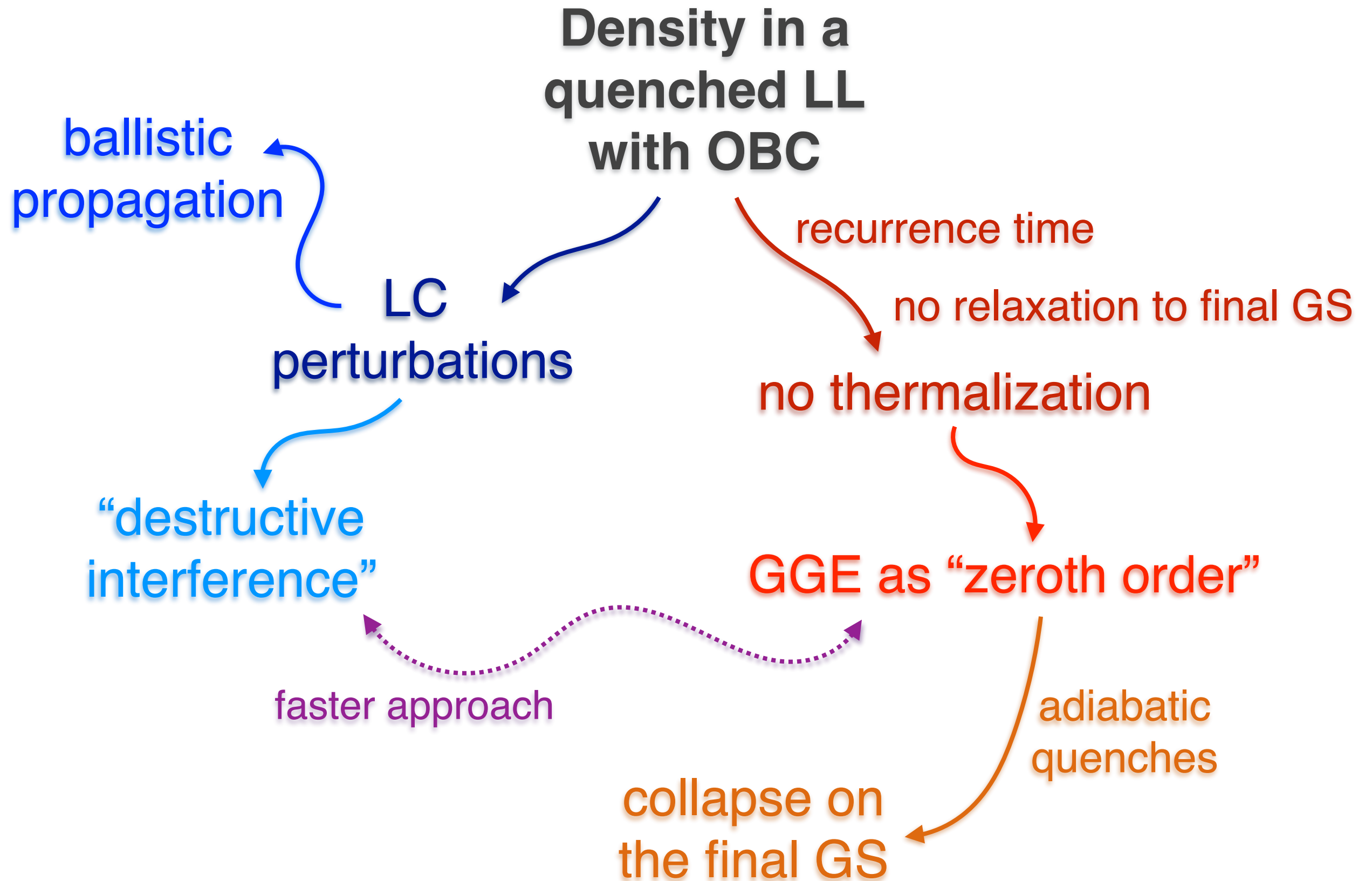
$$H_f = H(t \geq \tau) = \sum_{q>0} q v_f \beta_{q,f}^\dagger \beta_{q,f}$$

LCs move with the
instantaneous
 bosonic velocity

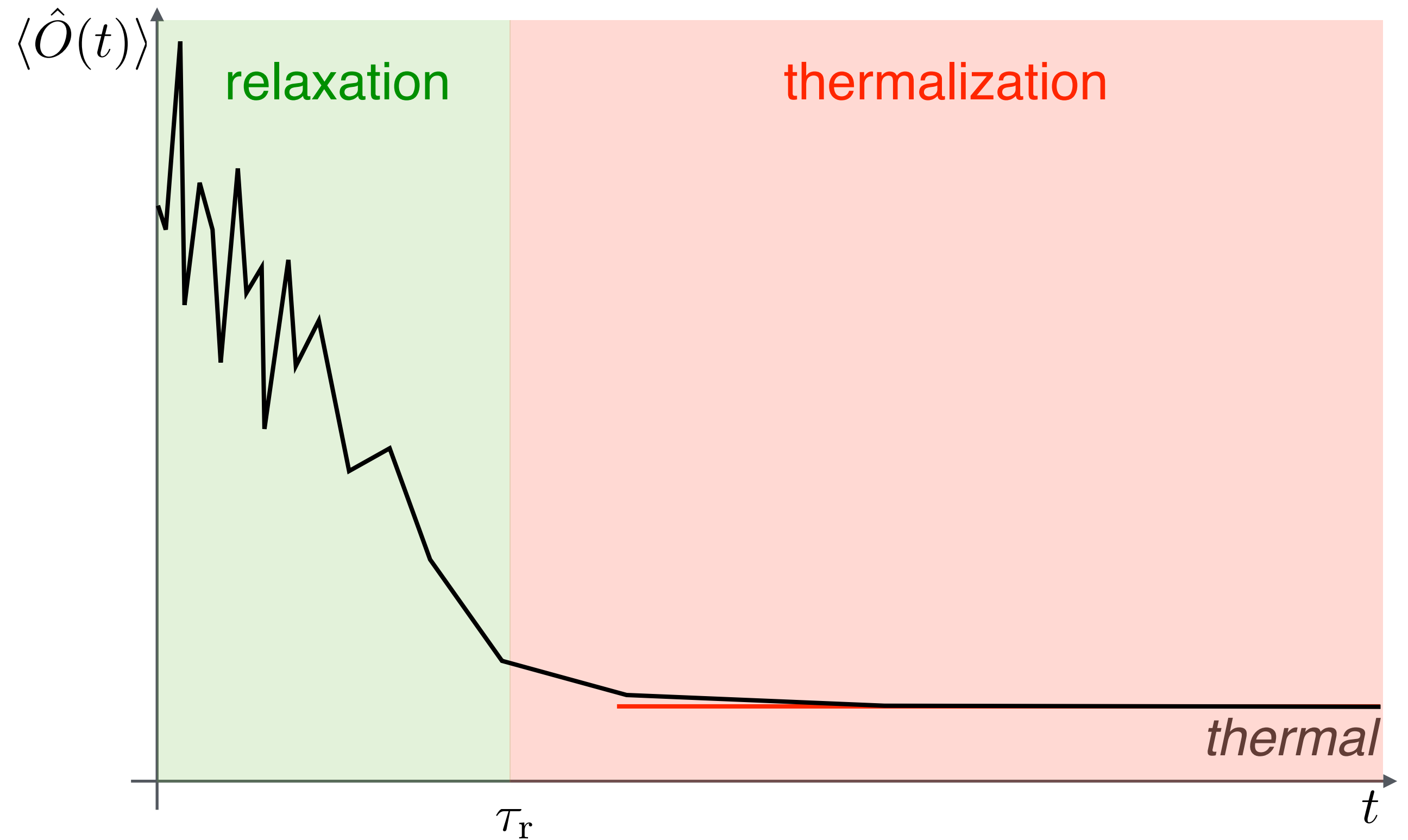
Interference between light cones



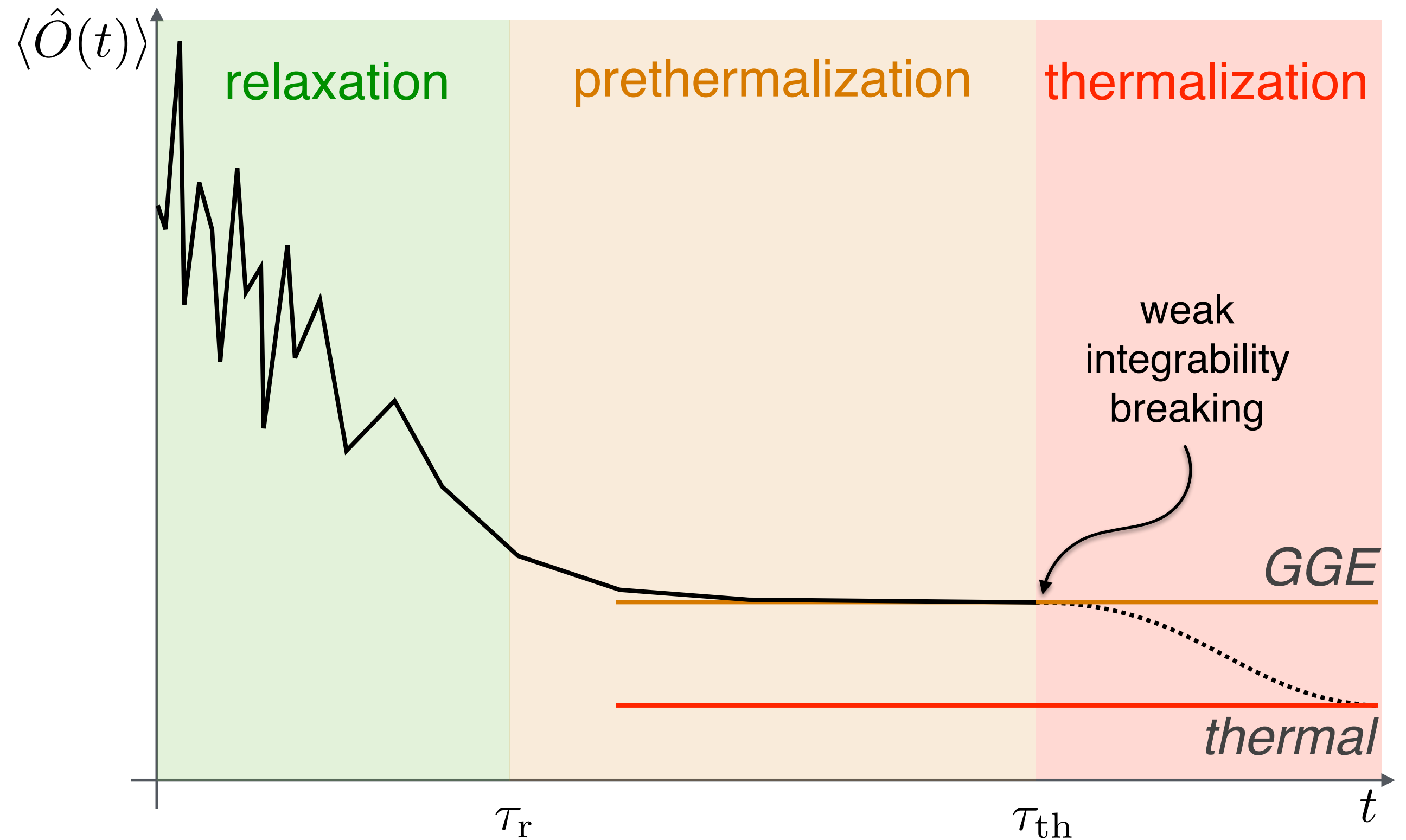
Conclusions



Thermalization



Prethermalization

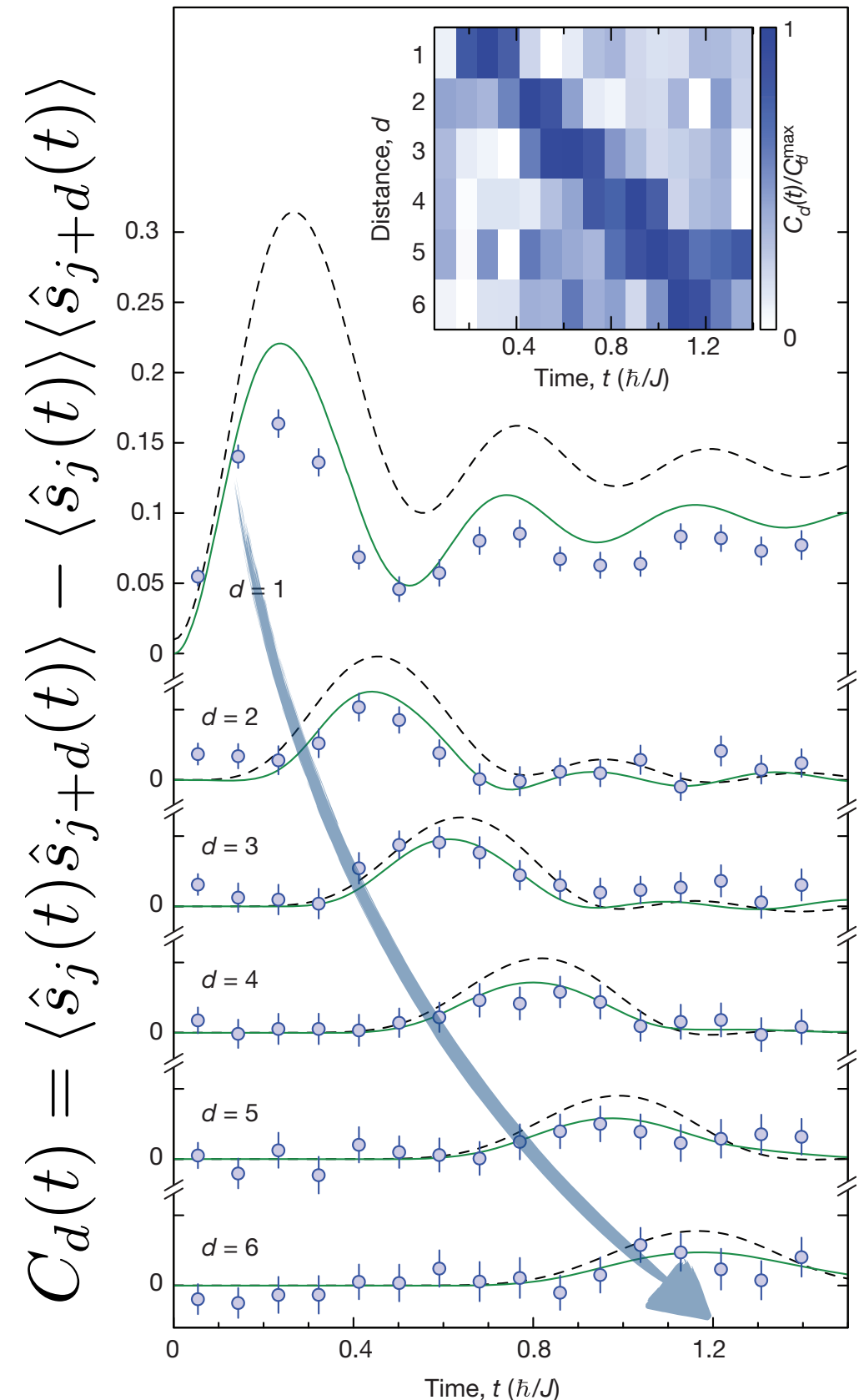
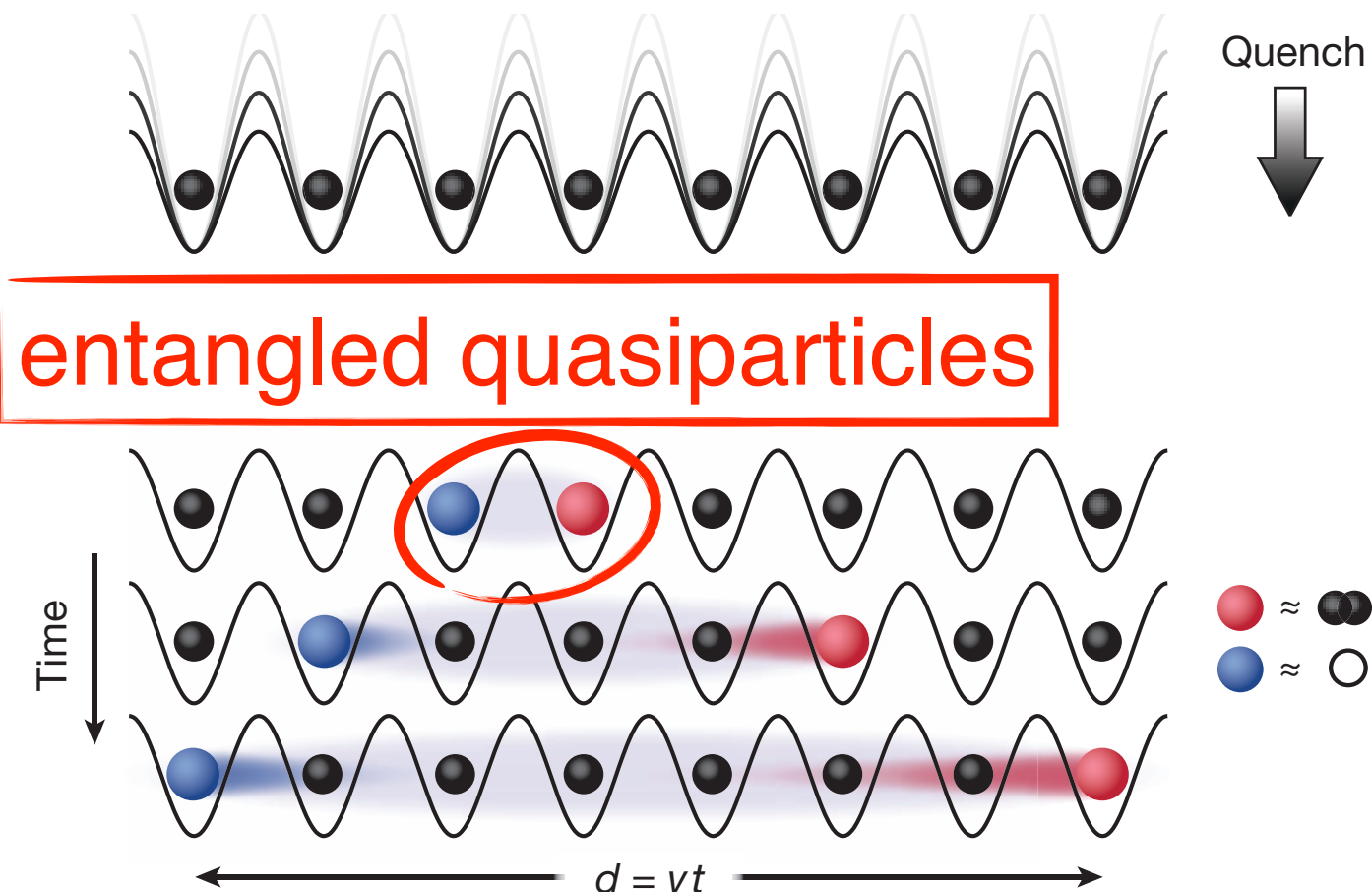


Light-cone spreading of correlations

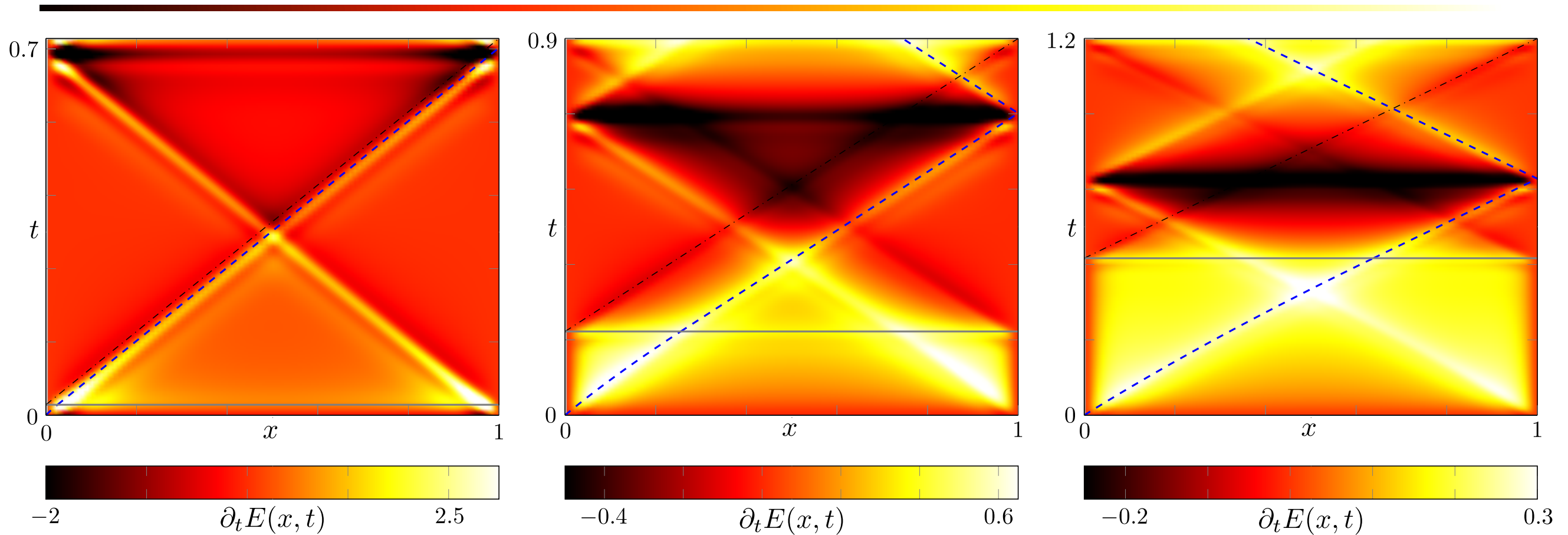
- Rel. QFT → information propagation bounded by $c \rightarrow$ **light-cone effect**

How fast “information” can spread in CM systems?

- Short-range interactions → **Lieb-Robinson bound**
- Quench in ultracold atomic gases



Where are the LCs?



► Sudden quench $E(x, t) = E_{\text{sq}}^{(\text{GGE})}(x) f_{\text{sq}}(t) \mathcal{C}_{\text{sq}}(x, v_f t)$

► Adiabatic quench ($\tau \gg \tau_{\text{ad}} \sim L|\eta|/v_i$)

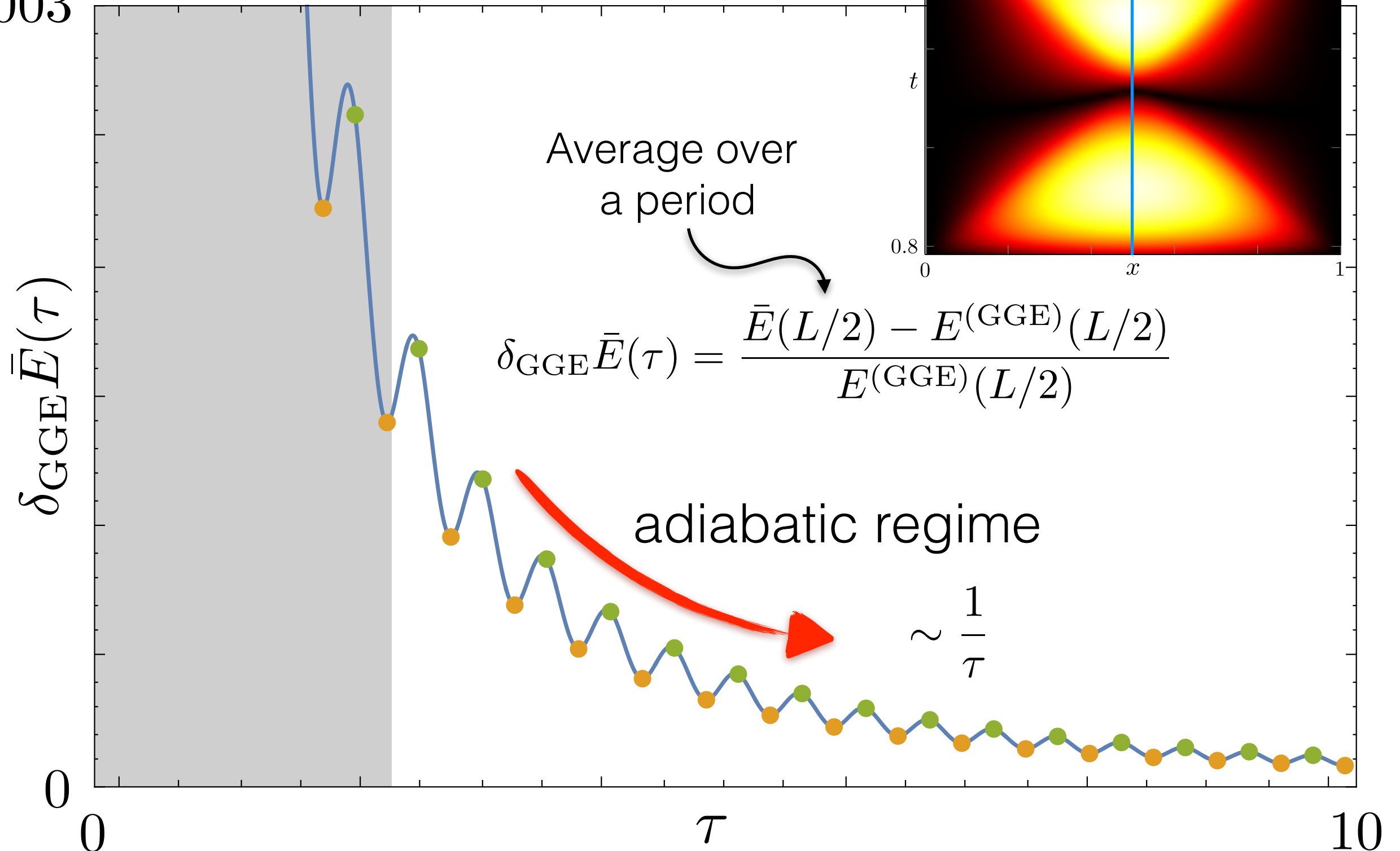
○ transient $E(x, t) = G(x, t) \mathcal{C}_{\text{ad}}(x, \ell(t))$

○ post-quench $E(x, t) = E_{\text{ad}}^{(\text{GGE})}(x) [f_{\text{ad}}(t) + A_1 \mathcal{C}_{\text{ad}}(x, v_f(t - \tau) + d) + A_2 \mathcal{C}_{\text{ad}}(x, v_f(t - \tau))]$

Approach to GGE

► How the GGE is approached?

0.003



Approach to final envelope

- ▶ How the GGE is approached?
- ▶ Is the final envelope approached by the GGE envelope?

